

Artificial gauge fields and Zitterbewegung in a BEC

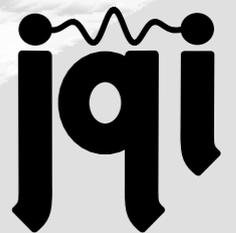
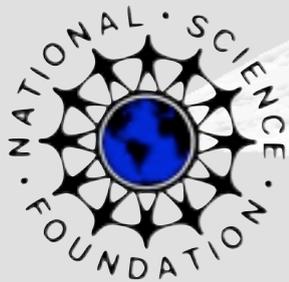
I. B. Spielman

Current team

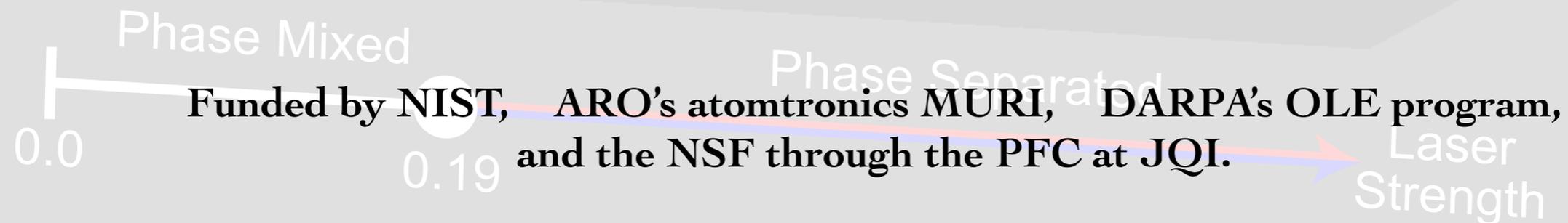
L. J. LeBlanc, M. C. Beeler, R. A. Williams, K. Jiménez-García, and A. R. Perry

Senior coworkers

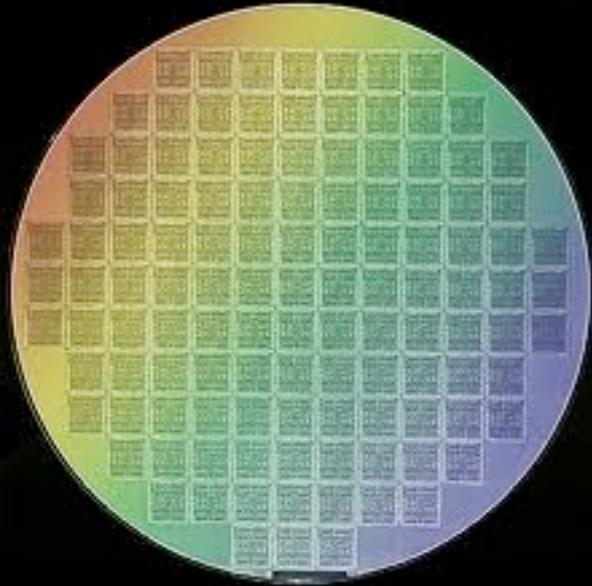
J. V. Porto, and W. D. Phillips



NIST

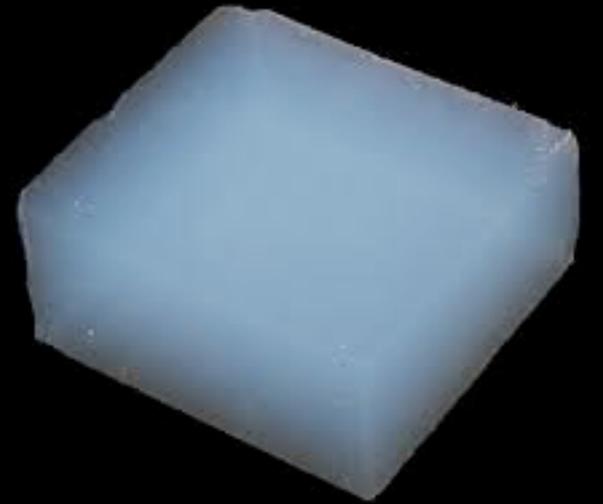


What are materials?



Si
 2.3 g/cm^3

Ian's answer: "chunks of stuff."



Aerogel
 1 mg/cm^3

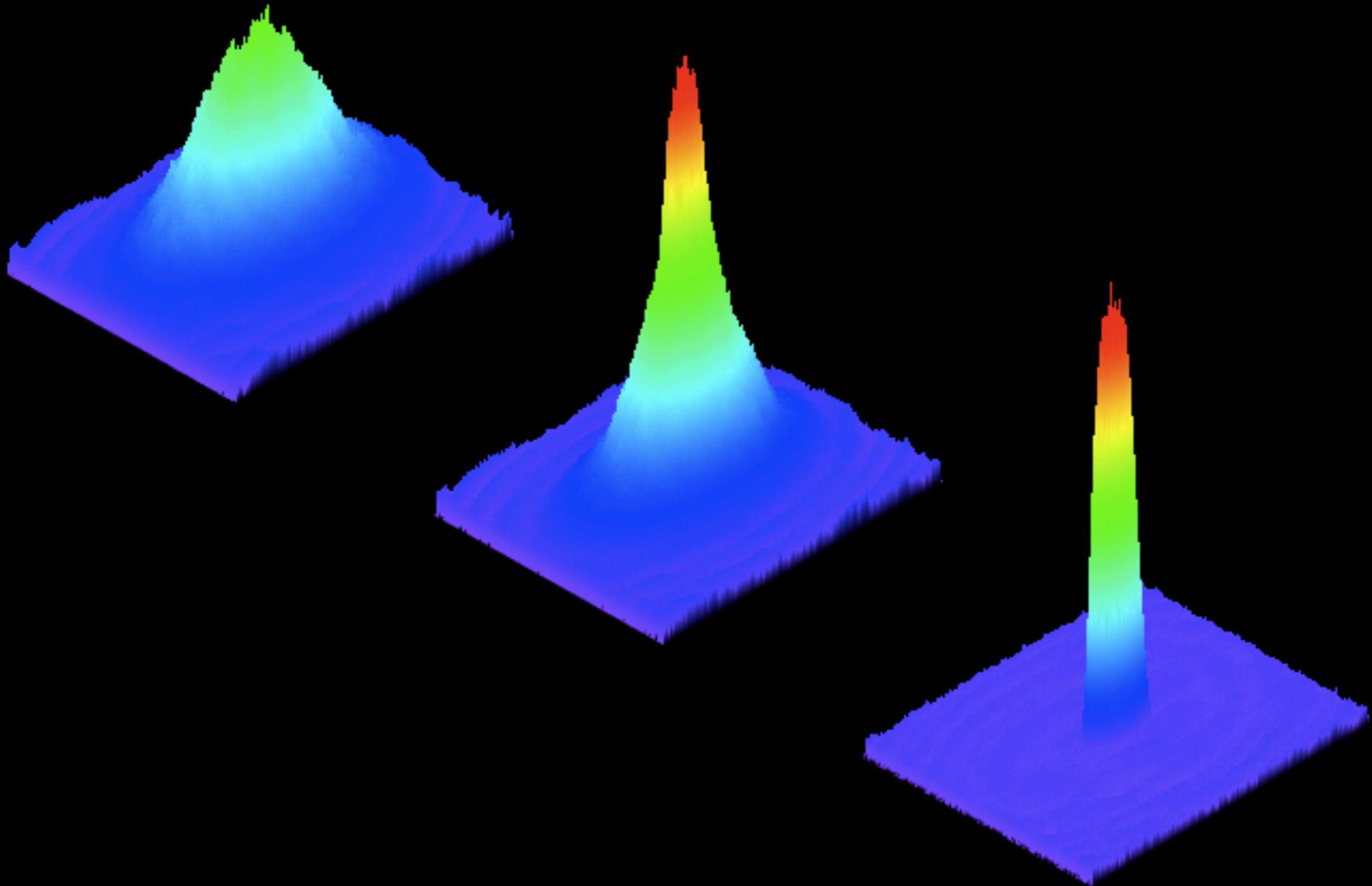
Liquid Helium
 125 mg/cm^3

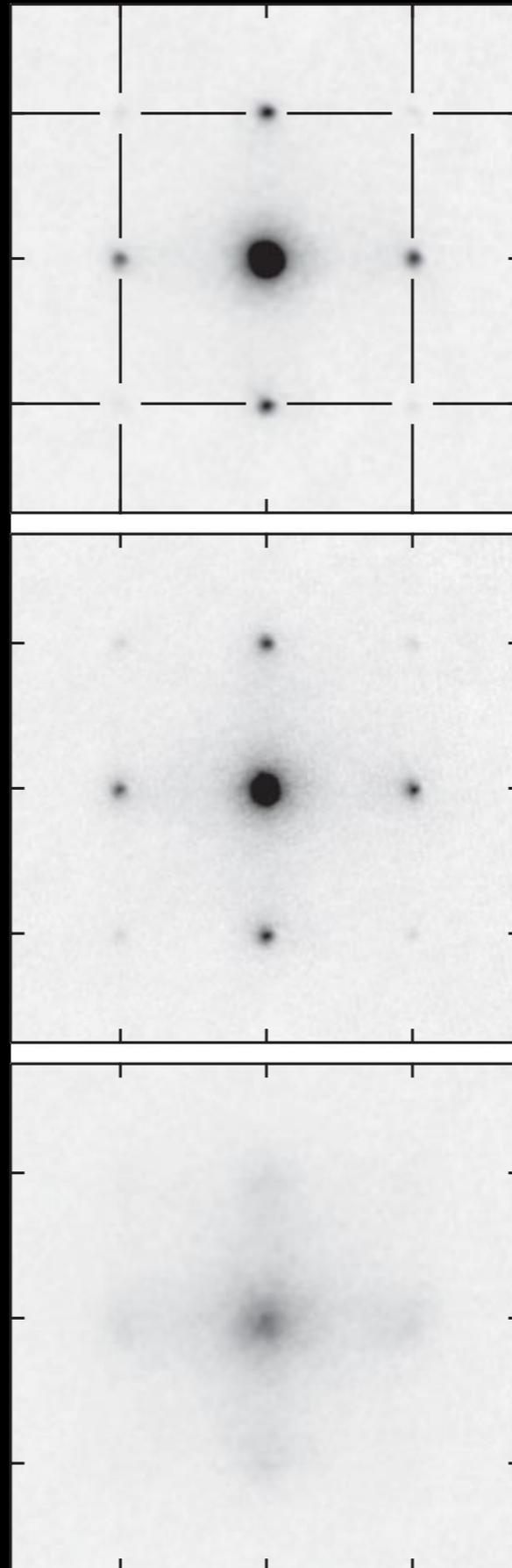
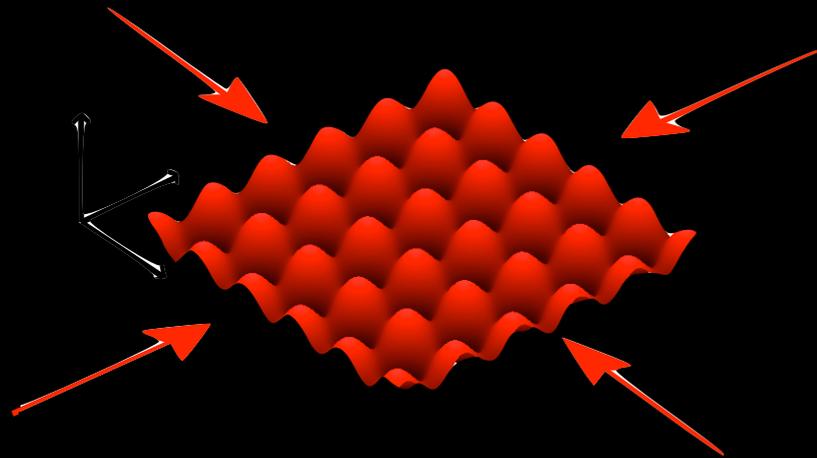


Ultracold neutral atoms

$\sim 10^{14} \text{ cm}^{-3}$ or 100 ng/cm^3
(air is $\sim 1 \text{ mg/cm}^3$)

Are these materials?





They can be fluids

They can be insulators

They can be bosons

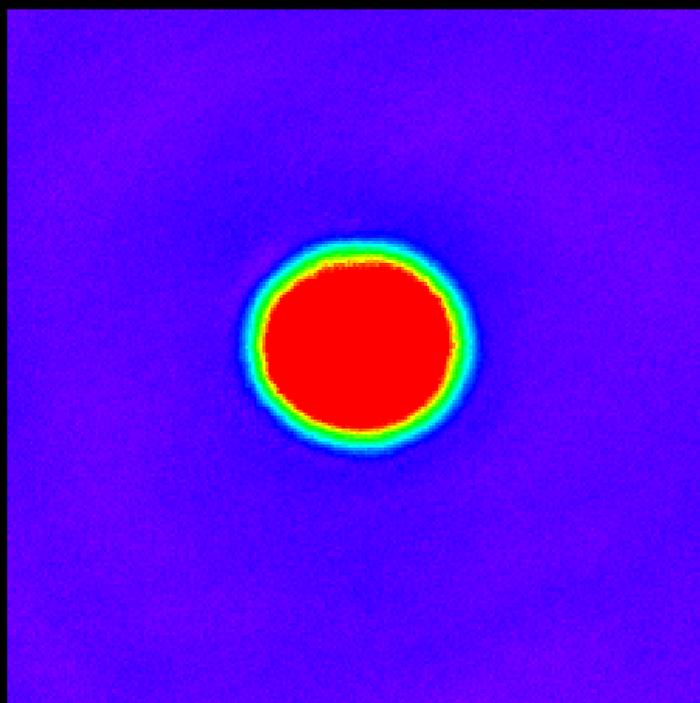
They can be fermions



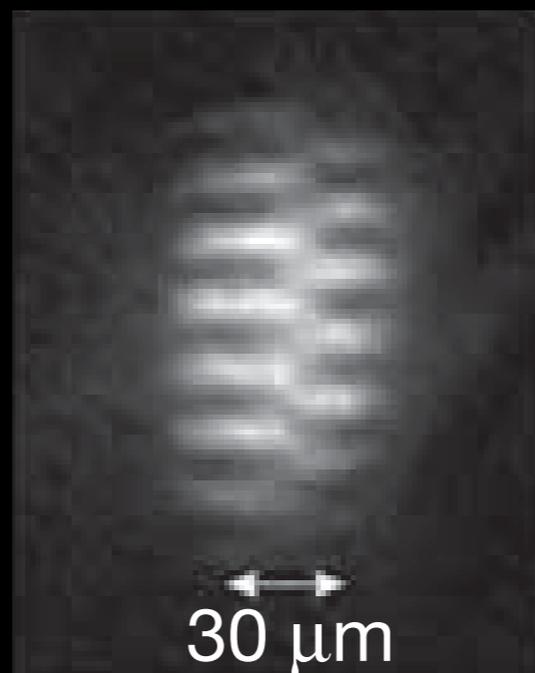
They can be molecules

They can be atoms

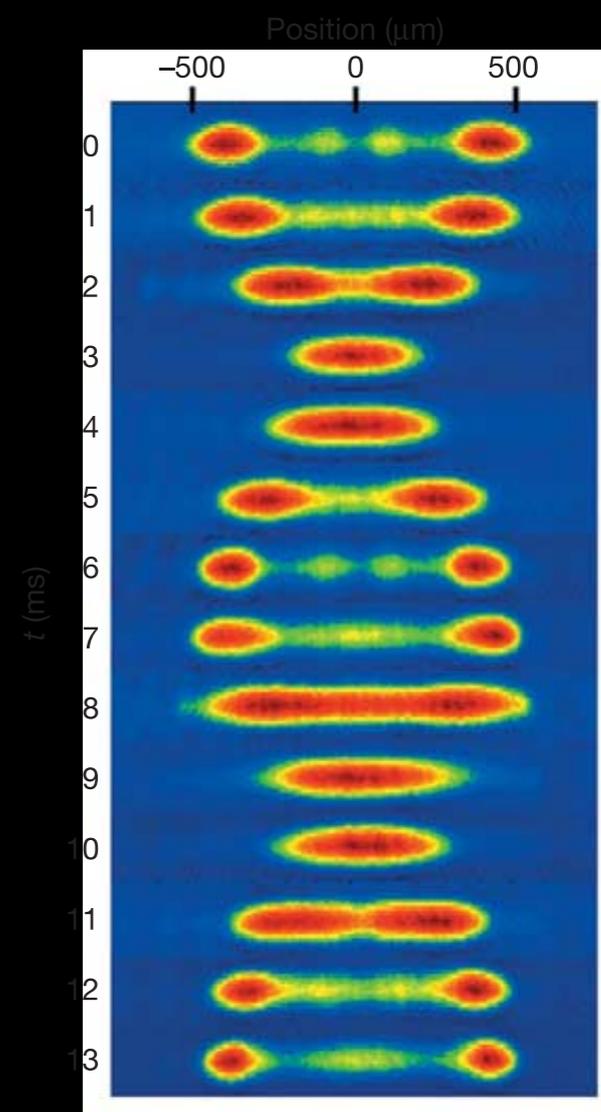
They can be 3D



They can be 2D



They can be 1D



e.g., Hadzibabic Nature (2006)

e.g., Kinoshita Nature (2006)

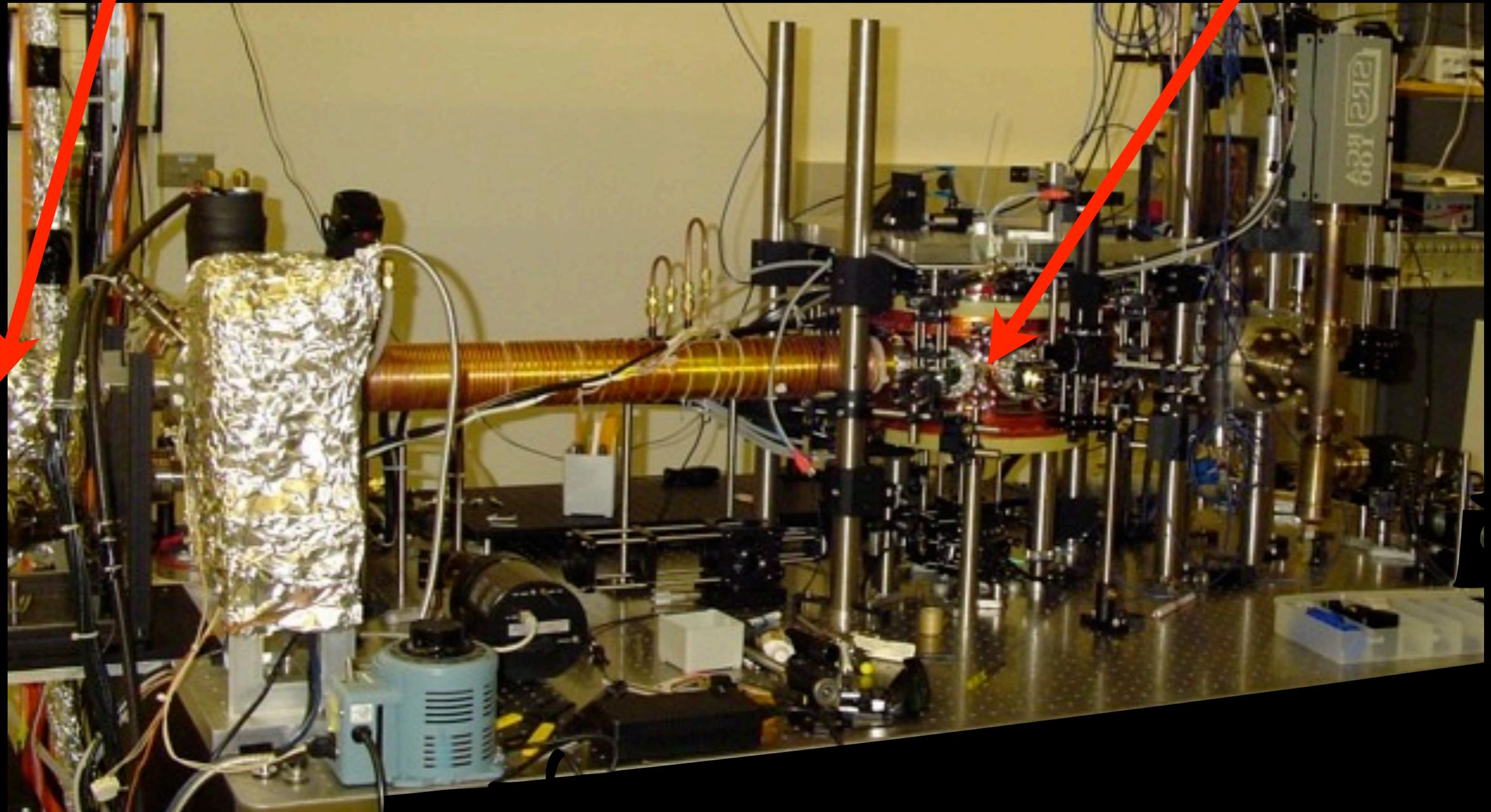
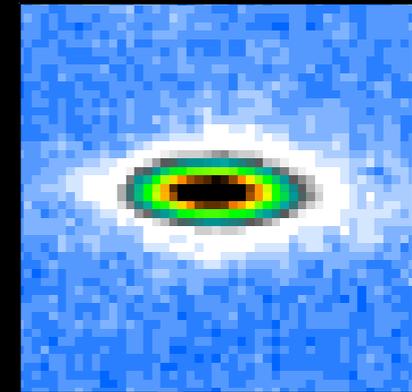
Starts like this

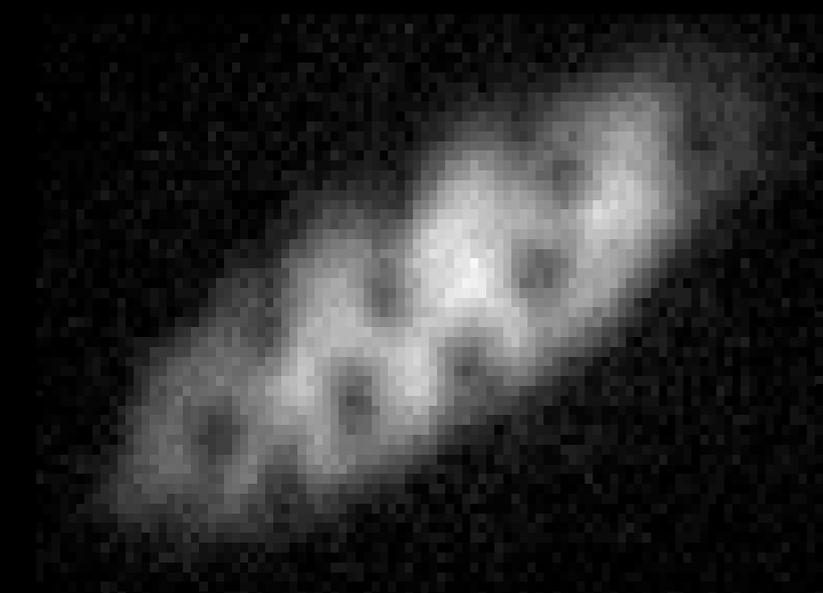
400 K



Ends here every 20 s

50 nK

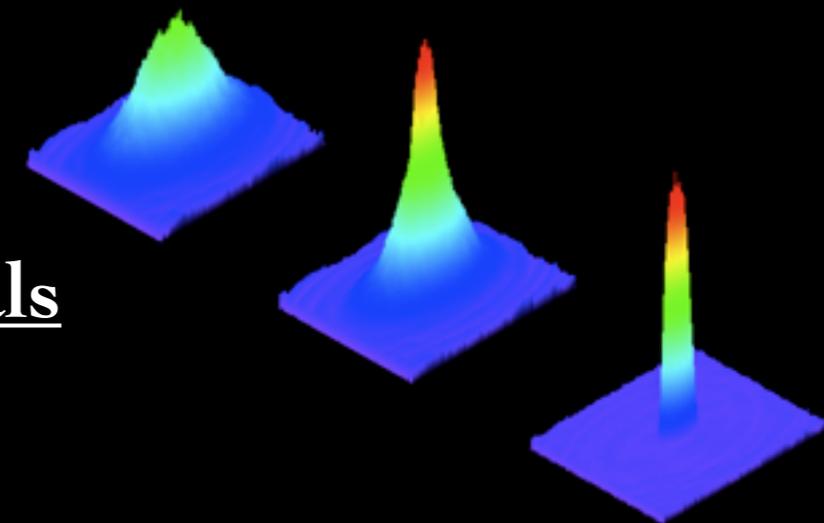




Cold atoms are good materials

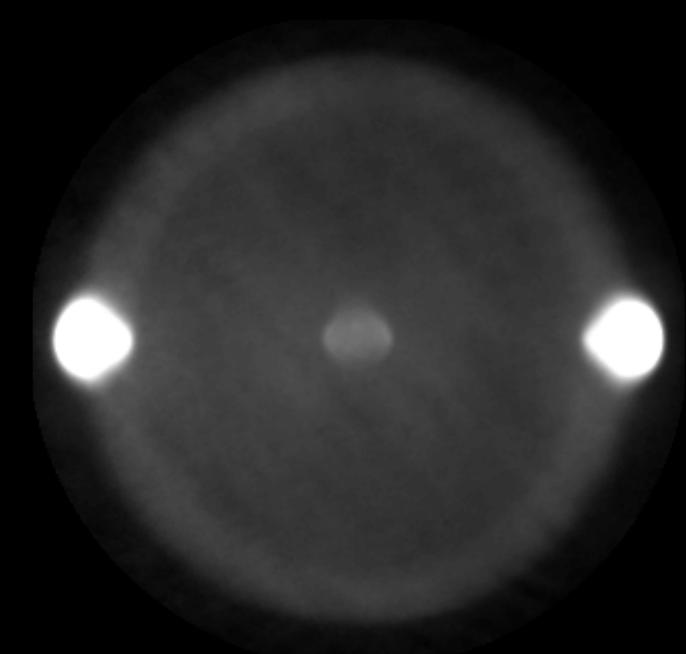
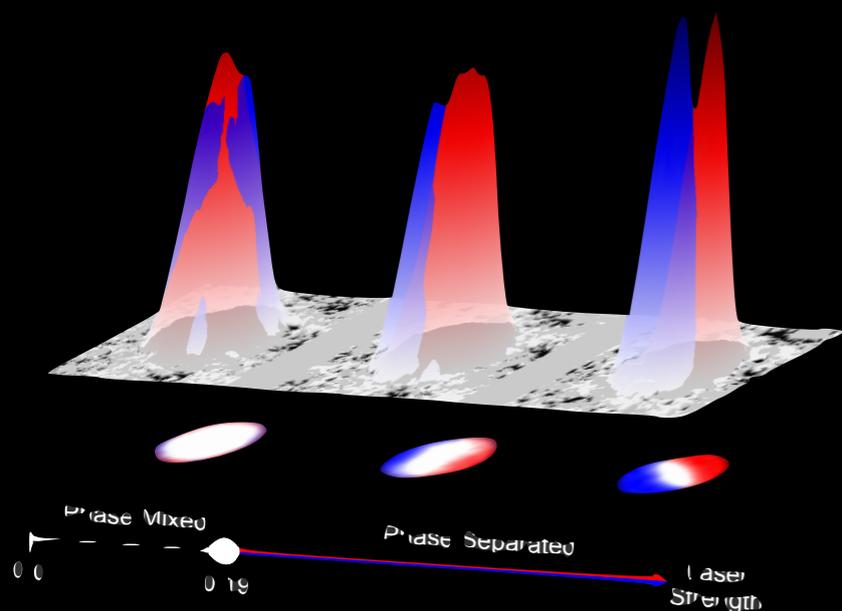
Numerous properties can be controlled and measured on all relevant timescales and in any lab

Very simple Hamiltonians

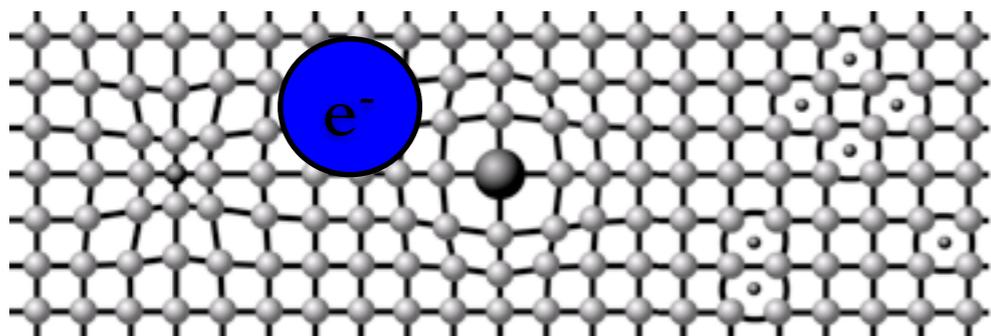


Cold atoms are bad materials Short lived, and do so in vacuum

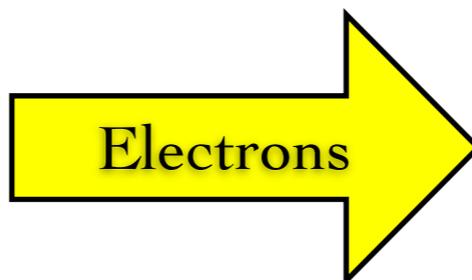
Interesting features all added by hand (complex experiments).



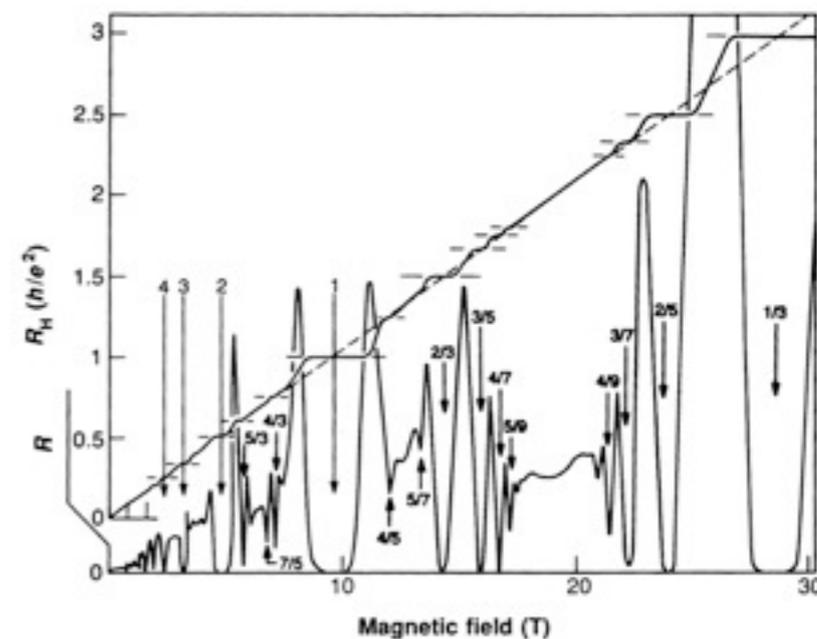
Electron



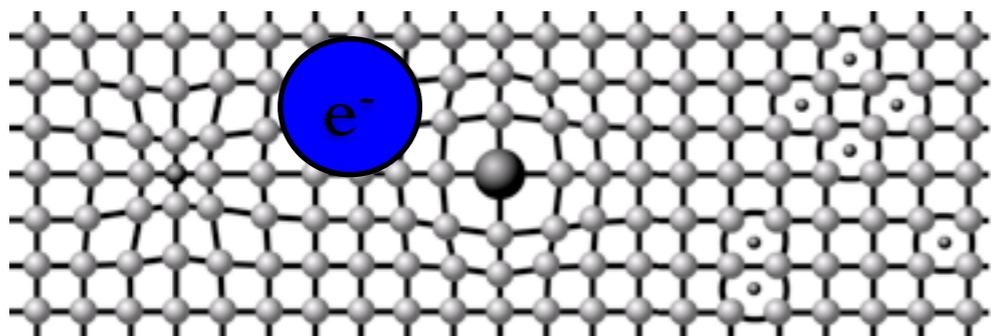
ab initio
understand is hard



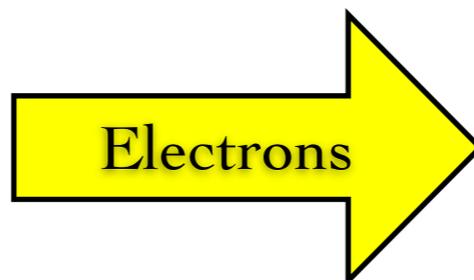
Materials



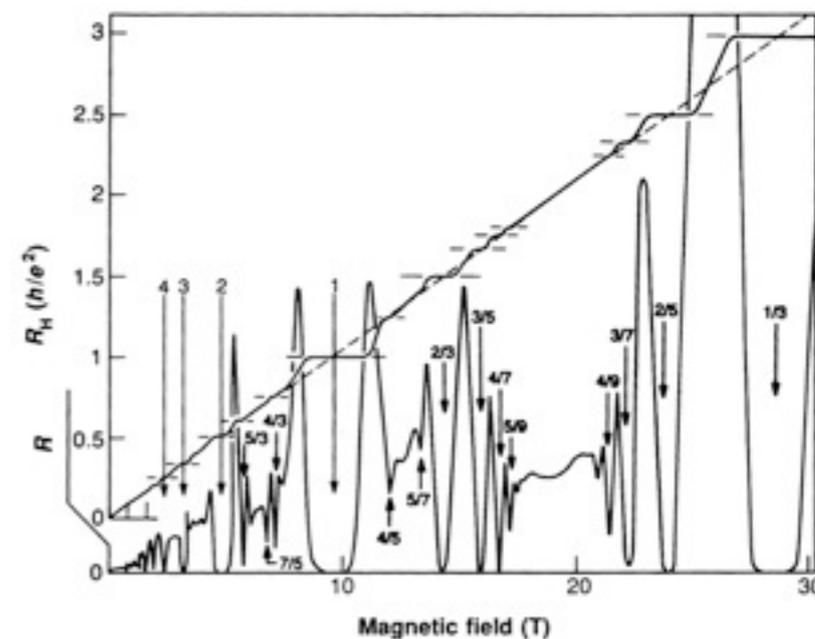
Electron



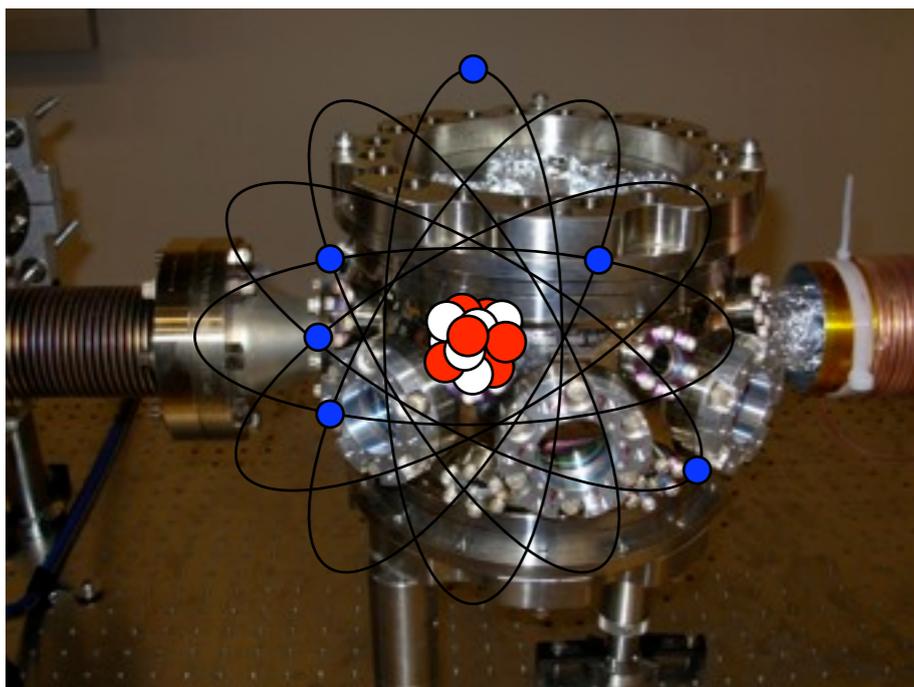
ab initio
understand is hard



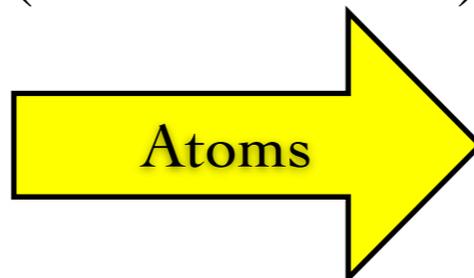
Materials



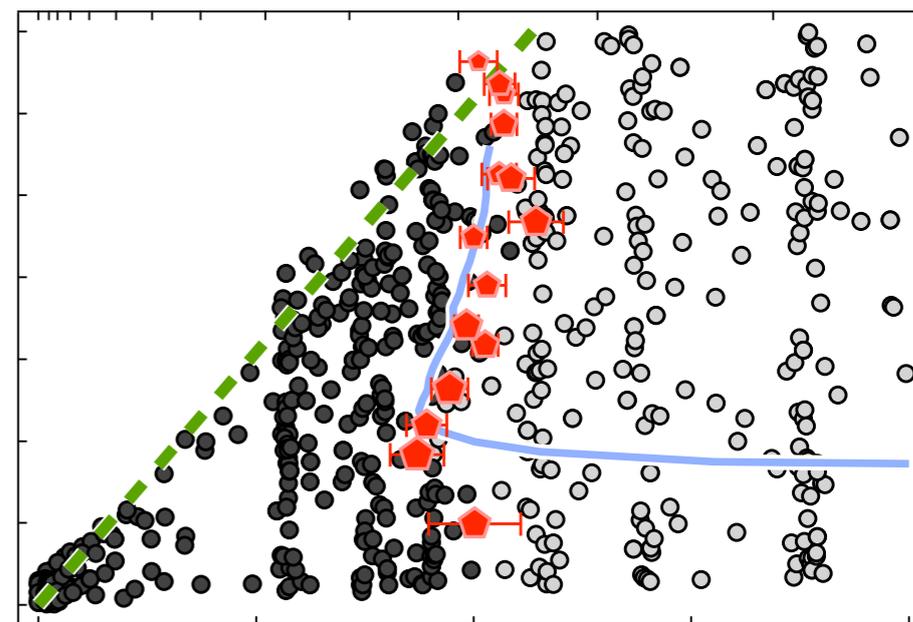
Atom



ab initio
is easier
(but still hard)



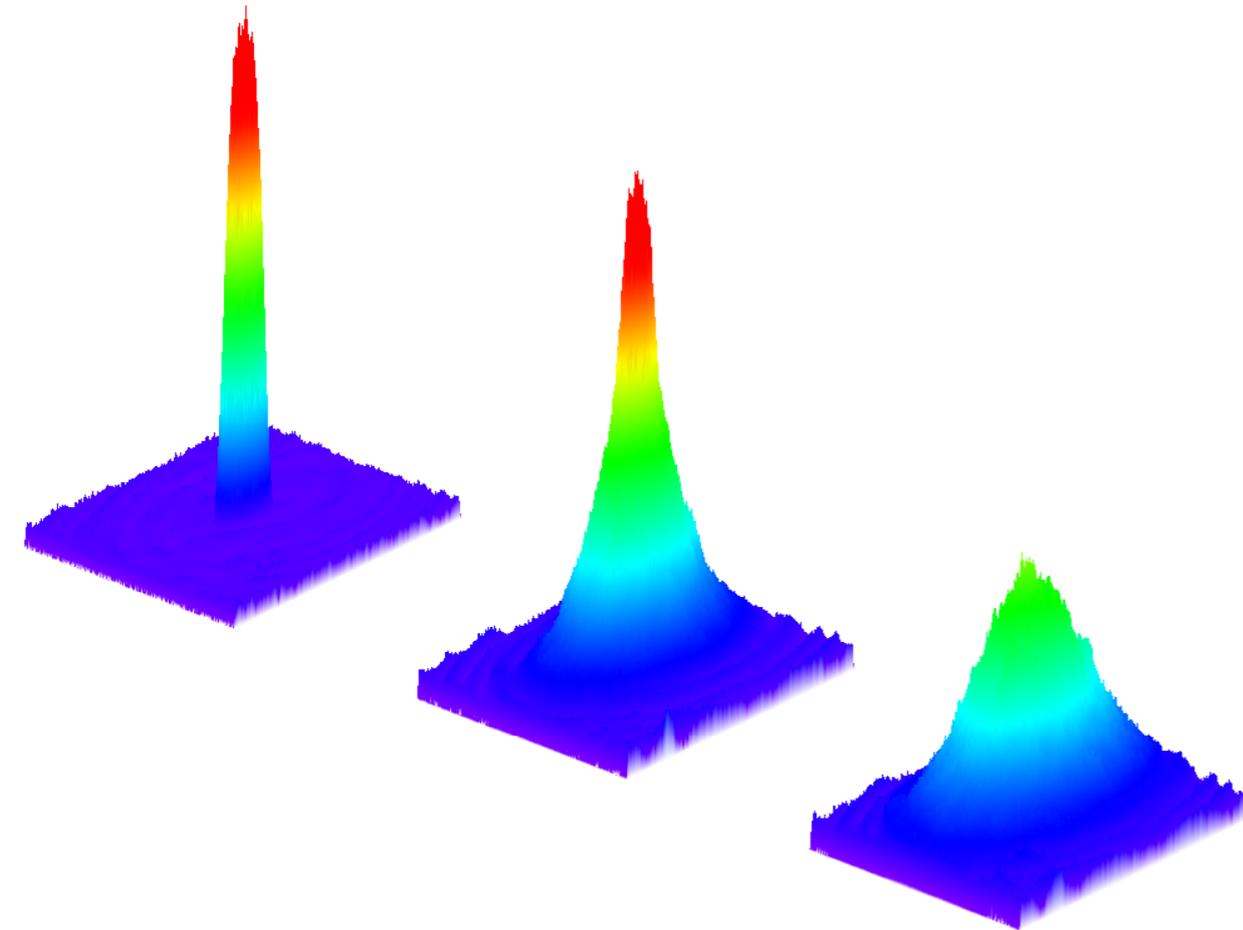
Atomic quantum materials



Vision: atoms + fields

Bose-Einstein condensation

Quantum Hall effects

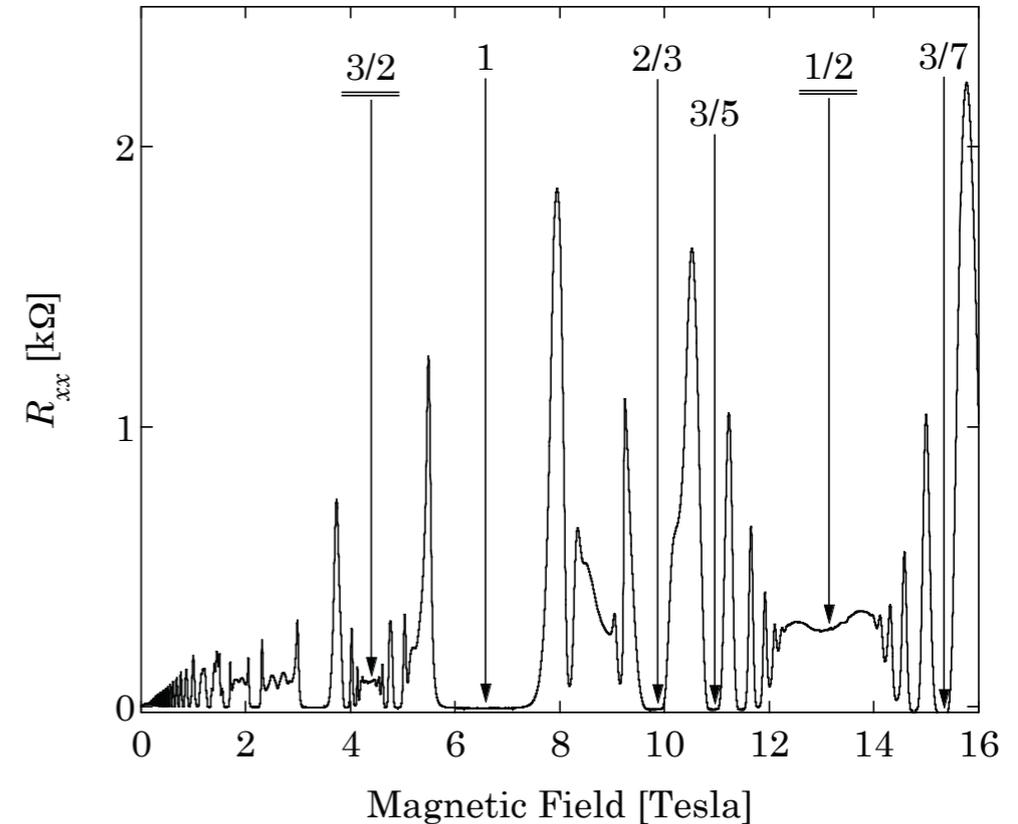


Charge neutral bosons

Macroscopically occupied quantum state

Weak interactions

Phonon-like excitations, bosons



Charged fermions in a large magnetic field

Macroscopically degenerate ground state

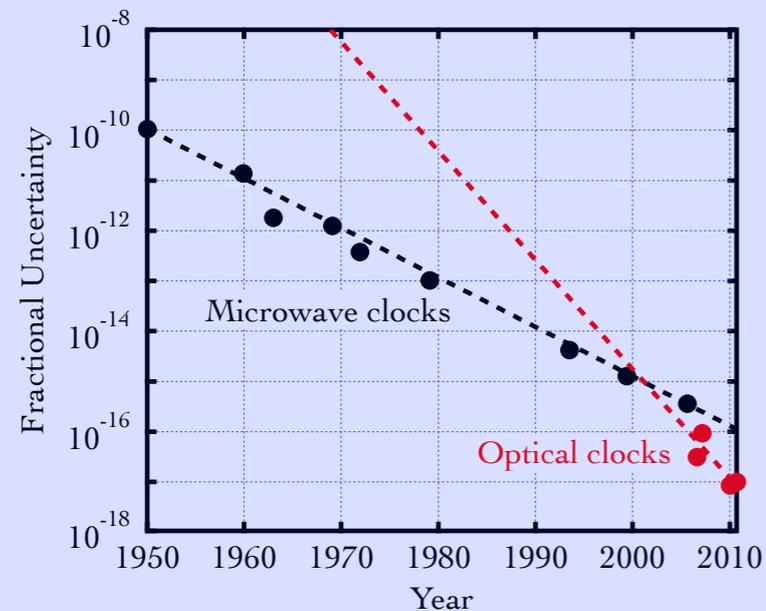
Strong interactions

Particle-like excitations with fractional quantum numbers

A brief history of atomic physics

Reduction

Very good understanding an atom



Assemble new systems

Make new quantum systems of **Atoms**

Control of external states

Many-body
physics

Quantum
simulation

**Origin of complexity
from simple components**

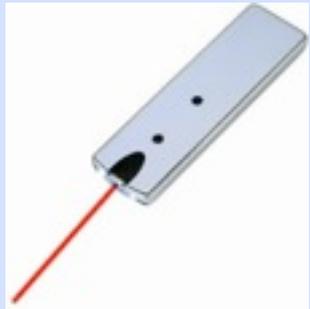
Quantum
information

Metrology

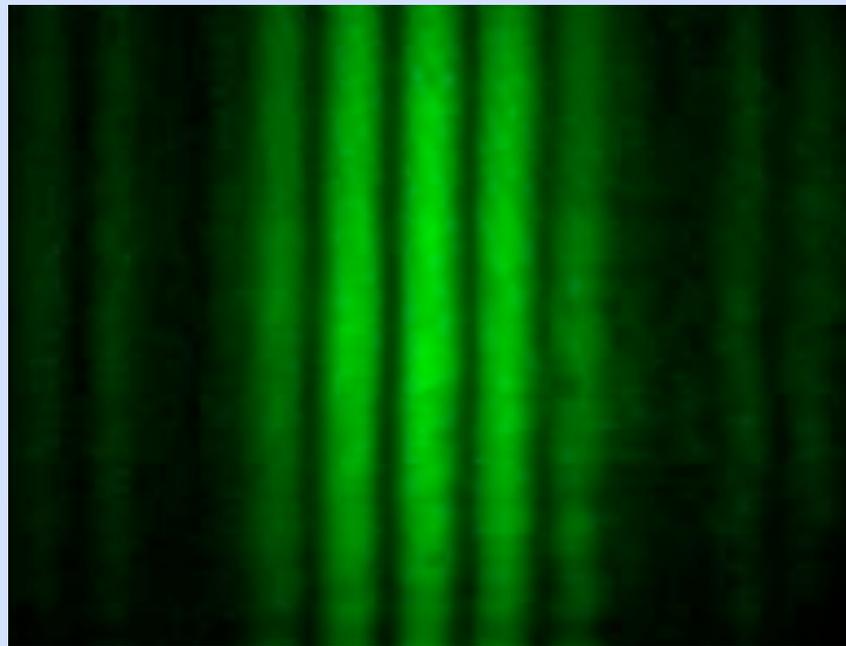
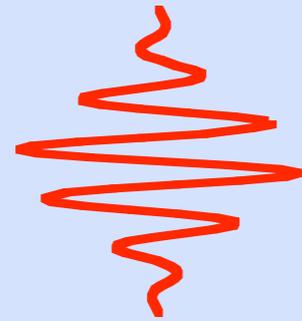
Control of internal states

Quantum mechanics: interference

Light
massless



=



Electrons

$$m = 9.8 \times 10^{-31} \text{ kg}$$

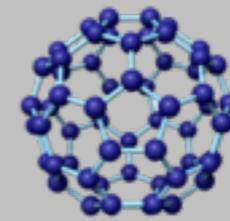
e^-

=

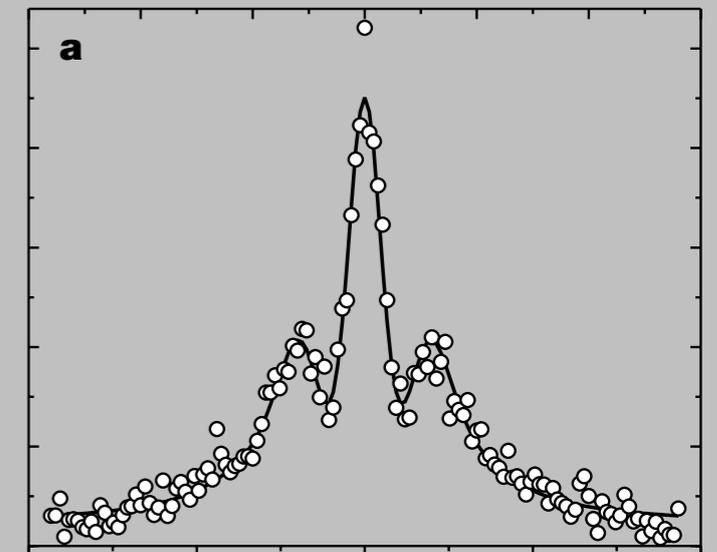


C-60

$$m = 1.2 \times 10^{-24} \text{ kg}$$



=

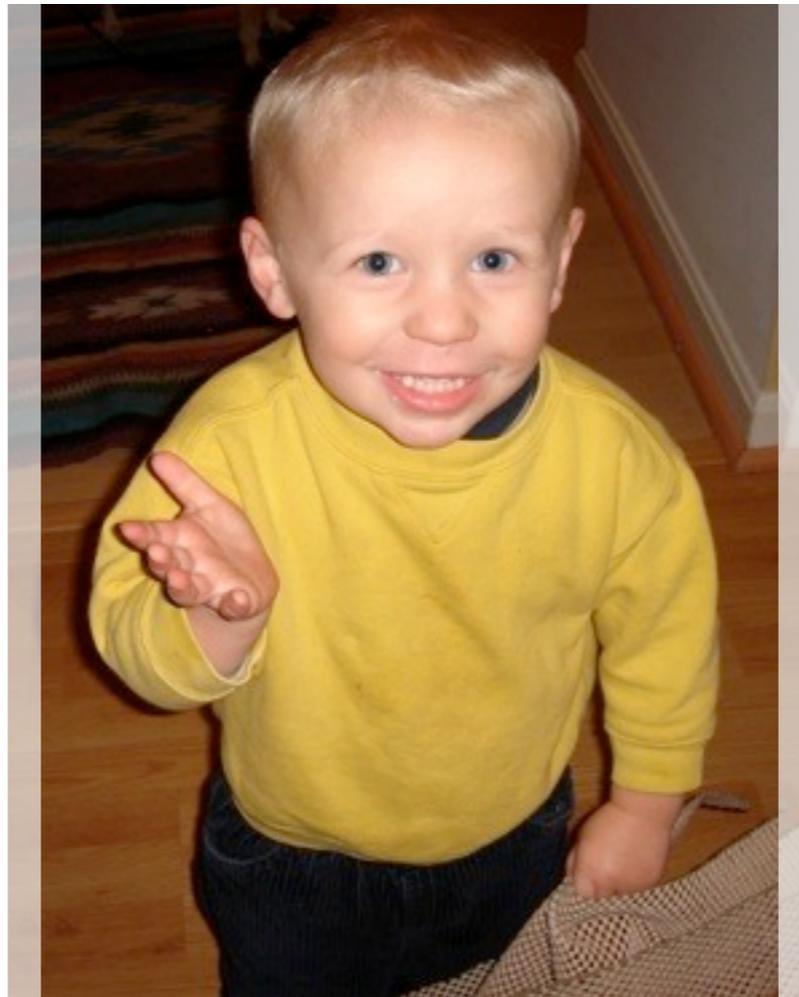


Quantum mechanics: when is it?

$$\Delta v = \frac{2\pi\hbar}{m\Delta x}$$

The uncertainty in the position of a 100 kg person at 1 m/s is just 7 nano-nano-nano-nano meters (7×10^{-36} meters).

For people this just doesn't matter (even a rapidly moving 15 kg one)



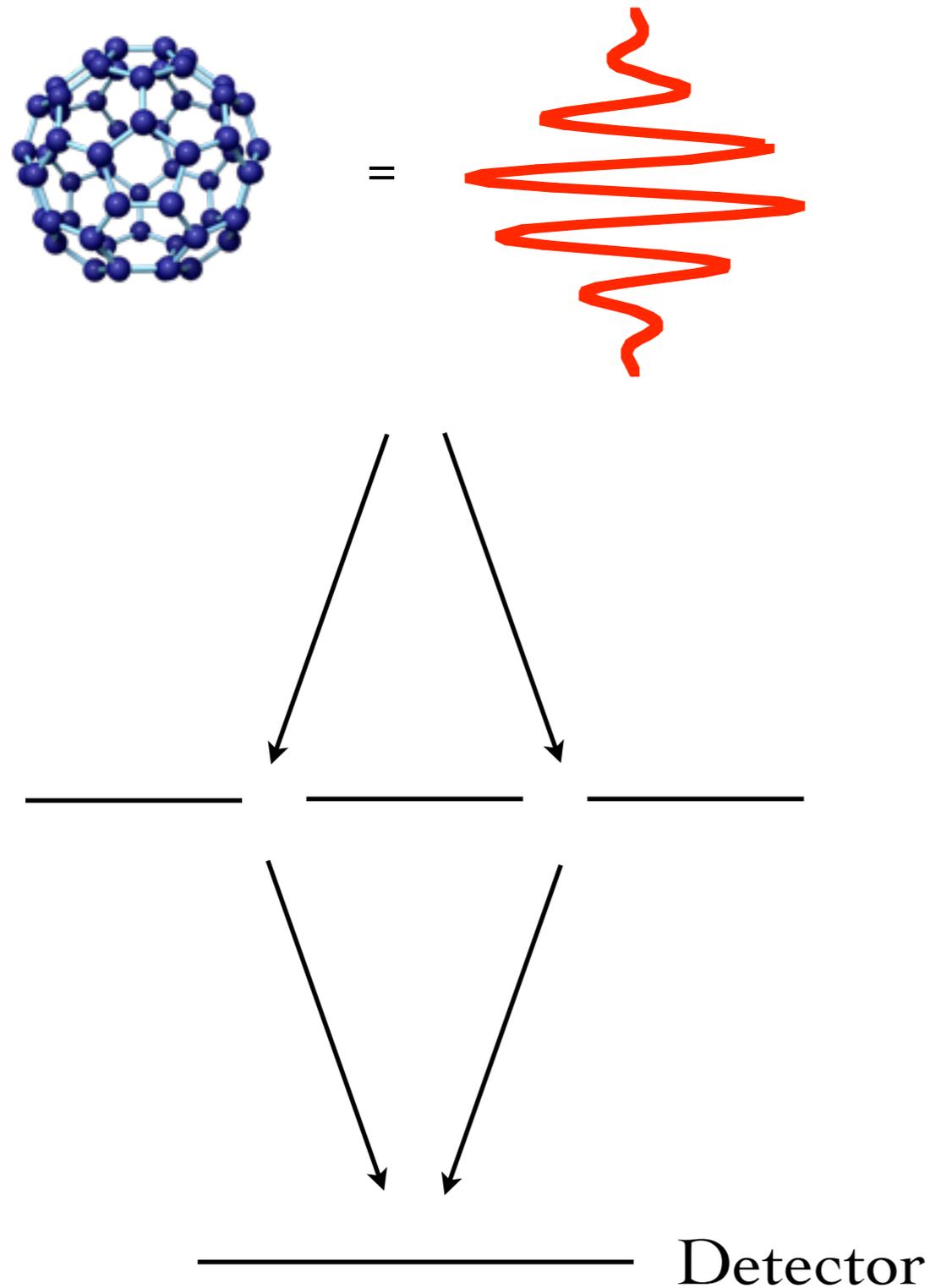
Back to quantum mechanics

Randomness

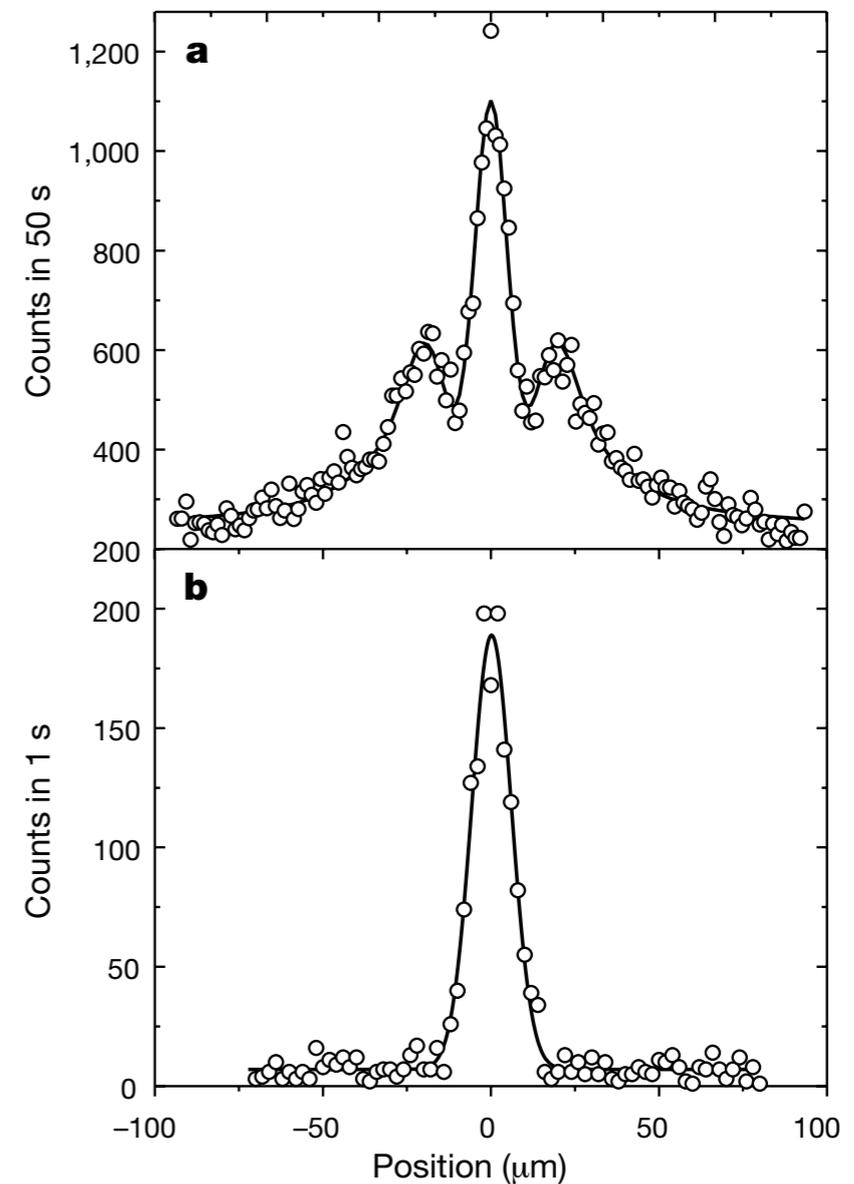
Quantum mechanics is a full deterministic theory (no randomness) **until** measurement



Quantum mechanics: interference (IV)



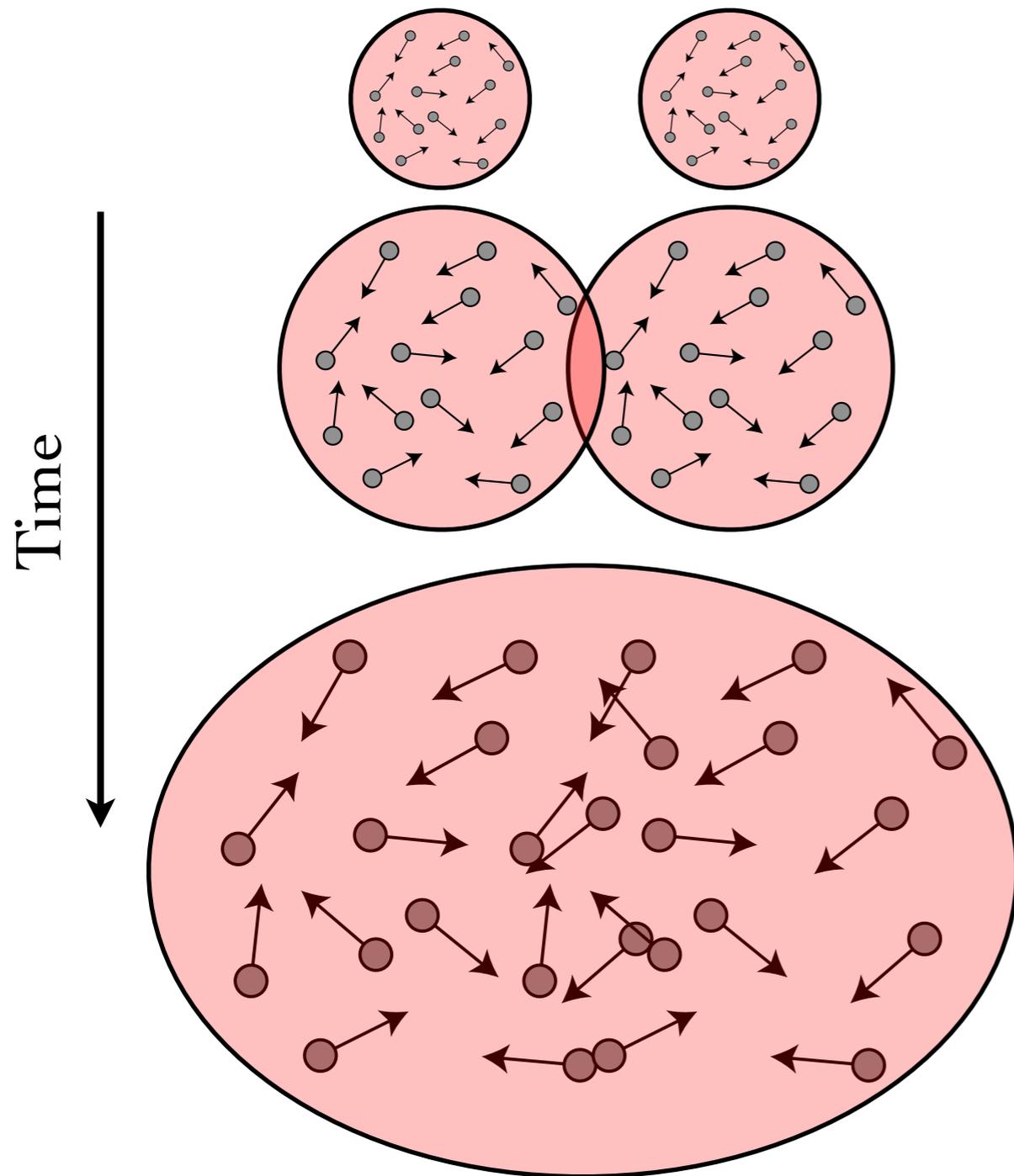
Interference of *individual* buckyballs.
One at a time.



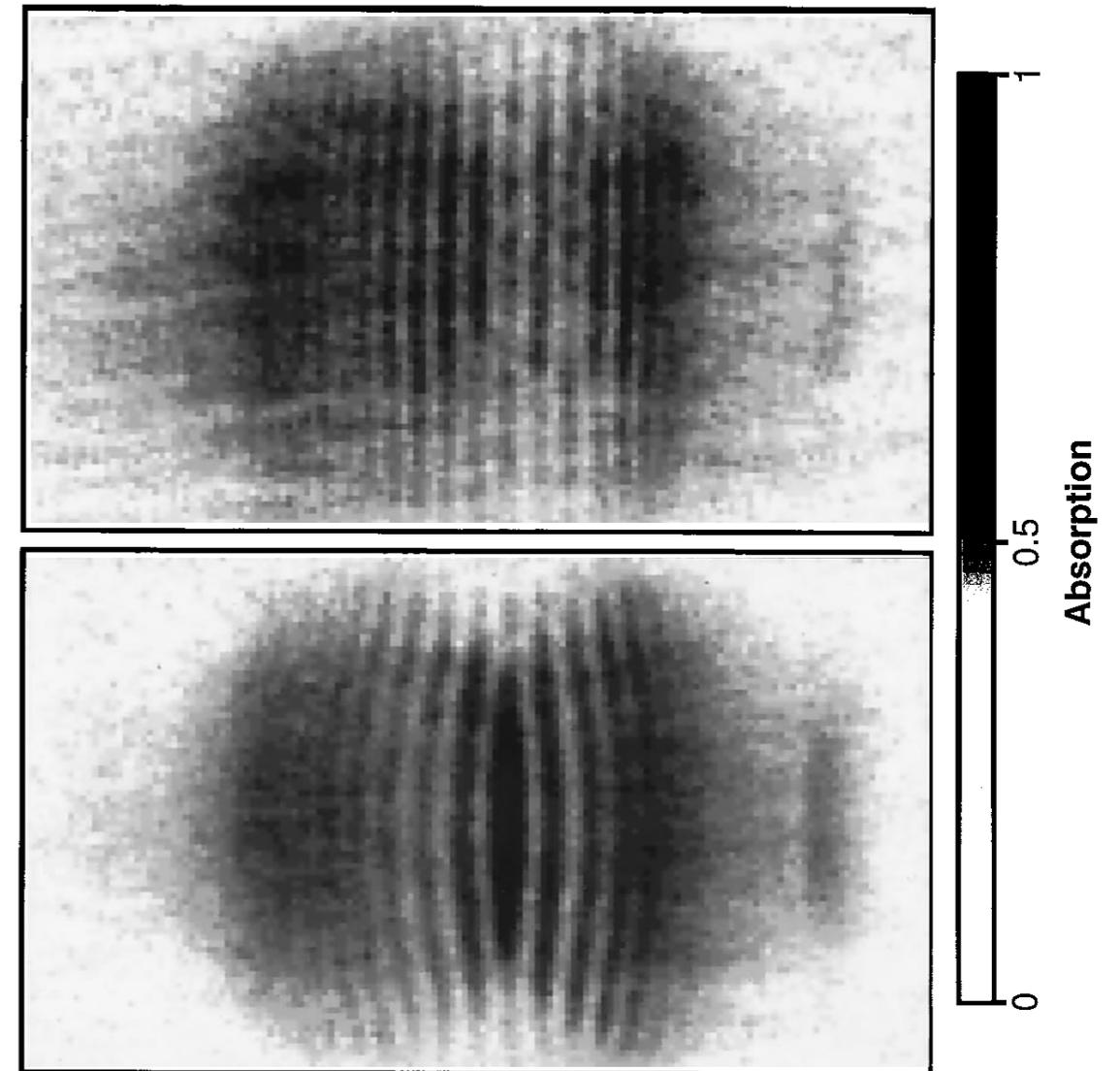
Arndt et al. Nature (1999)

Quantum mechanics: interference (II)

Two usual gases

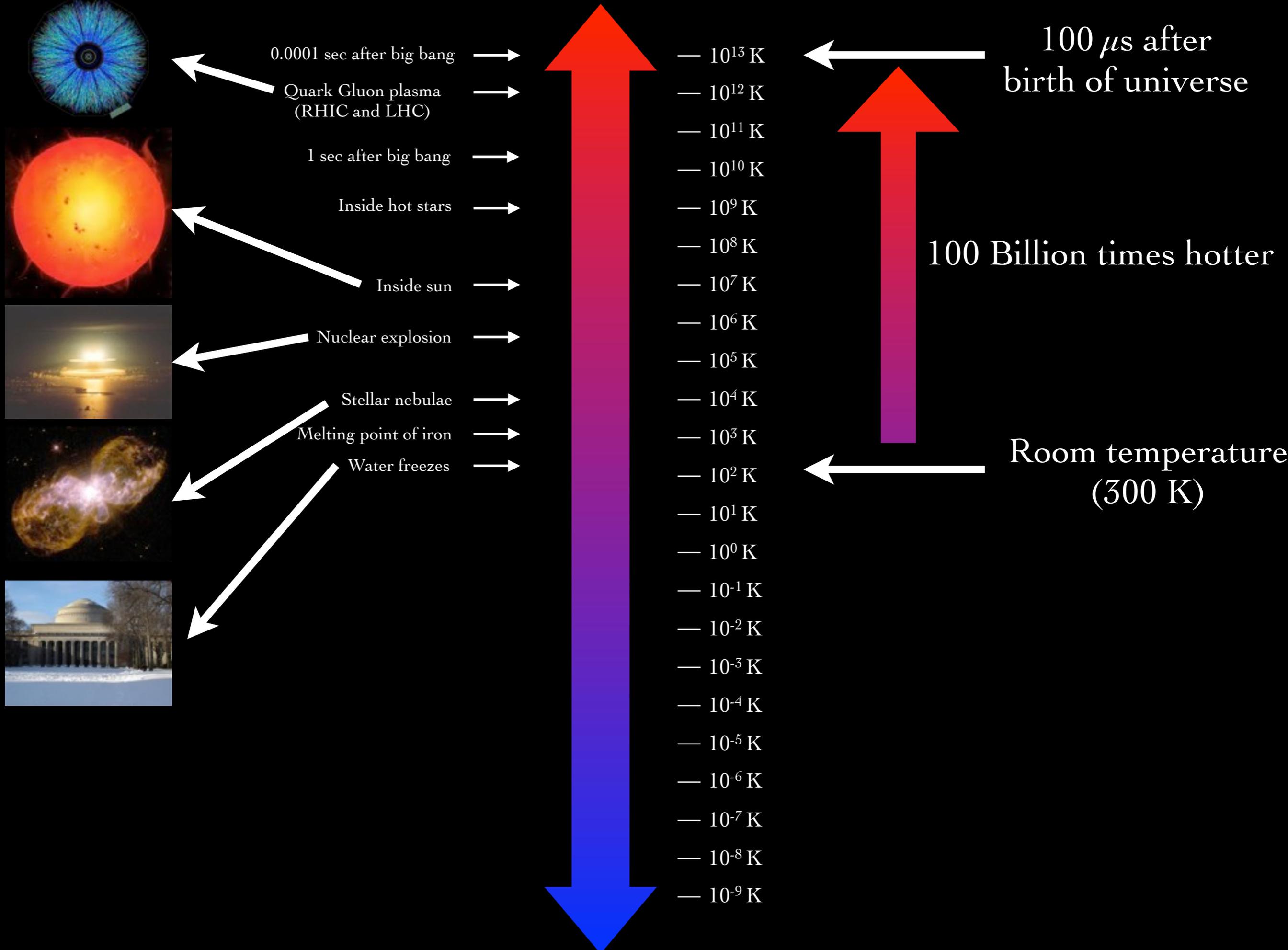


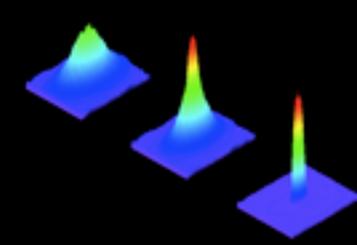
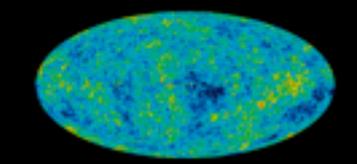
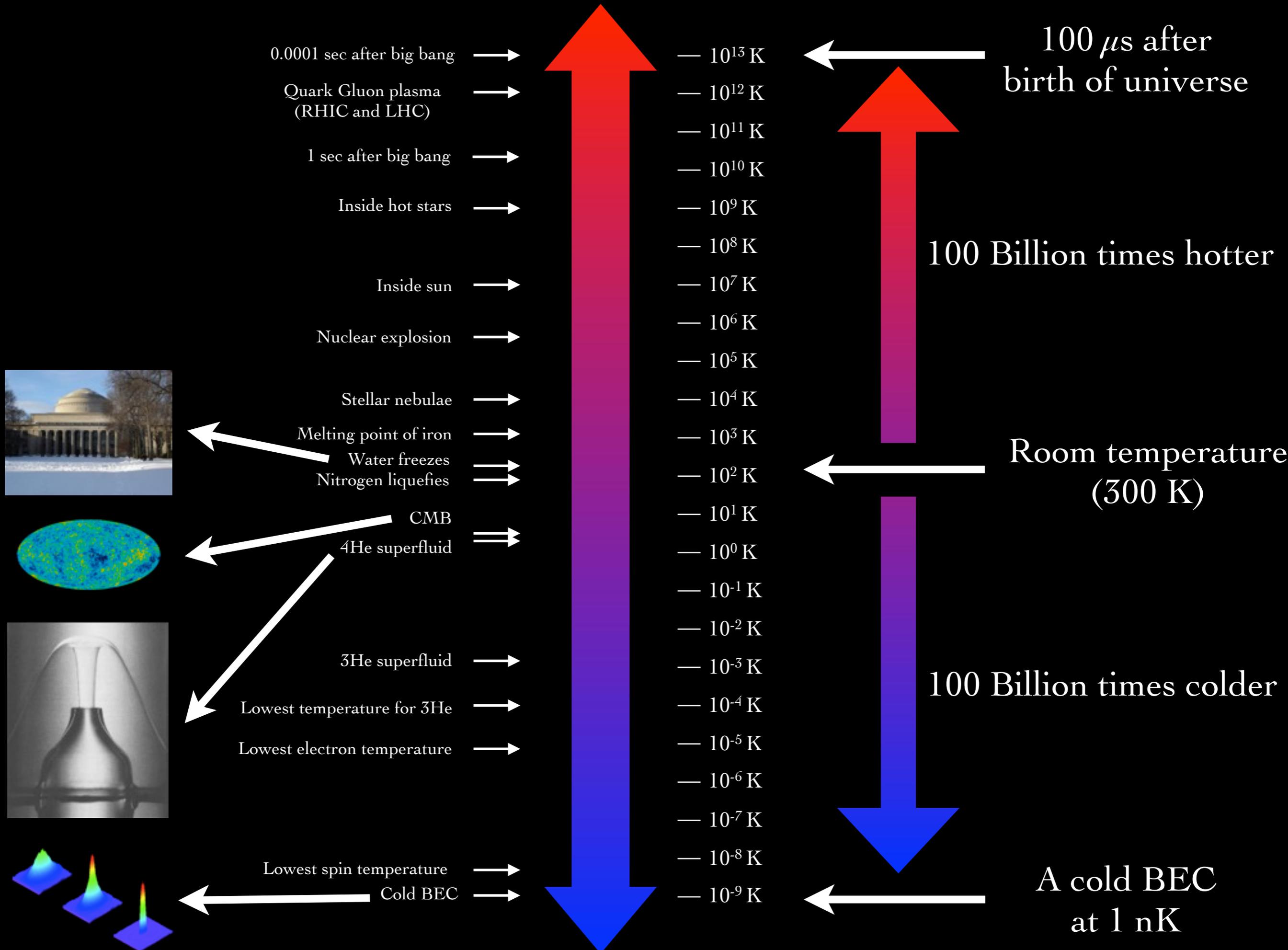
Interference between two BEC's
each with 10^7 atoms



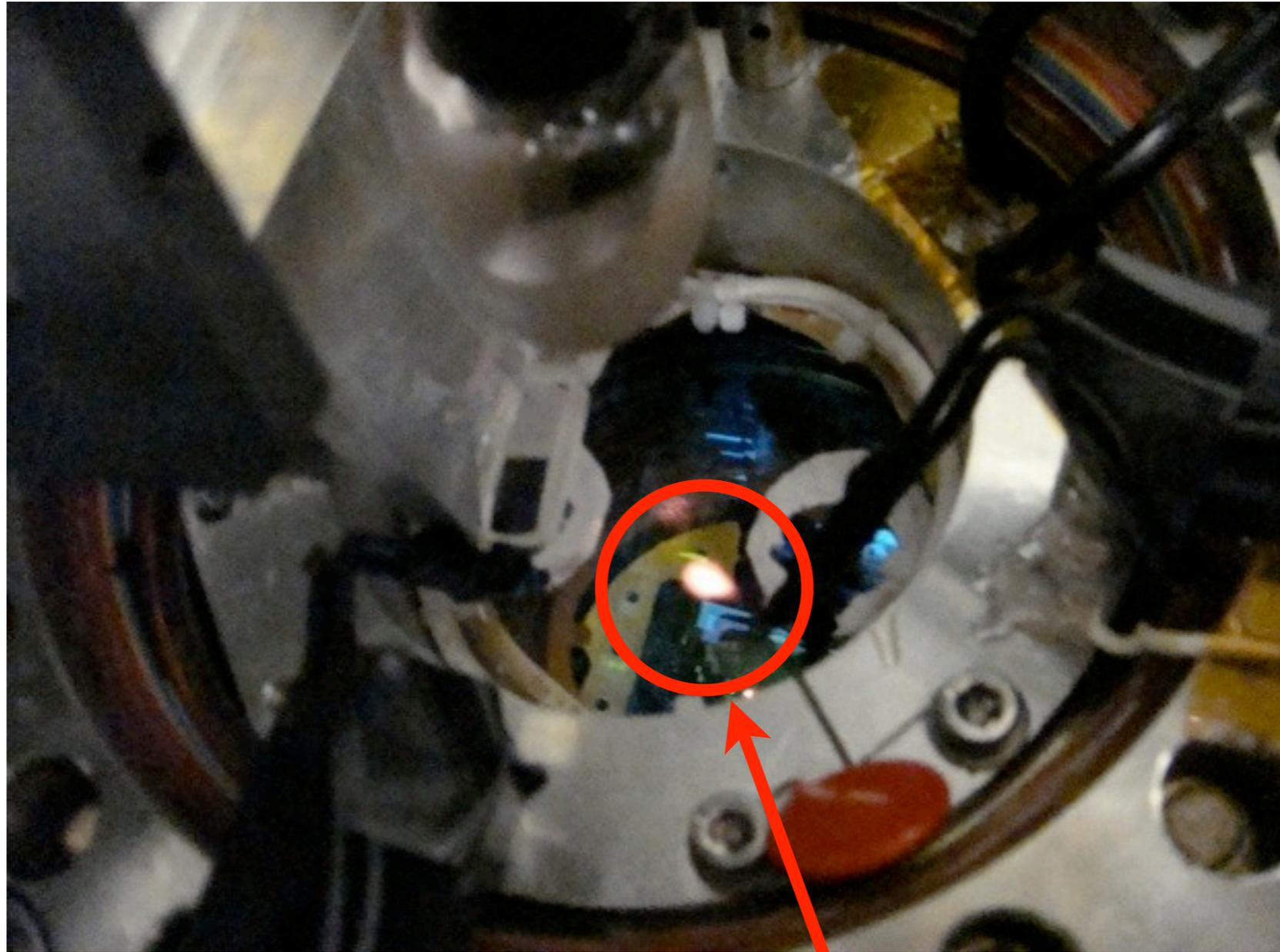
MIT: Andrews et al. Science (1997)

Mass
 4×10^{-19} kg





Anatomy of an experiment: laser cooling



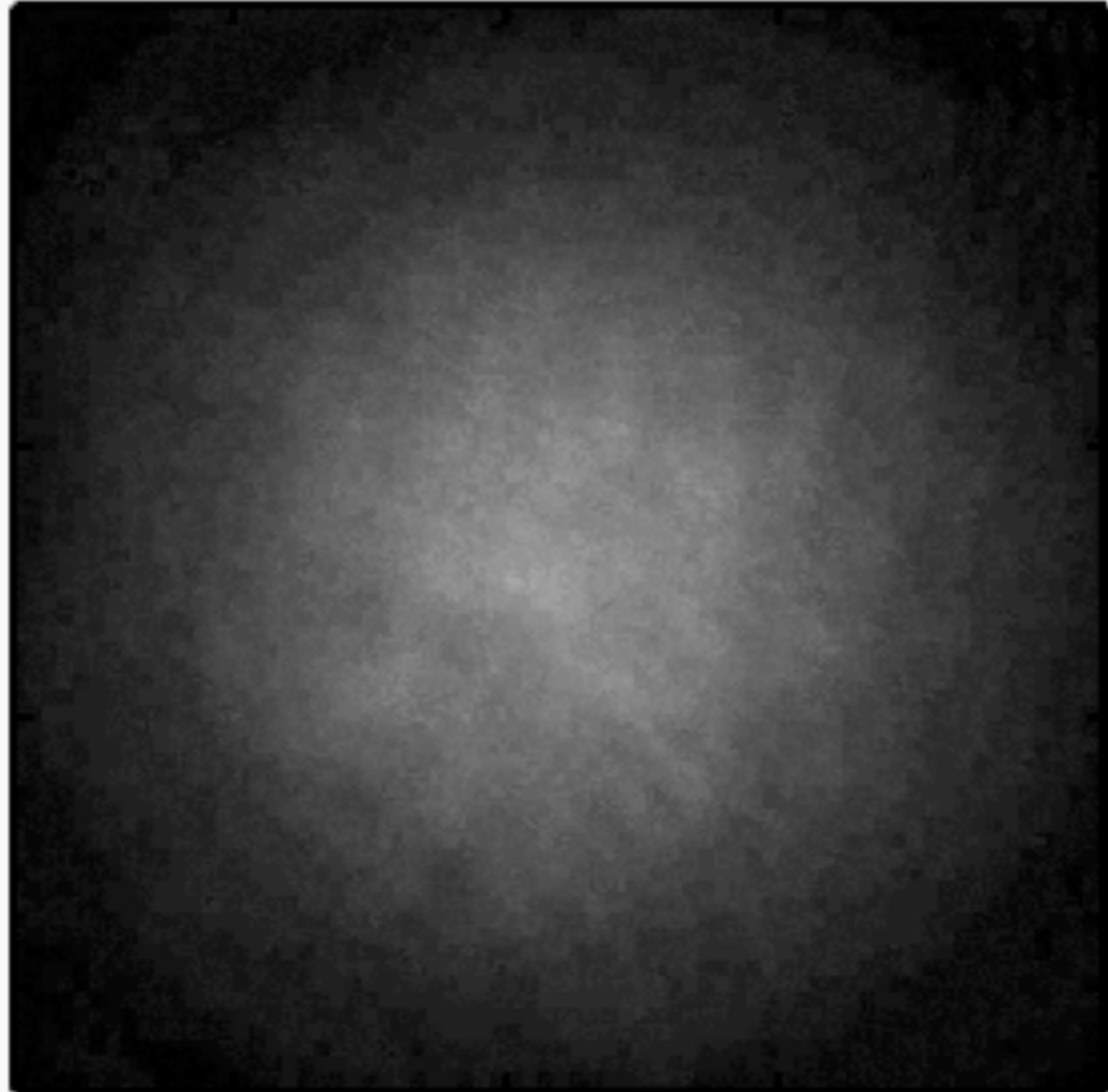
~3 billion laser cooled atoms
 $T \sim 500 \mu\text{K}$

Anatomy of an experiment: evaporation

Laser cooling



Evaporation



Anatomy of an experiment: detection

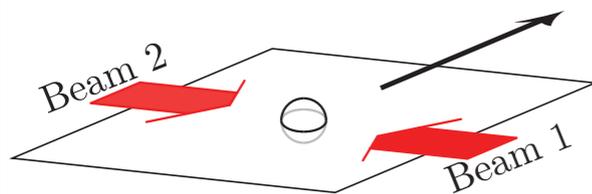
Laser cooling



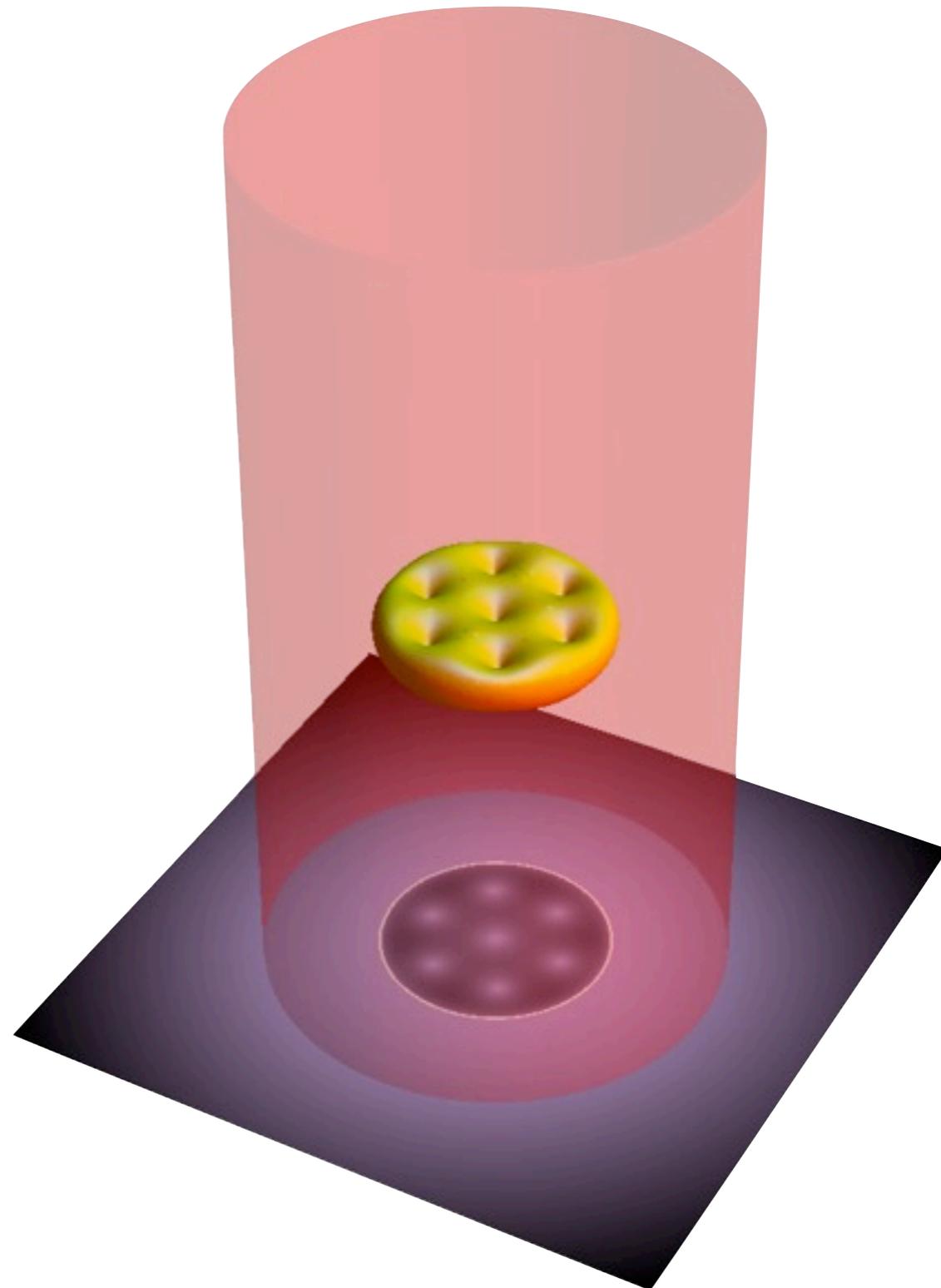
Evaporation



Experiment

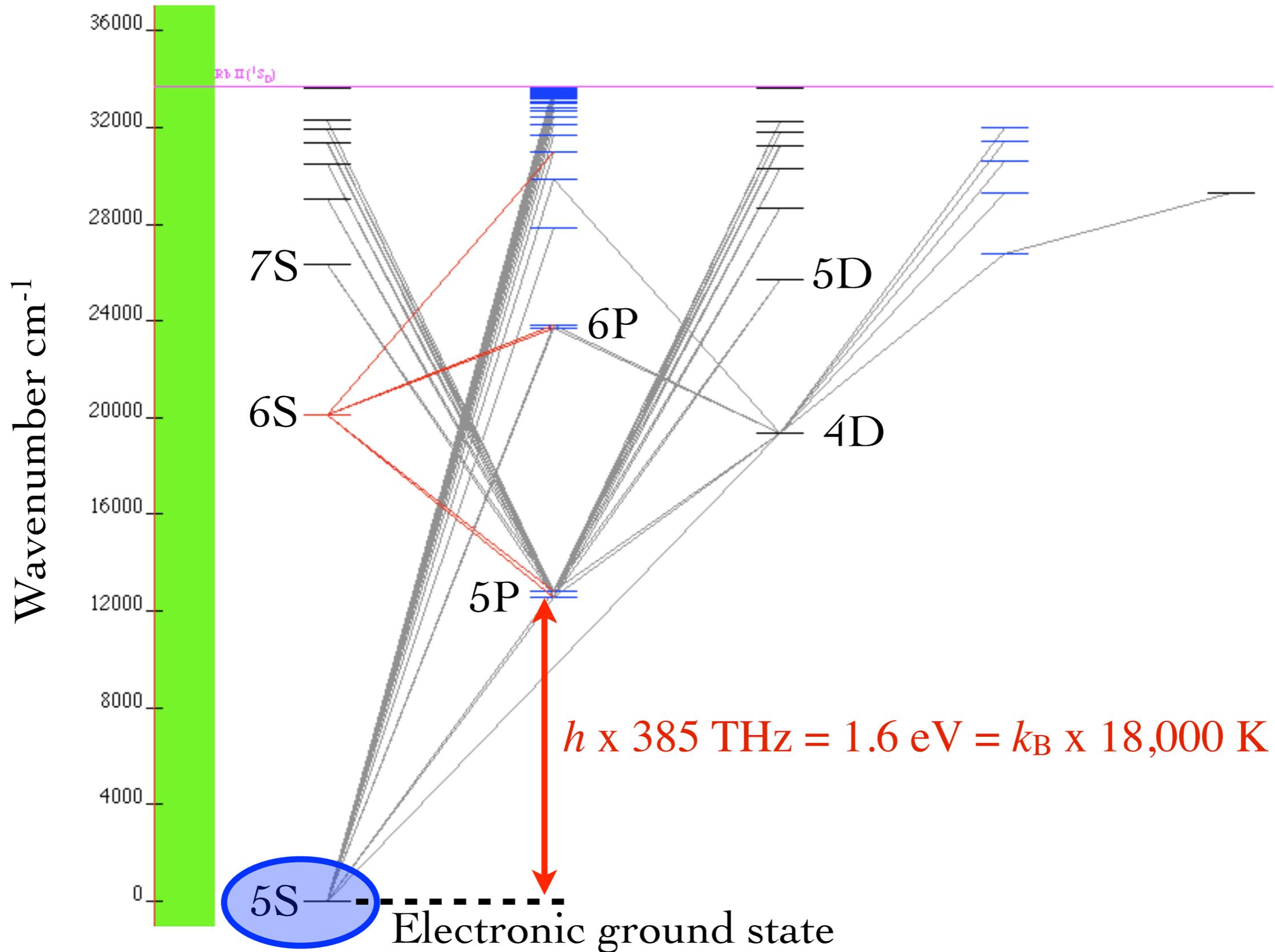


Absorption imaging

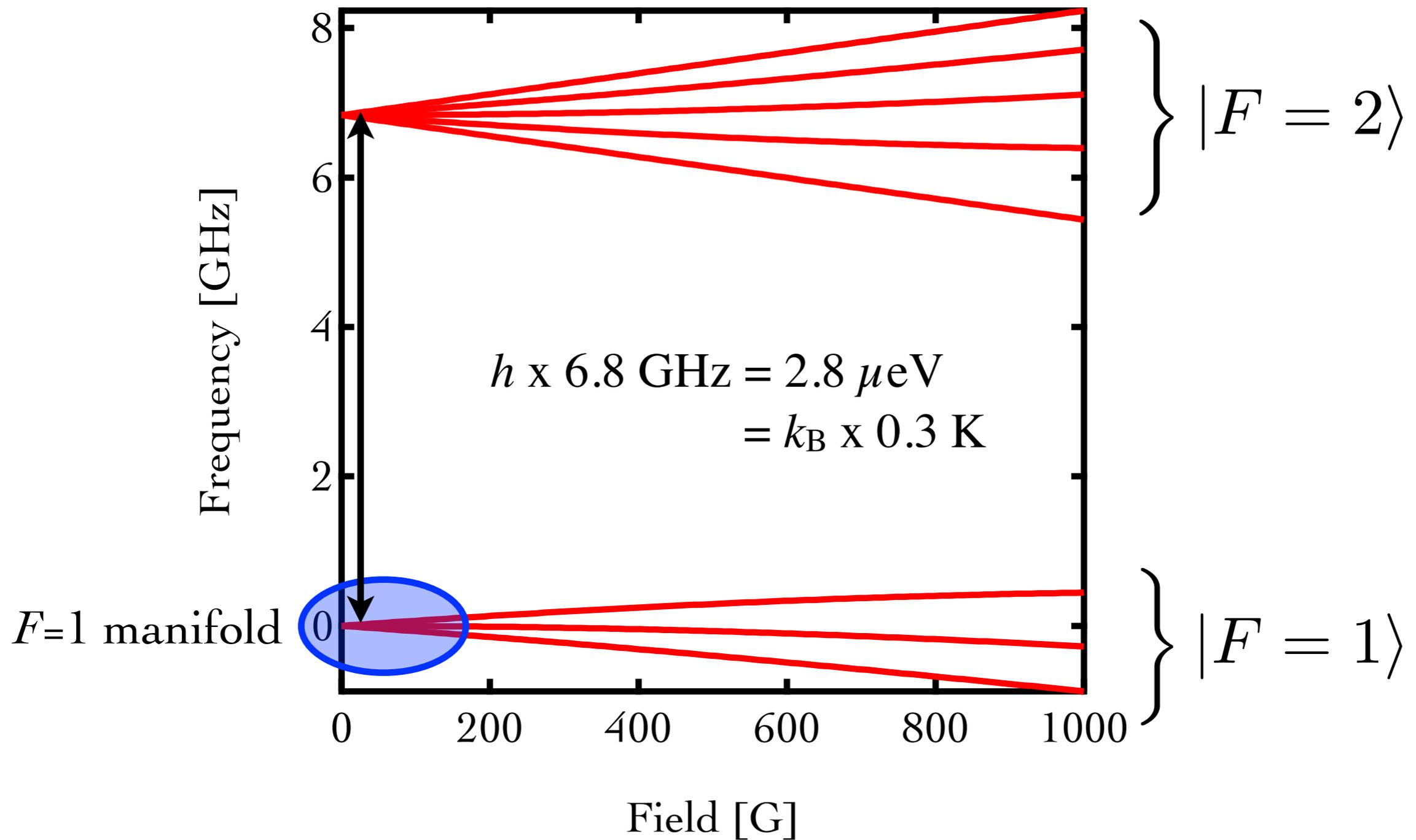


Rubidium 87: “The GaAs of atoms”

An atom is perhaps *the* quintessential quantum system



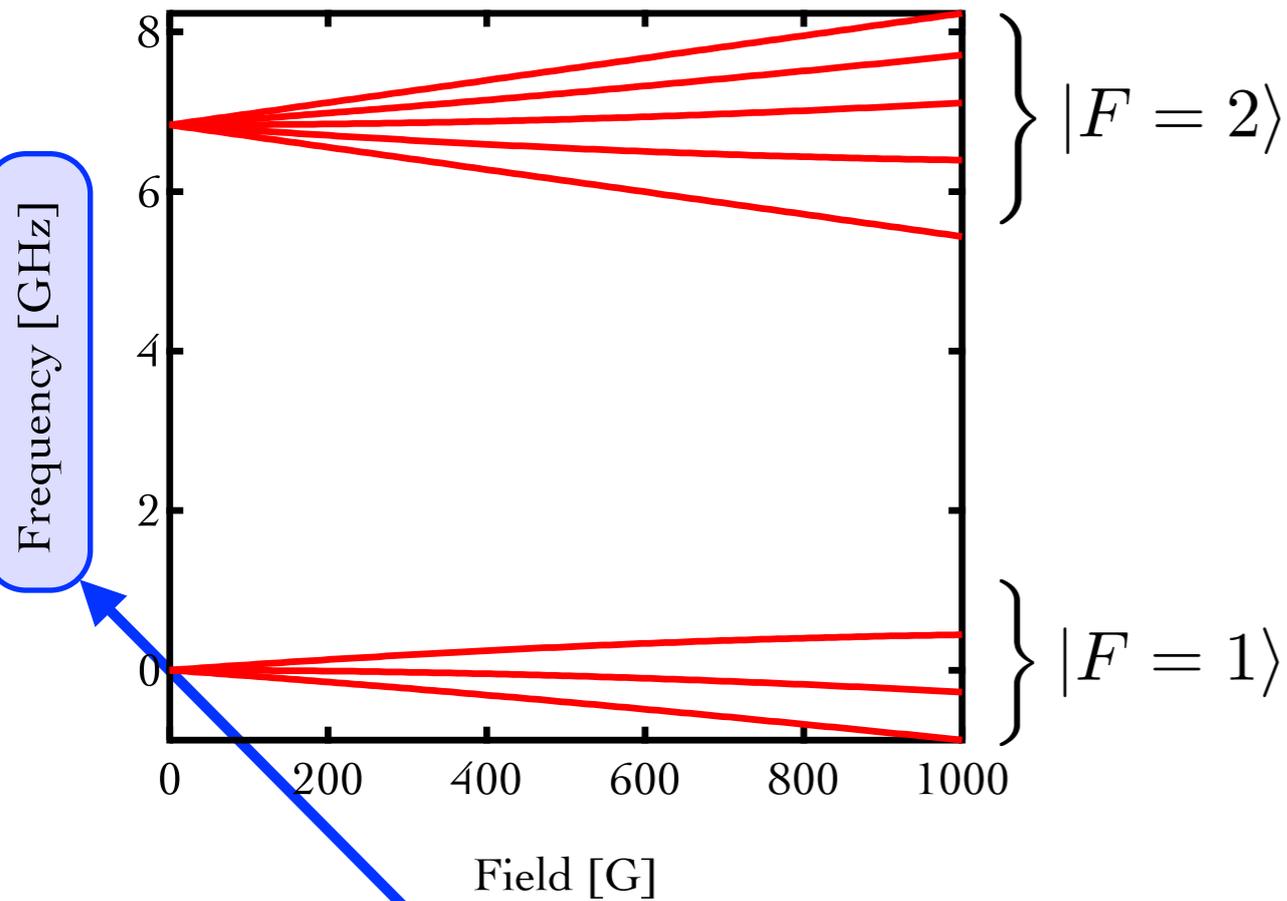
Rubidium 87: $5S_{1/2}$ ground state



What do magnetic fields do?

To spins

Zeeman effect, example of ^{87}Rb



Energy in dimensions of Hz

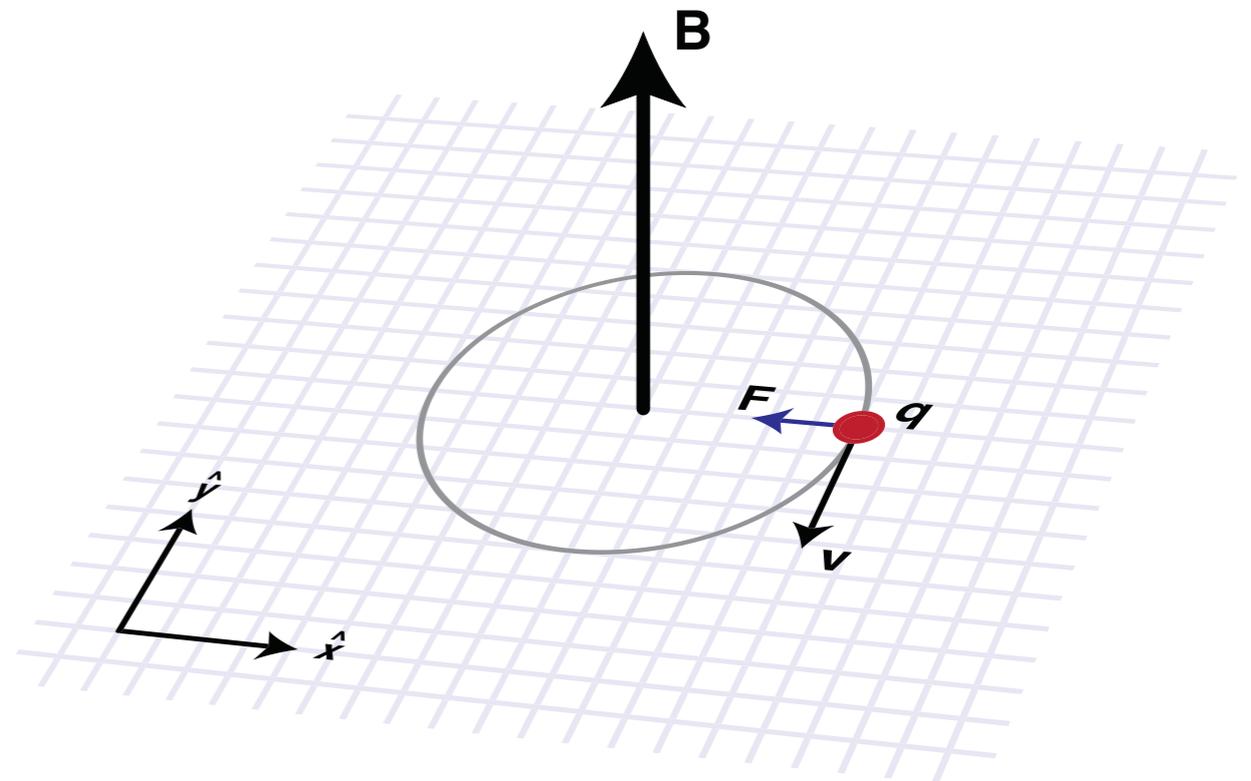
To charges

Lorentz force

Freshman mechanics

Mechanical variables and forces

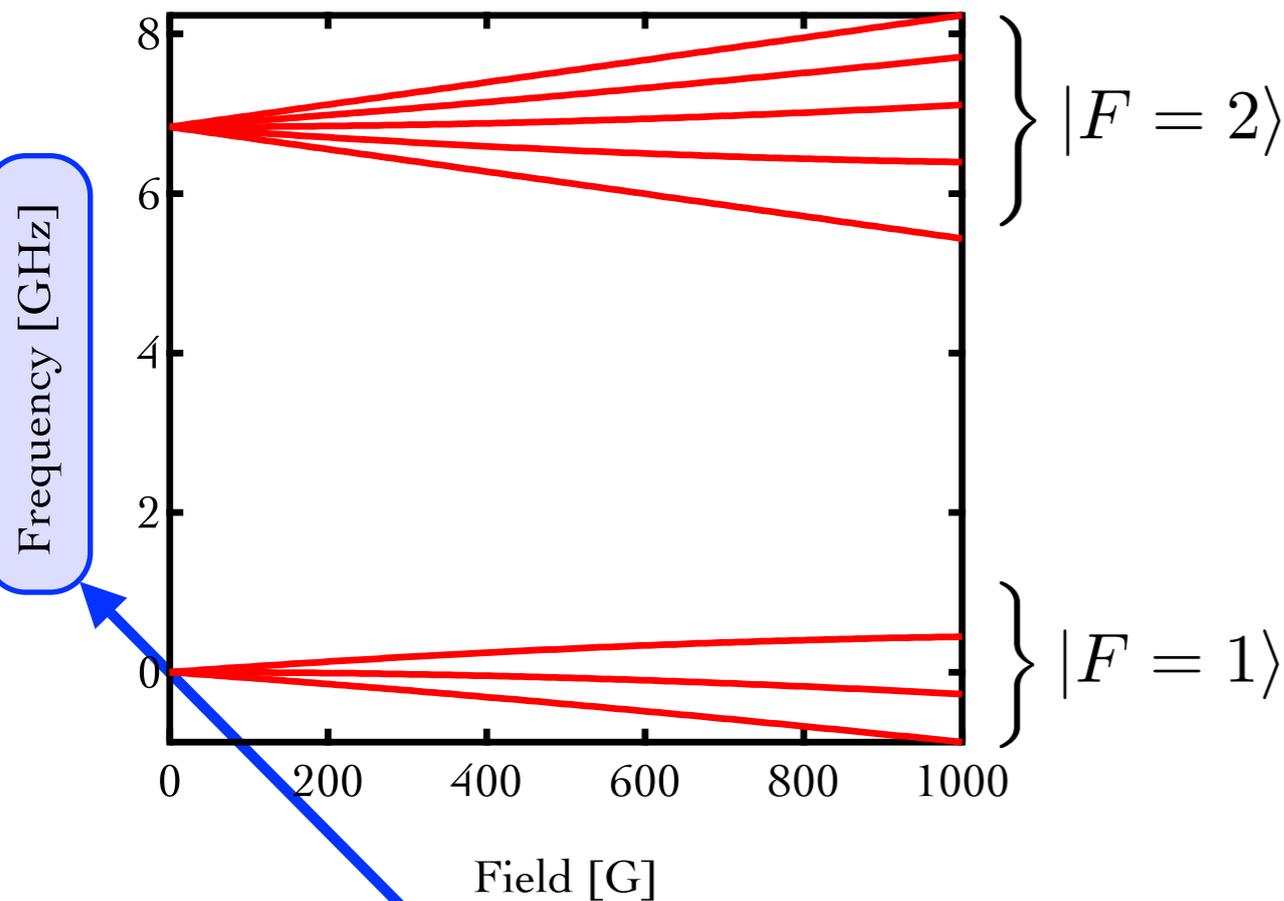
$$\mathbf{F} = q\mathbf{v} \times \mathbf{B}$$



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Lorentz force

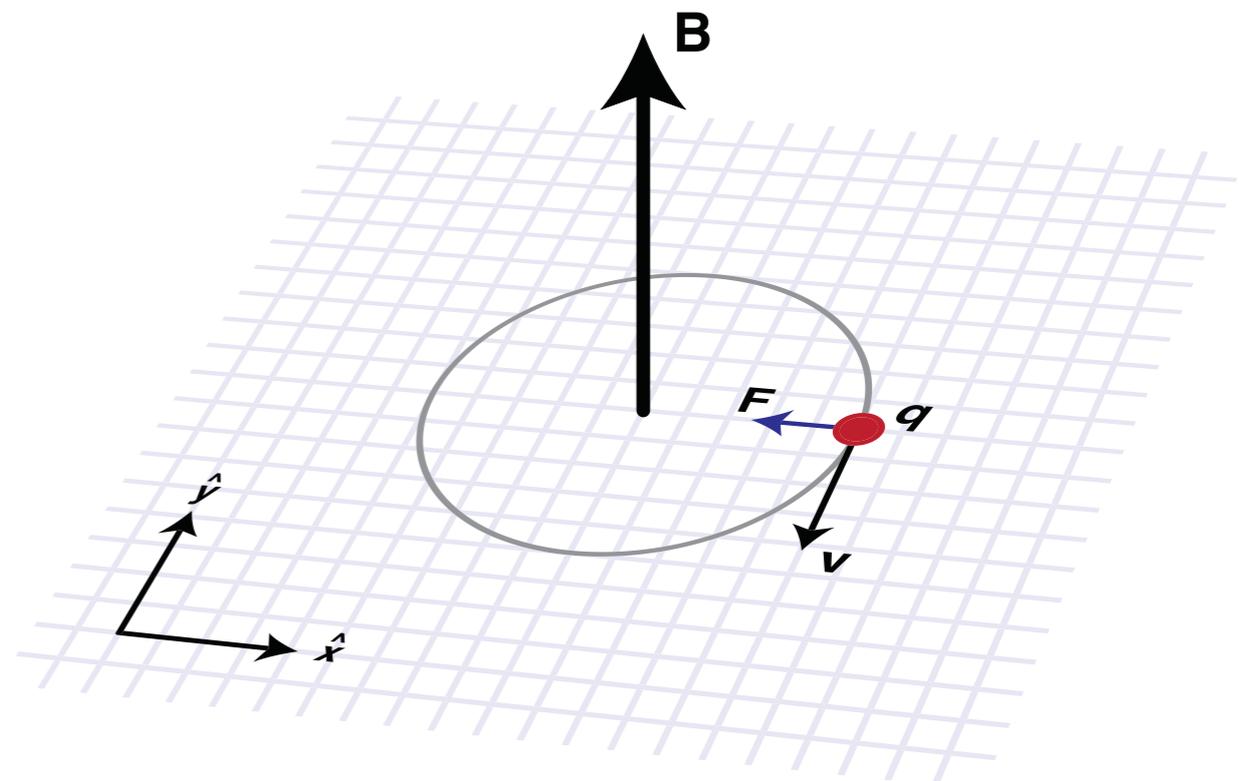
Junior mechanics

Canonical variables and vector potential

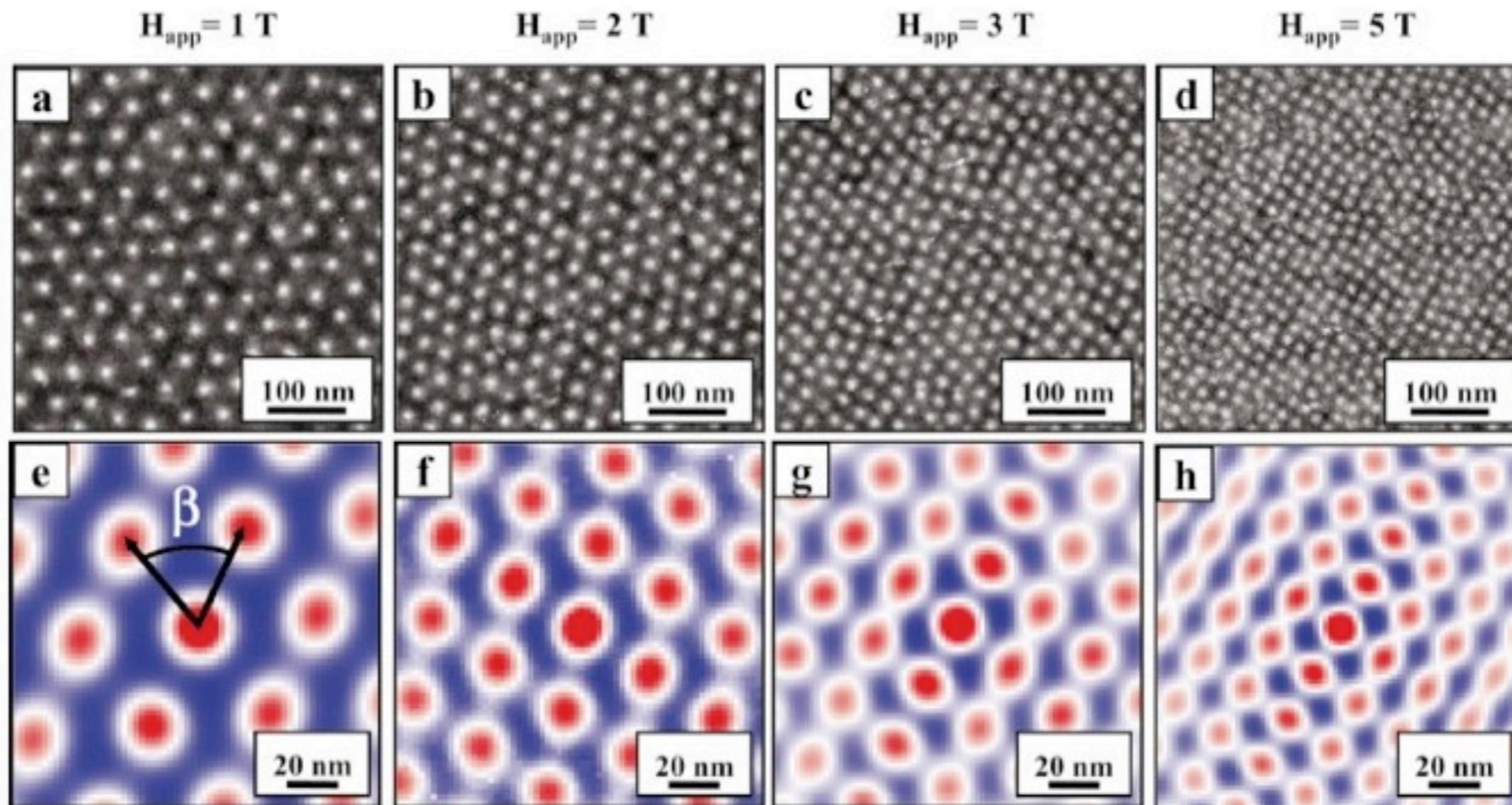
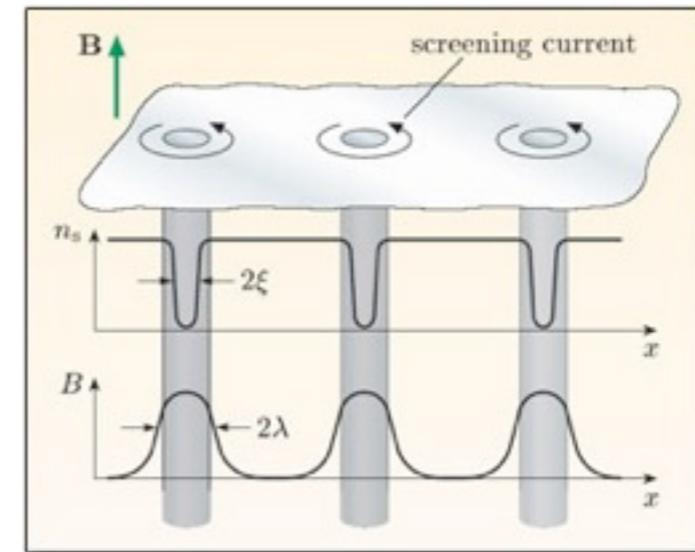
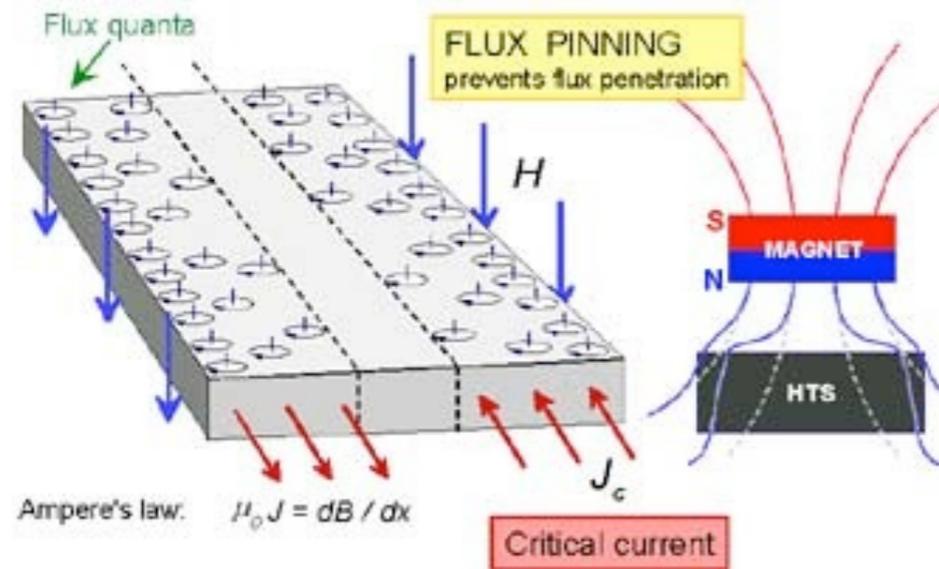
$$\mathbf{B} = \nabla \times \mathbf{A} \quad \text{e.g., } \mathbf{A} = \frac{B}{2}(x\hat{y} - y\hat{x})$$

$$H = \frac{(\mathbf{p} - q\mathbf{A})^2}{2m}$$

$$\dot{p}_j = -\frac{\partial}{\partial x_j} H, \quad \dot{x}_j = \frac{\partial}{\partial p_j} H$$



Type-2 superconductor in a B field: vortices



References

H. Hess *et al*, PRL (1989); C. E. Sosolik *et al*, PRB (2003)

Vortices

Angular momentum is quantized in units of $\hbar = 1.05 \times 10^{-34}$ J s.
Small!

This disposal unit has $\sim 10^{33}$ quanta (10^{10} per water molecule)!



Vortices

About 10^{52} quanta (10^{14} per molecule)



Vortex: big

About 10^{95} quanta (10^{27} per hydrogen)



Vortex: little

10^4 quanta (1 per atom)



Simple minded example: rotation

Rotation

$$H_r = \frac{\hbar^2}{2m} \left[\left(\hat{k}'_x - \frac{m\Omega}{\hbar} y \right)^2 + \left(\hat{k}'_y + \frac{m\Omega}{\hbar} x \right)^2 \right] + \frac{m(\omega^2 - \Omega^2)}{2} \mathbf{r}^2$$

JILA

MIT

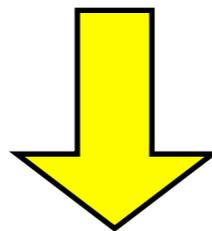
How to simulate magnetic fields

The Hamiltonian in the rotating frame has an effective field.

For high fields fine tuning is required.

ENS, JILA, MIT, ...

$$\hat{R}(\theta = \Omega t) = \exp \left(i\Omega t \hat{L}_z / \hbar \right)$$

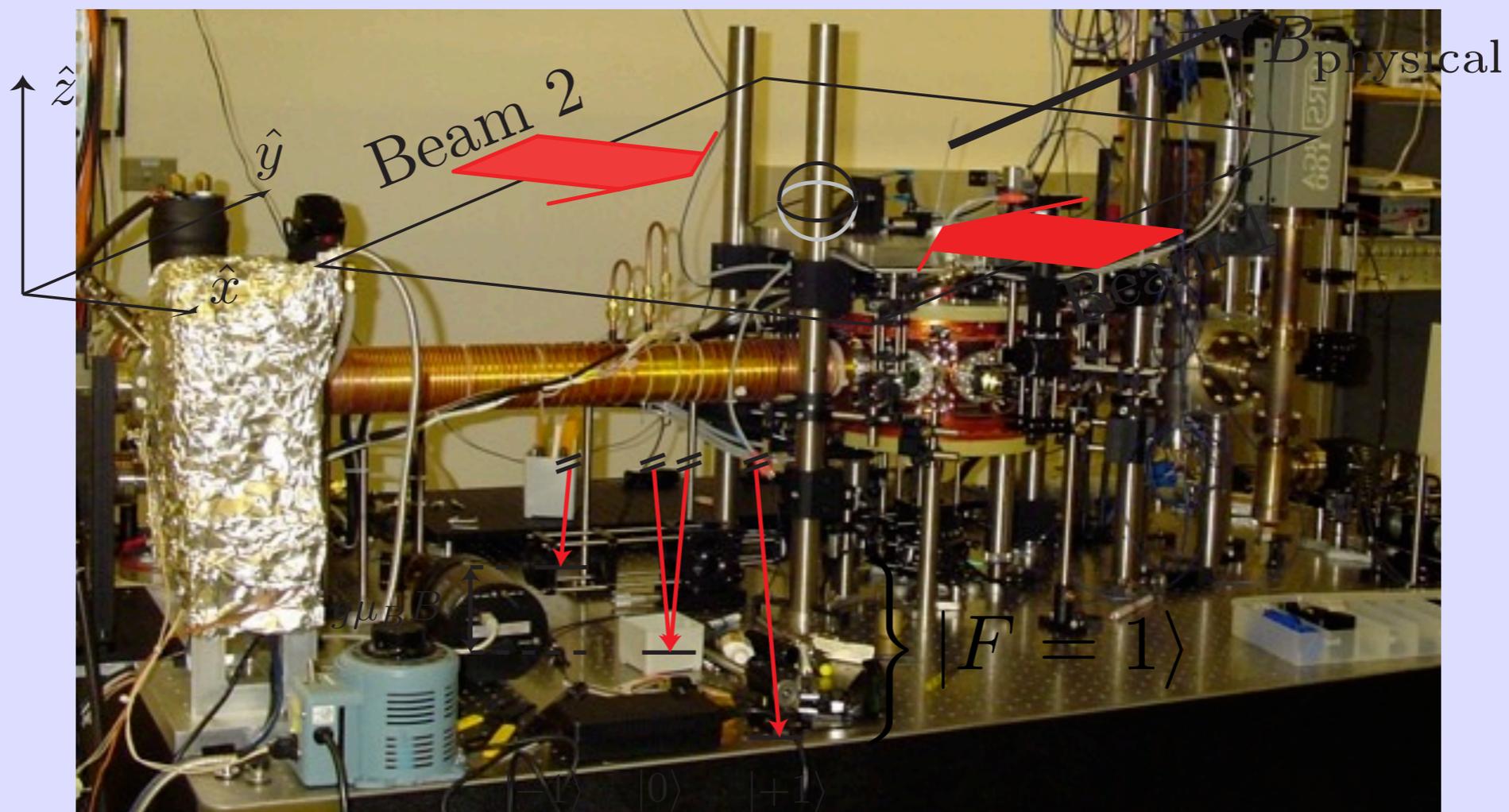


$$H_r = \frac{\hbar^2}{2m} \left[\left(\hat{k}'_y + \frac{m\Omega}{\hbar} x \right)^2 + \left(\hat{k}'_x - \frac{m\Omega}{\hbar} y \right)^2 \right] + \frac{m(\omega^2 - \Omega^2)}{2} (x^2 + y^2)$$

Simplicity from complexity

Raman dressed states

Brief description of implementation and theory

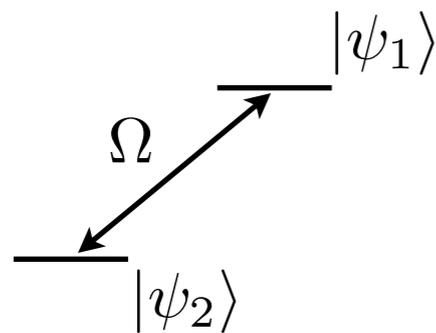


Coupled systems intuition

Schematic

Two levels coupled together

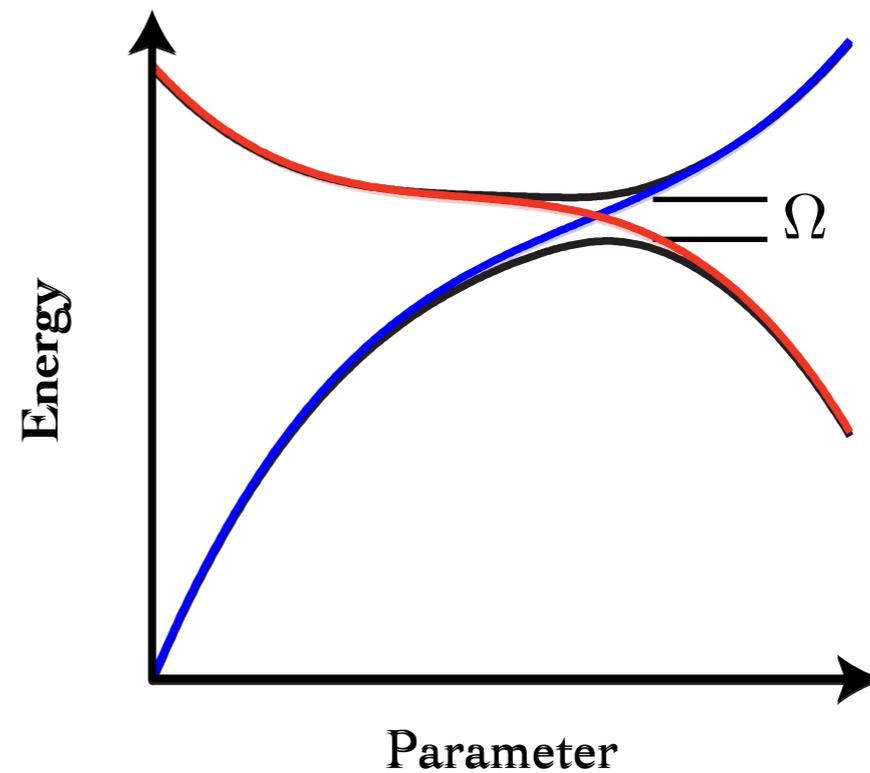
Energy



$$H = \begin{pmatrix} E_1 & \Omega/2 \\ \Omega/2 & E_2 \end{pmatrix}$$

Graphic result

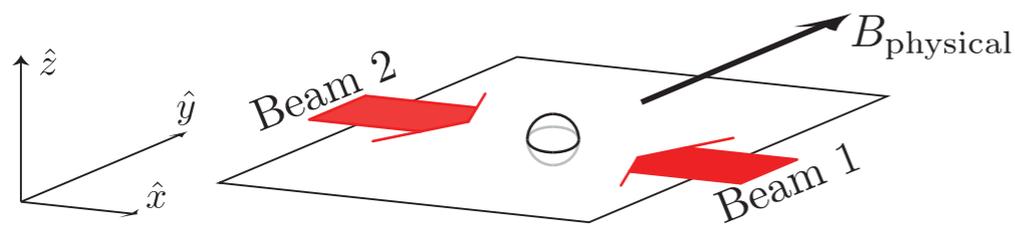
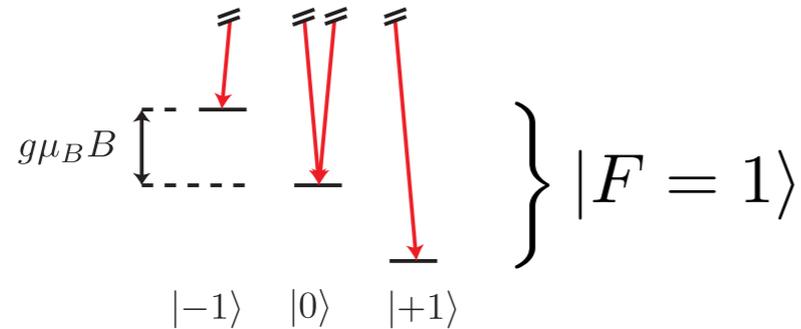
Avoided crossings



Atom light interaction: pictures

Atom light interaction

Given the following geometry and levels

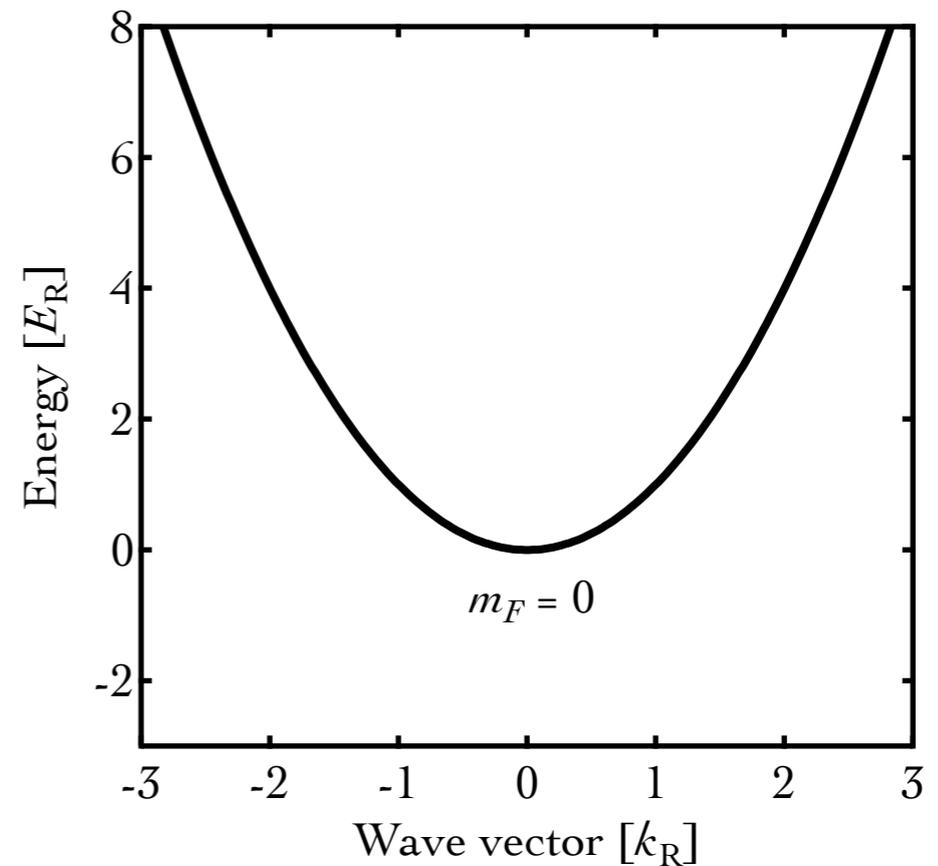


Dimensions

$$k_R = \frac{2\pi}{\lambda}, \quad E_R = \frac{\hbar^2 k_R^2}{2m}$$

$$E_R \approx h \times 3 \text{ kHz} = k_B \times 140 \text{ nK}$$

Coupled States



References

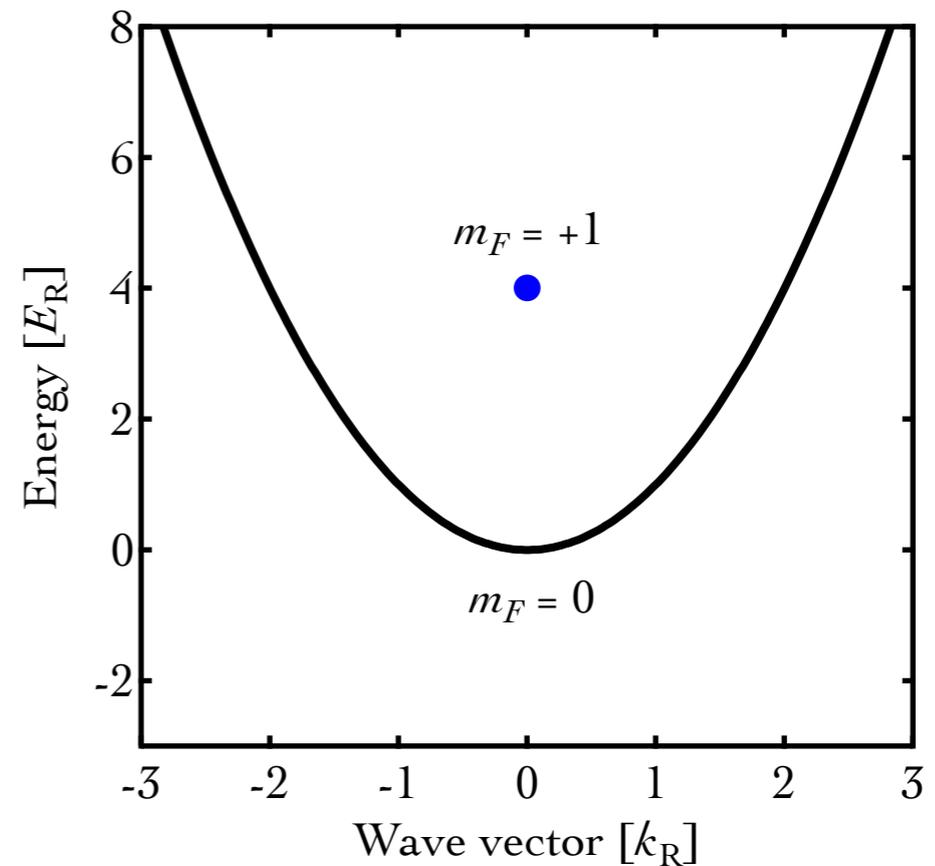
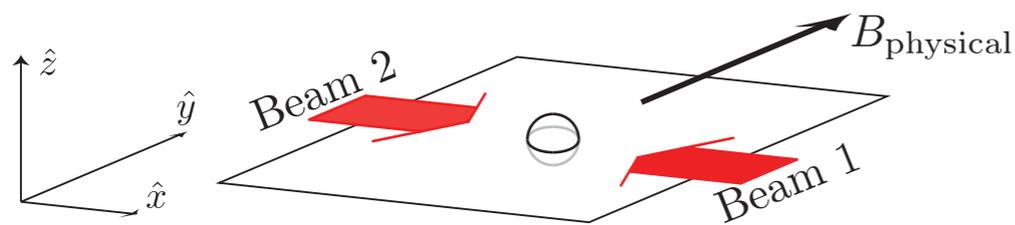
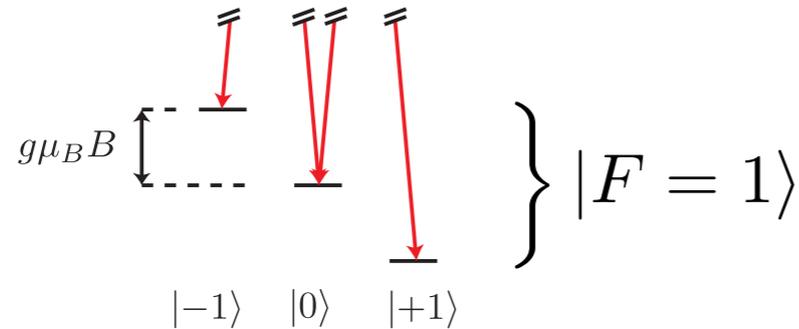
- [1] Juzeliūnas, et al., PRA 025602 **73** (2006), + earlier pubs
- [2] S.-L. Zhu, et al., PRL 240401 **97** (2006)
- [3] Günter et al, PRA **79** 011604 (2009)
- [4] IBS, PRA 063613 **79** (2009)

Atom light interaction: pictures

Atom light interaction

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Coupled States



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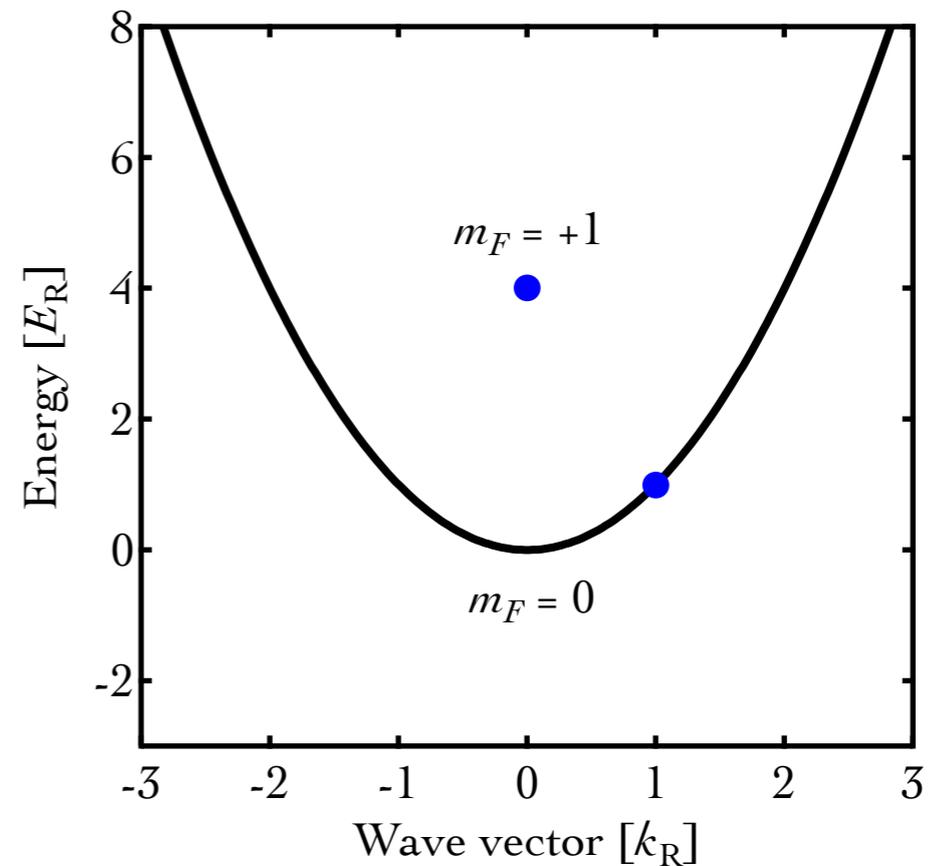
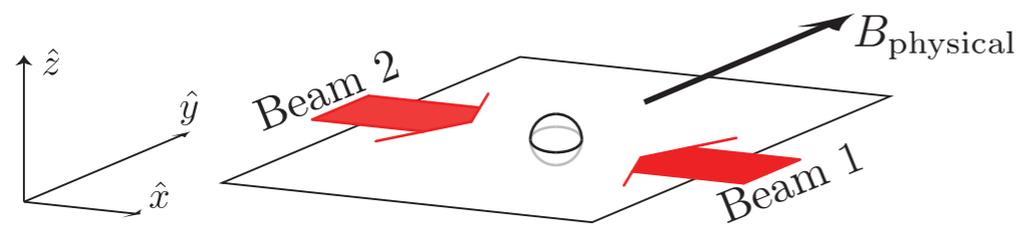
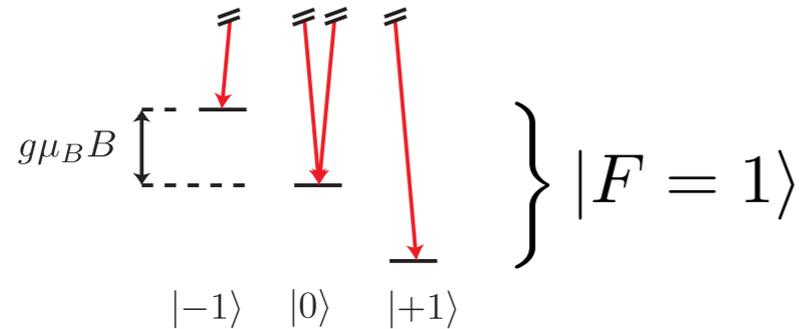
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Atom light interaction: pictures

Atom light interaction

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Coupled States



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References

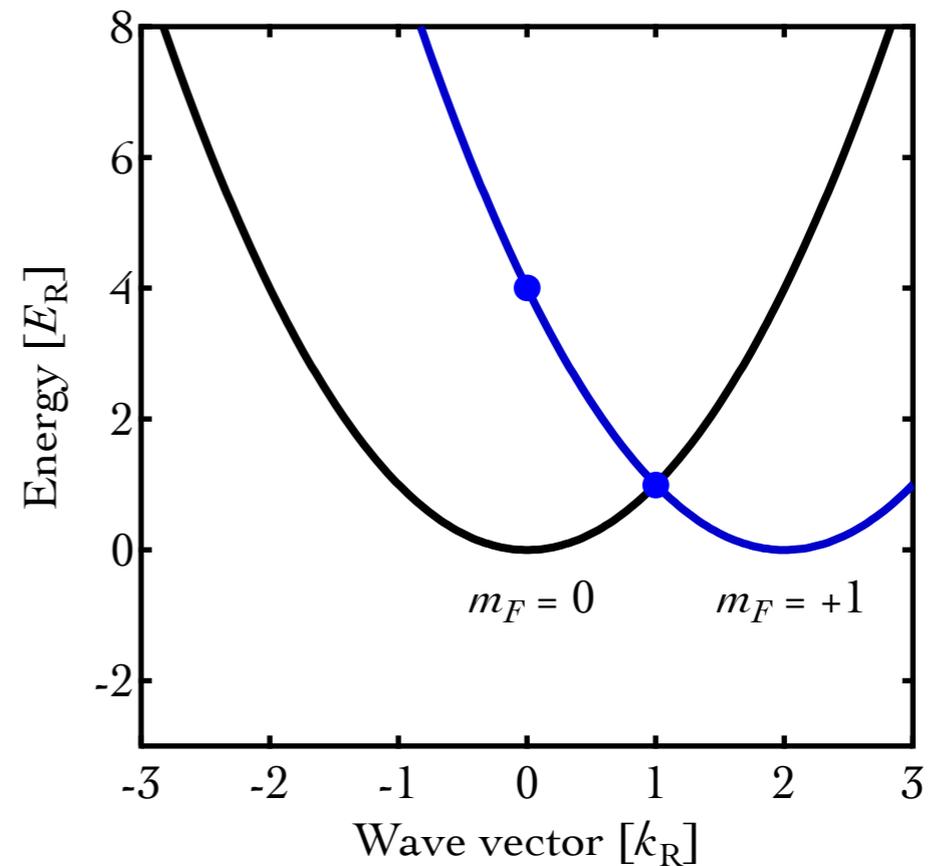
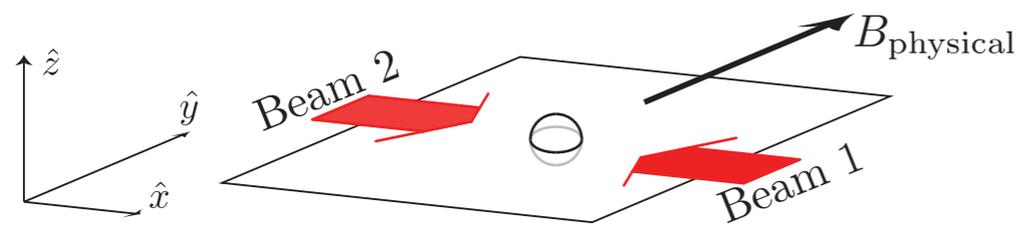
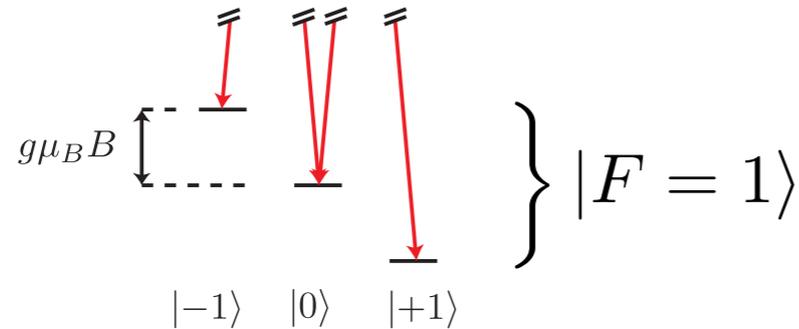
- [1] Juzeliūnas, et al., PRA 025602 **73** (2006), + earlier pubs
- [2] S.-L. Zhu, et al., PRL 240401 **97** (2006)
- [3] Günter et al, PRA **79** 011604 (2009)
- [4] IBS, PRA 063613 **79** (2009)

Atom light interaction: pictures

Atom light interaction

Given the following geometry and levels

Coupled States



Dimensions

$$k_R = \frac{2\pi}{\lambda}, \quad E_R = \frac{\hbar^2 k_R^2}{2m}$$

$$E_R \approx h \times 3 \text{ kHz} = k_B \times 140 \text{ nK}$$

References

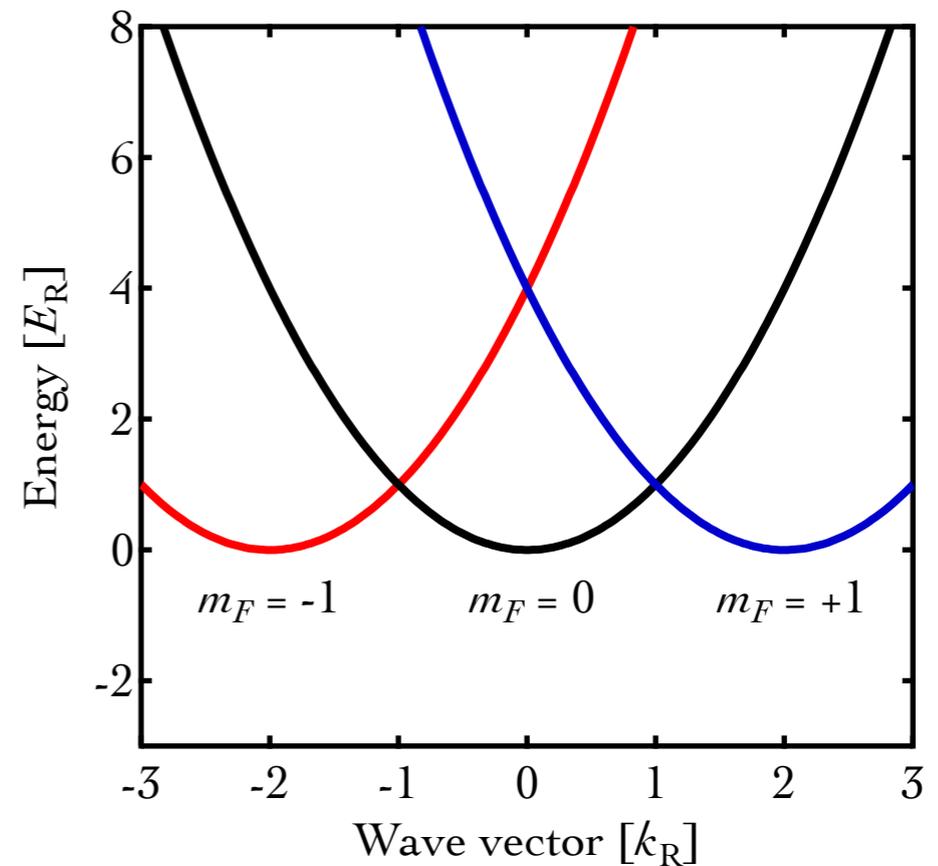
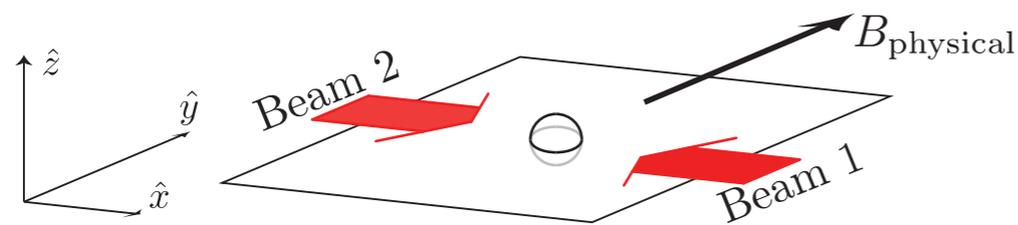
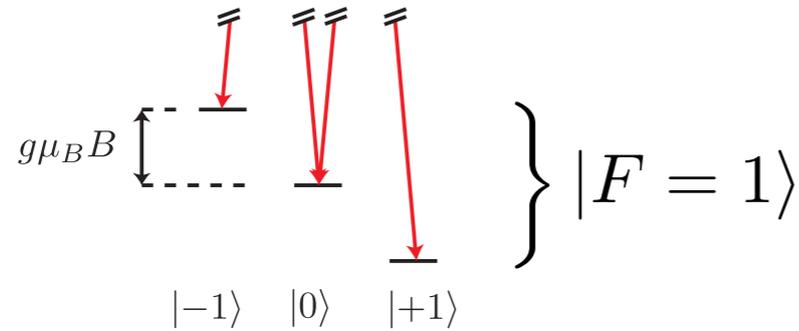
- [1] Juzeliūnas, et al., PRA 025602 **73** (2006), + earlier pubs
- [2] S.-L. Zhu, et al., PRL 240401 **97** (2006)
- [3] Günter et al, PRA **79** 011604 (2009)
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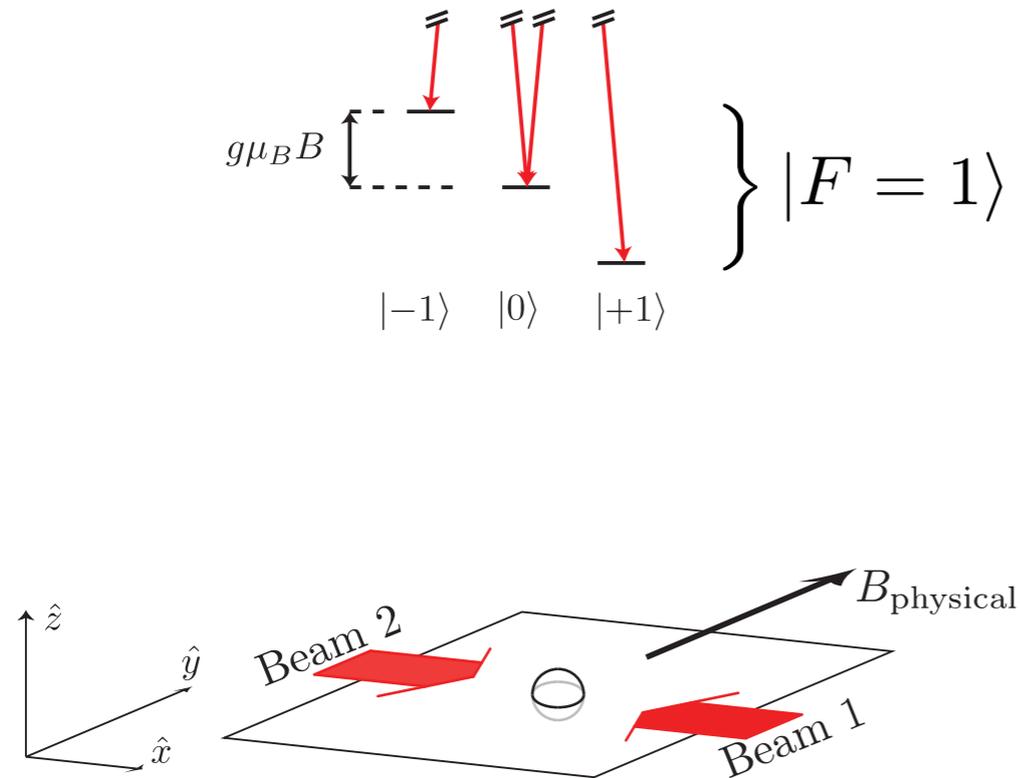
References

- [1] Juzeliūnas, et al., PRA 025602 **73** (2006), + earlier pubs
- [2] S.-L. Zhu, et al., PRL 240401 **97** (2006)
- [3] Günter et al, PRA **79** 011604 (2009)
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Atom light interaction: pictures

Atom light interaction

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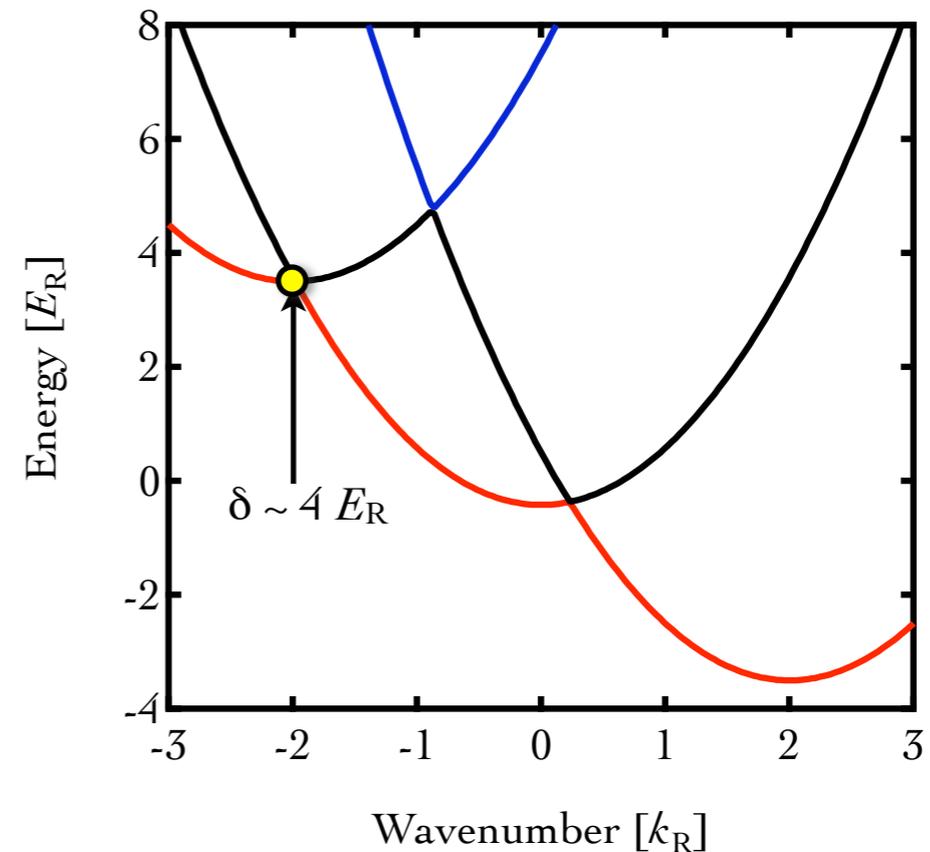
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$$E_R \approx h \times 3 \text{ kHz} = k_B \times 140 \text{ nK}$$

Coupled States

States will be labeled by:
 (1) the “band index” and by
 (2) a quasi-momentum k



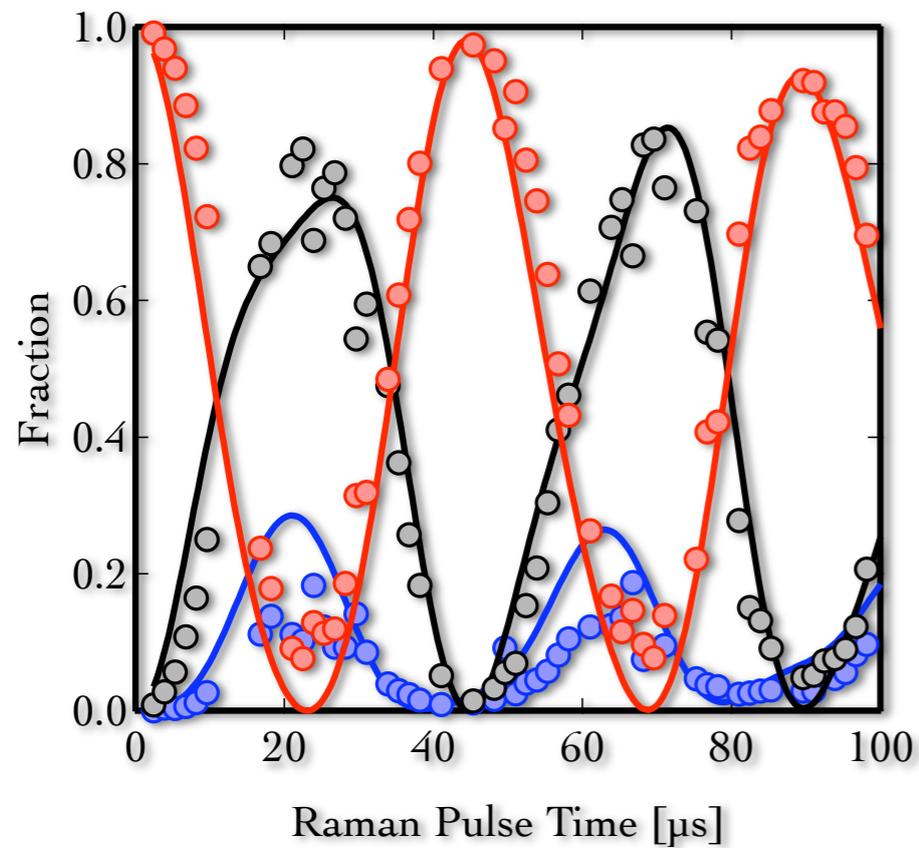
References

- [1] Juzeliūnas, et al., PRA 025602 **73** (2006), + earlier pubs
- [2] S.-L. Zhu, et al., PRL 240401 **97** (2006)
- [3] Günter et al, PRA **79** 011604 (2009)
- [4] IBS, PRA 063613 **79** (2009)

Atom light interaction: pictures

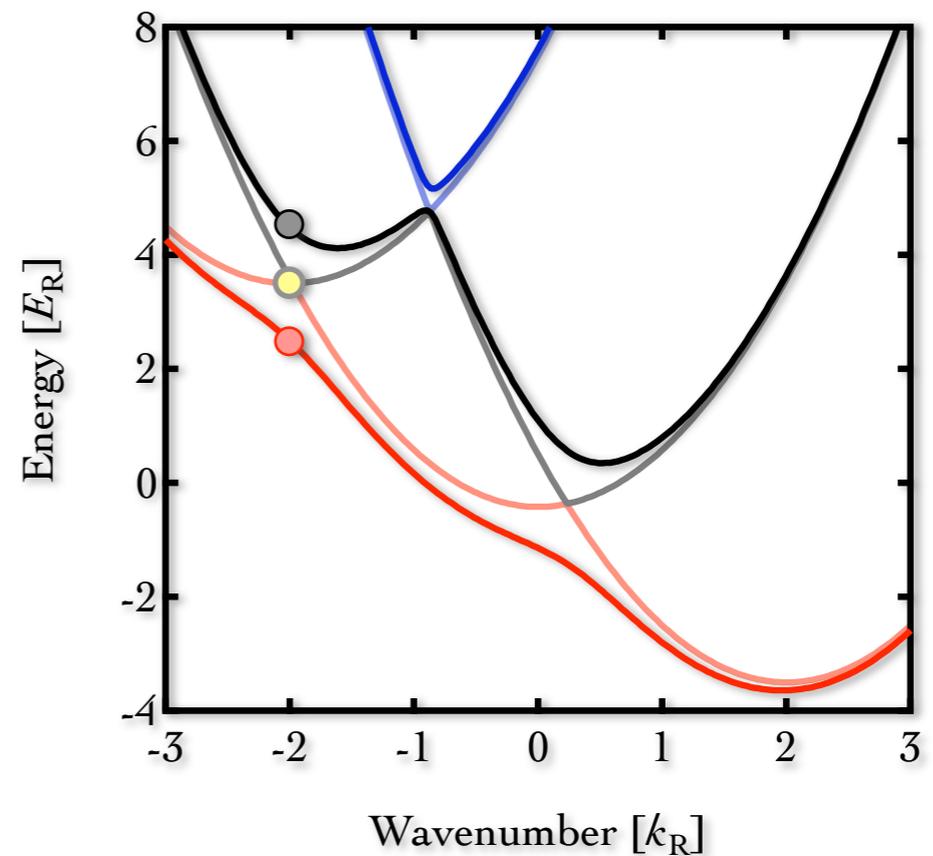
Time evolution

In the sudden limit (Raman-Nath)
Population oscillations yield coupling



Coupled States

States will be labeled by:
(1) the “band index” and by
(2) a quasi-momentum k



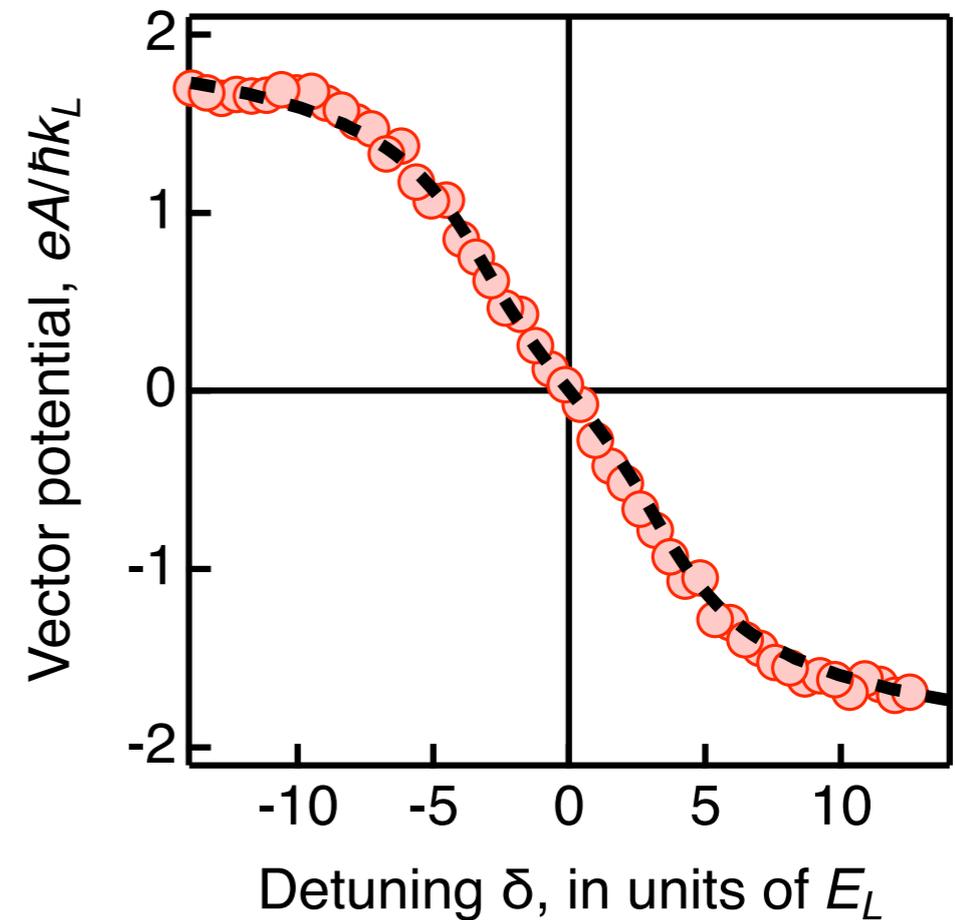
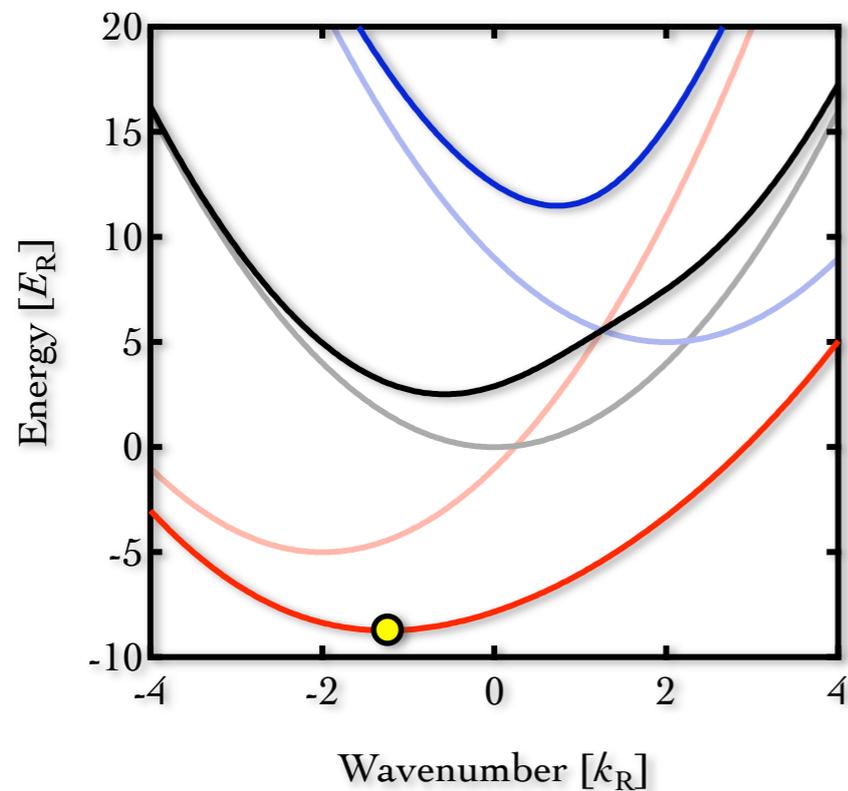
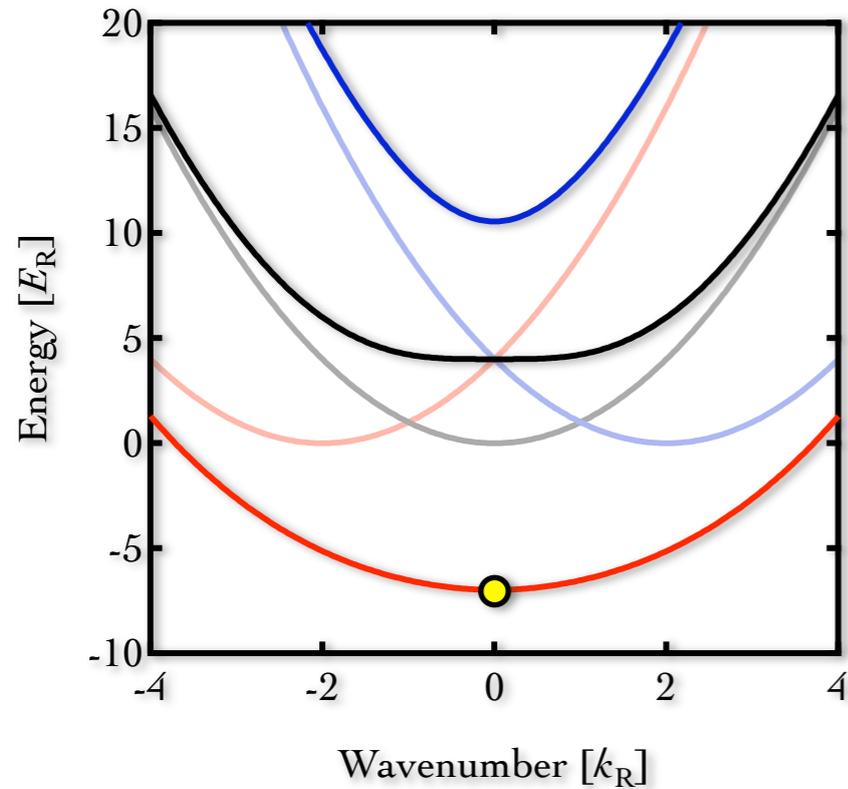
Fundamental intuition

Transfer function

$$\hat{H} = \frac{\hbar^2}{2m} \left\{ \left[k_x - \frac{qA_x(\delta, \Omega)}{\hbar} \right]^2 + k_y^2 \right\} + V(\mathbf{x})$$

where $\delta(x, y, t)$ and $\Omega(x, y, t)$

The detuning and coupling specify the local synthetic vector potential



References

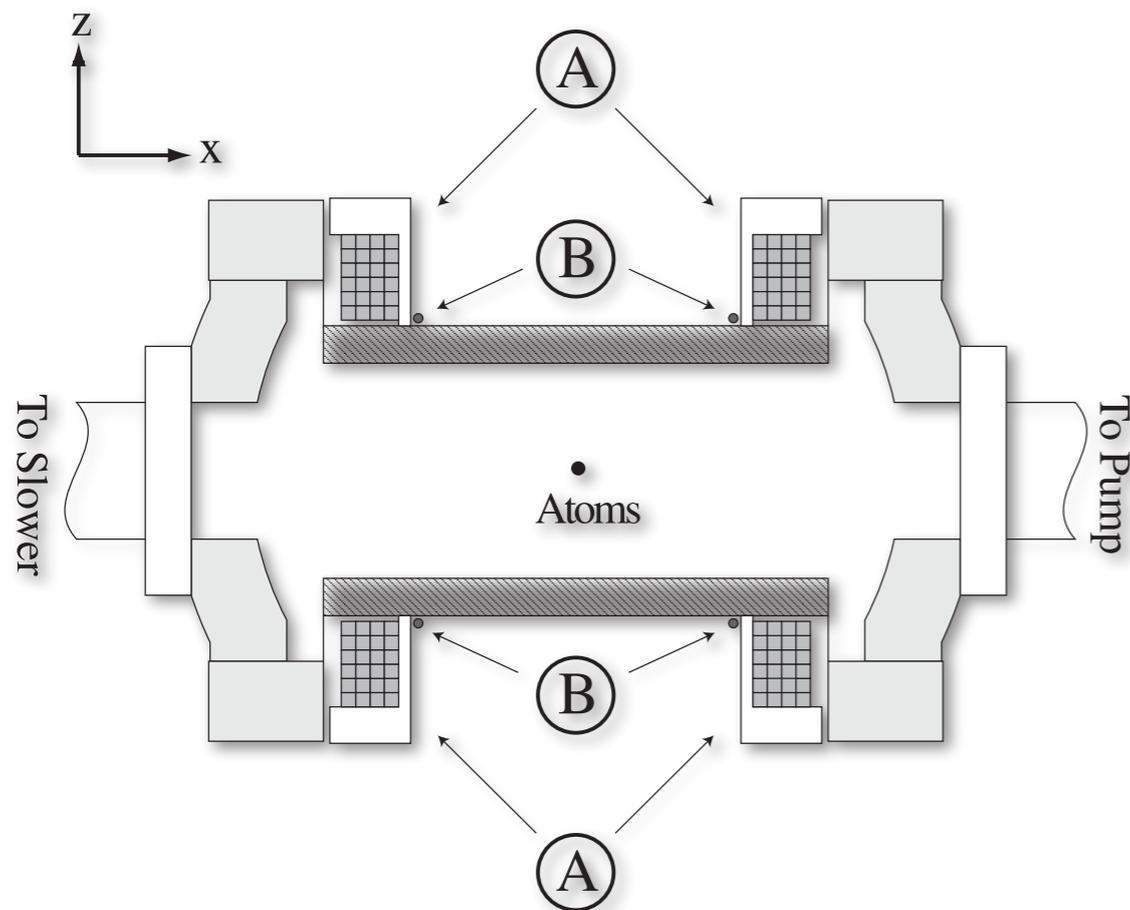
Y.-J. Lin et al, PRL (2009)

A laboratory tunable vector potential

Idea

We can control the *engineered* vector potential in time and space giving *synthetic* \mathbf{E} and \mathbf{B} fields.

Bias and quadrupole \mathbf{B} fields = offset and gradient in detuning.

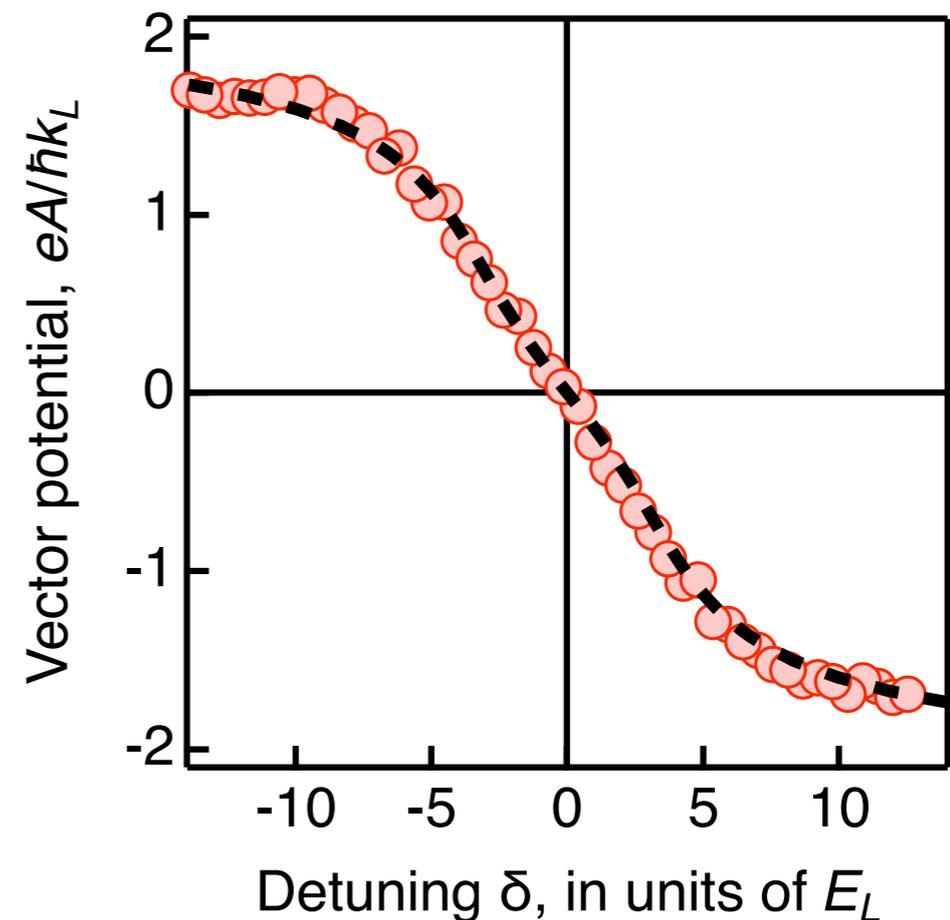


Transfer function

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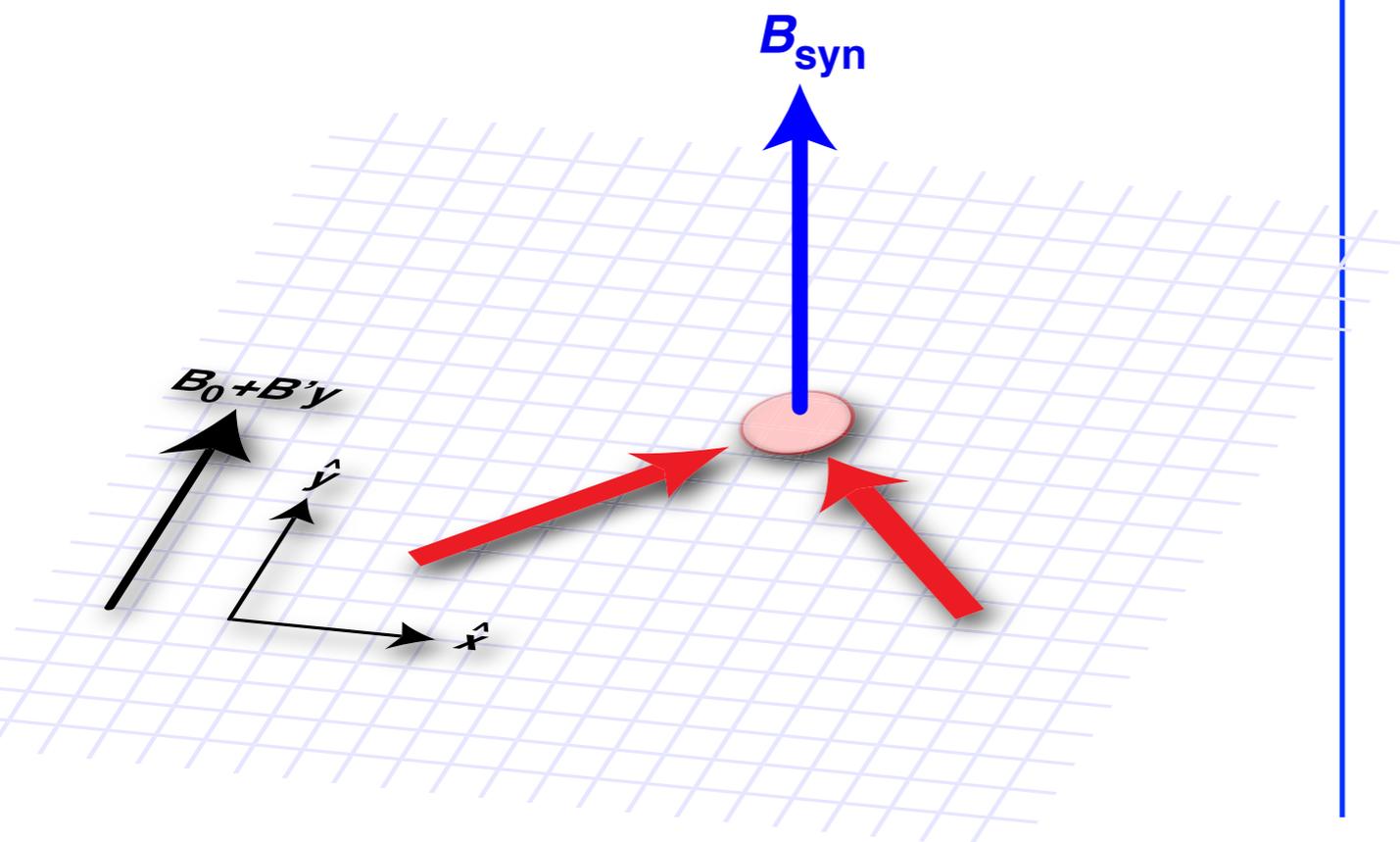
Y.-J. Lin et al, PRL (2009)

Synthetic magnetic field

A non-uniform vector potential

Spatial dependence gives magnetic fields and forces

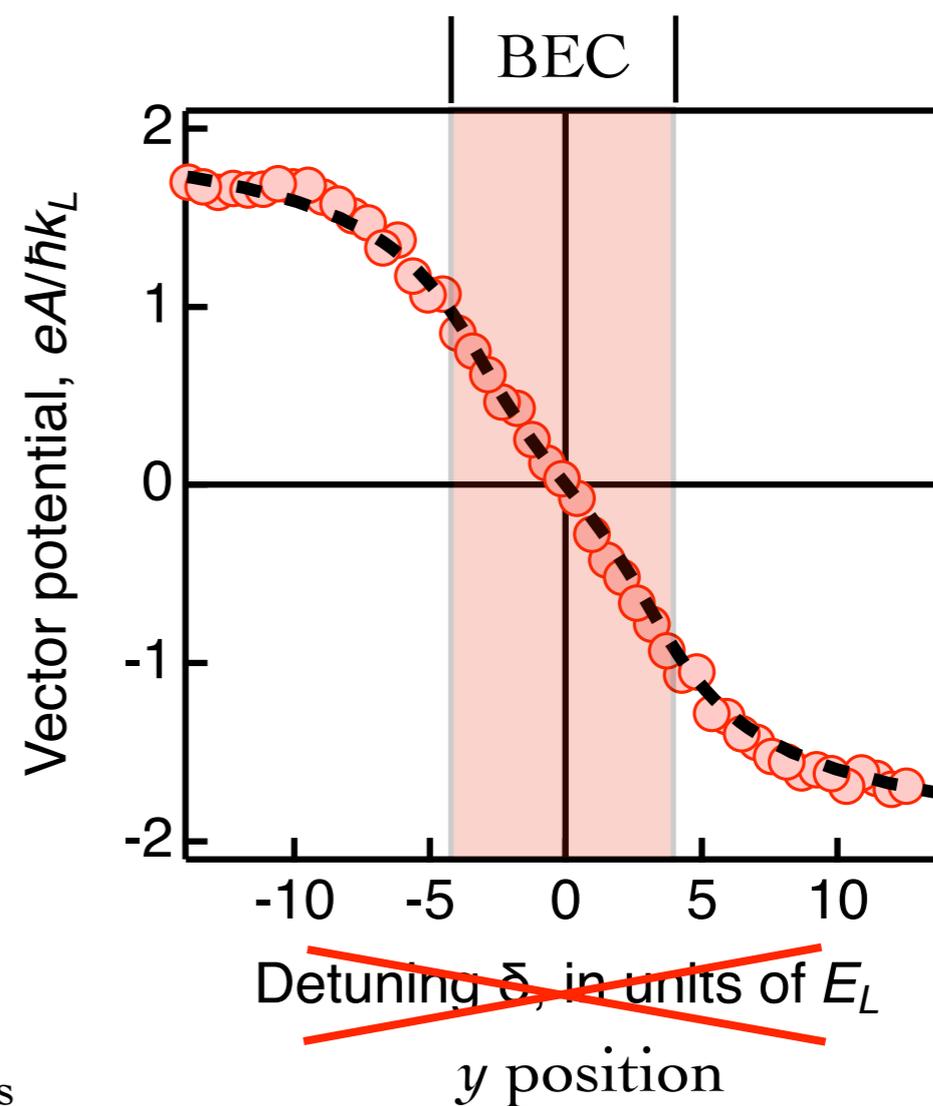
$$B = \nabla \times A$$



Transfer function

$$\hat{H} = \frac{\hbar^2}{2m} \left\{ \left[k_x - \frac{qA_x(\delta, \Omega)}{\hbar} \right]^2 + k_y^2 \right\} + V(\mathbf{x})$$

where $\delta(x, y, t)$ and $\Omega(x, y, t)$



References

Y.-J. Lin *et al*; Nature (2009)

Synthetic magnetic field

Loading procedure

Initial state

RF dressed state (RF on, ramp B to resonance)

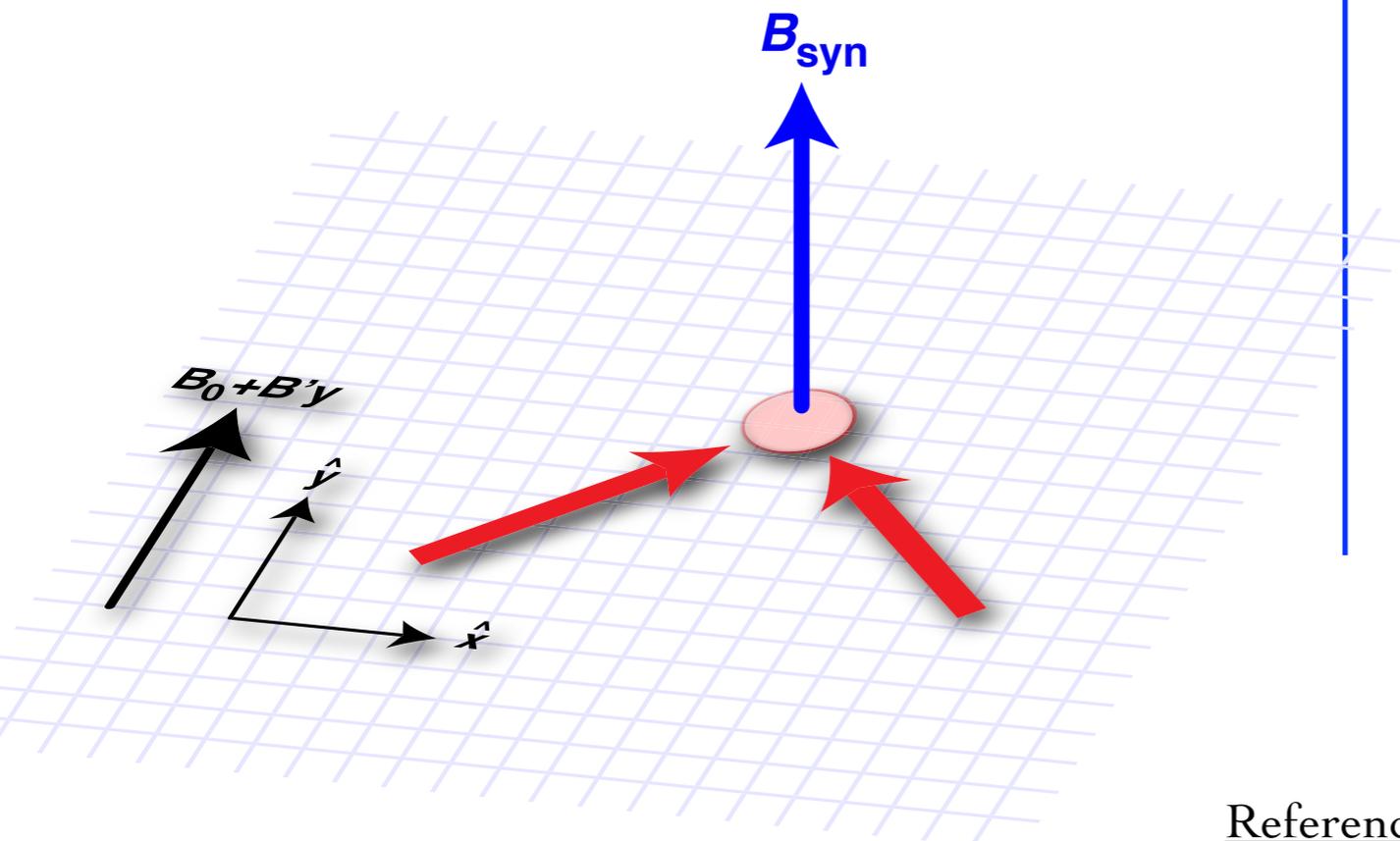
Raman + RF dressed state (Ramp Raman on)

Raman only dressed state (Ramp RF off)

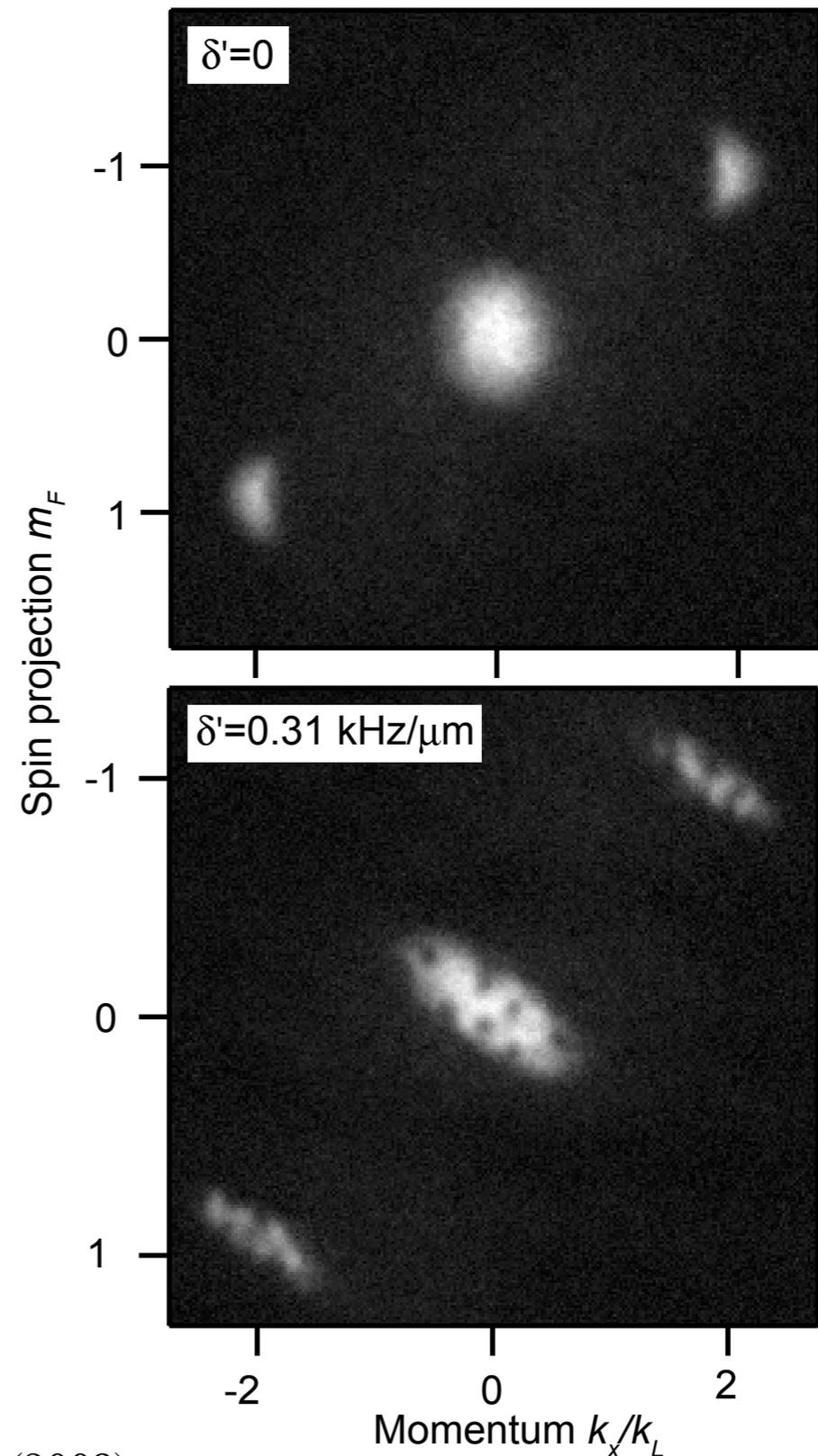
Ramp field gradient on (from 0 to 500 Hz/ μm)

Equilibrate for 500 ms

TOF imaging



Outcome



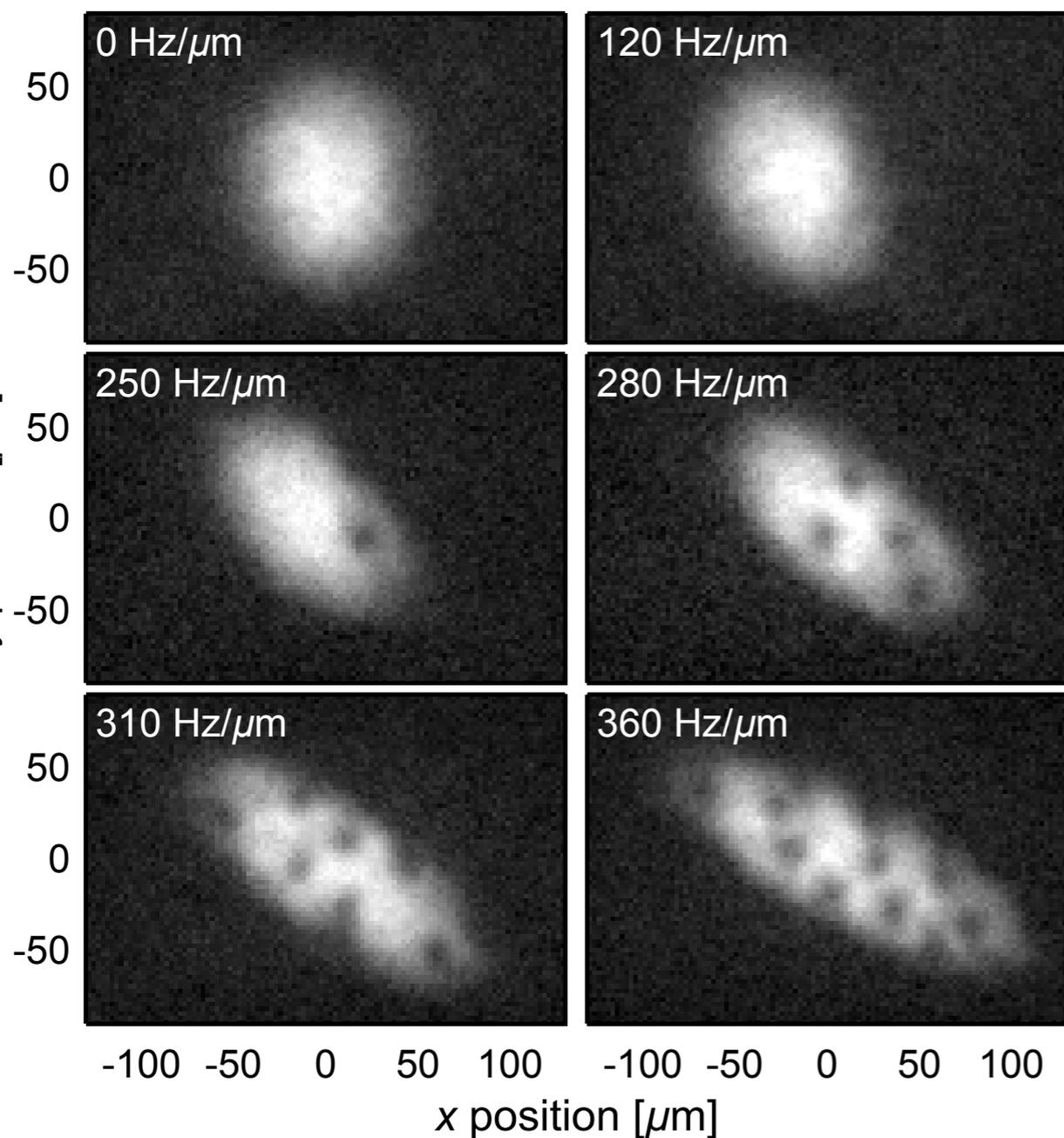
References

Y.-J. Lin *et al* Nature (2009)

Expected properties of BEC with fields

Vortex number

Spatial dependence gives magnetic fields and forces

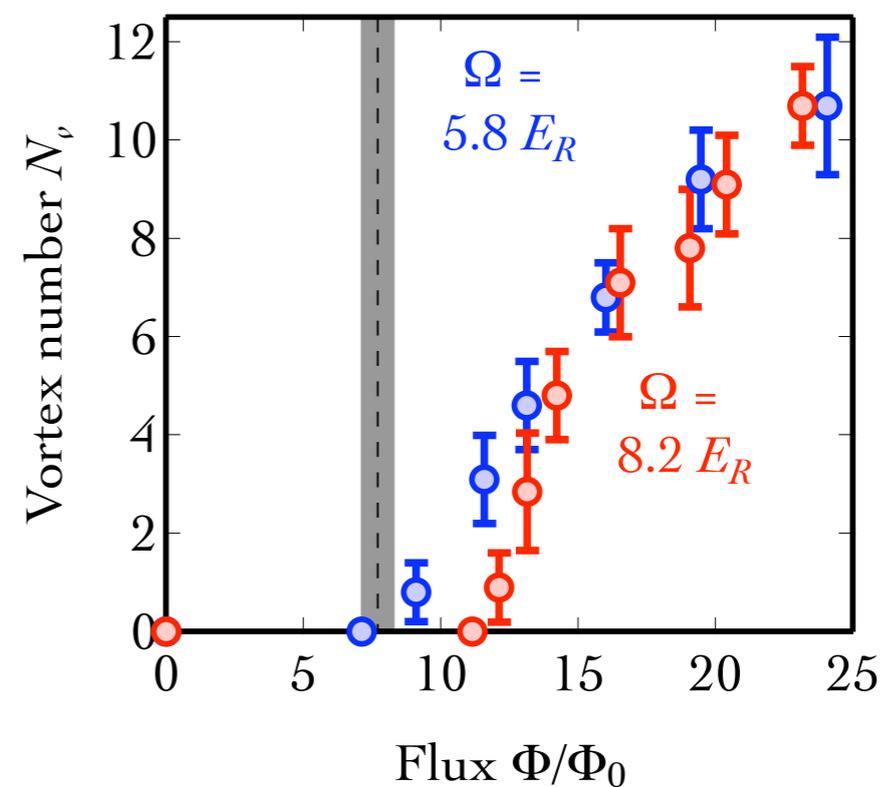
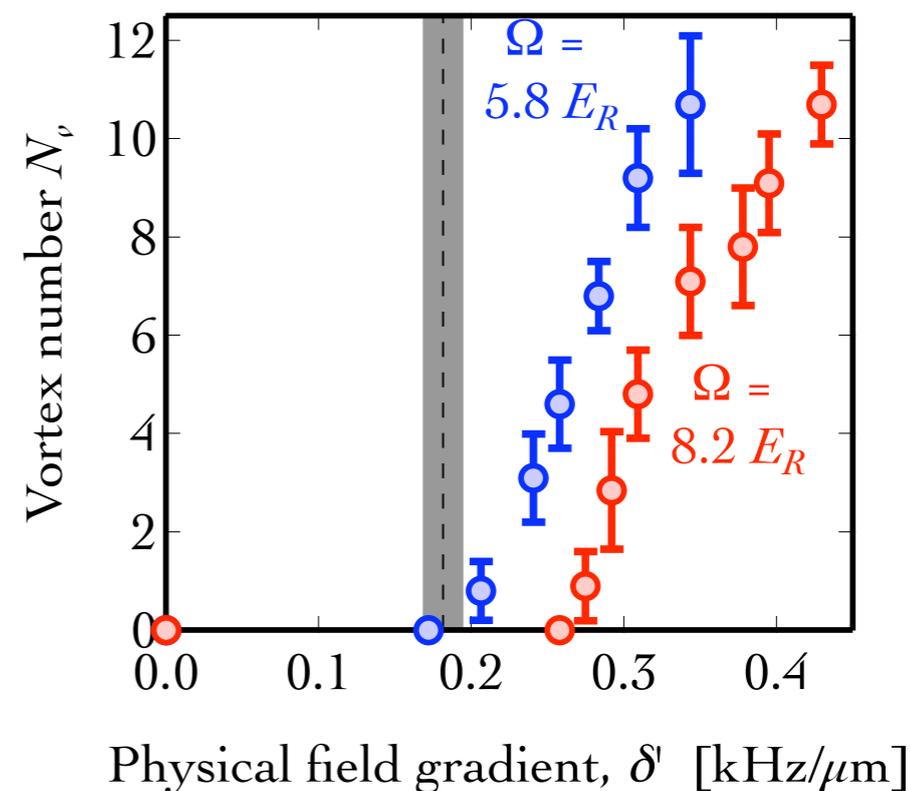


References

Y.-J. Lin *et al*; Nature (2009)

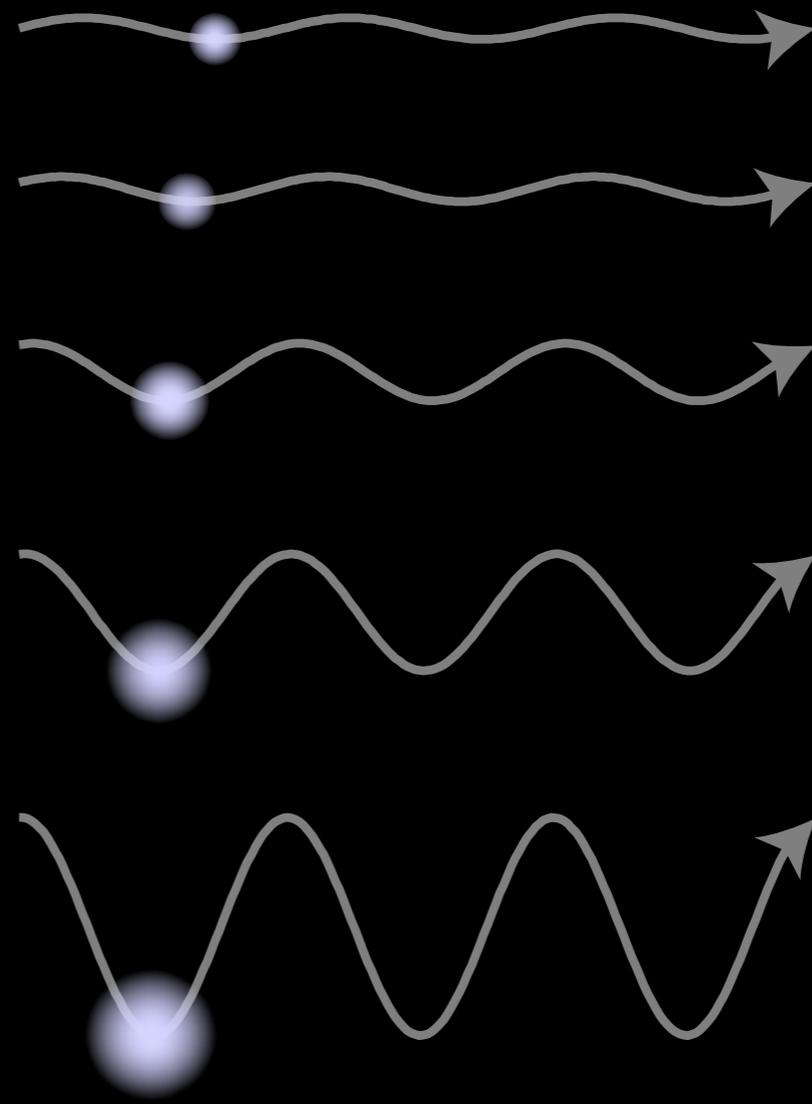
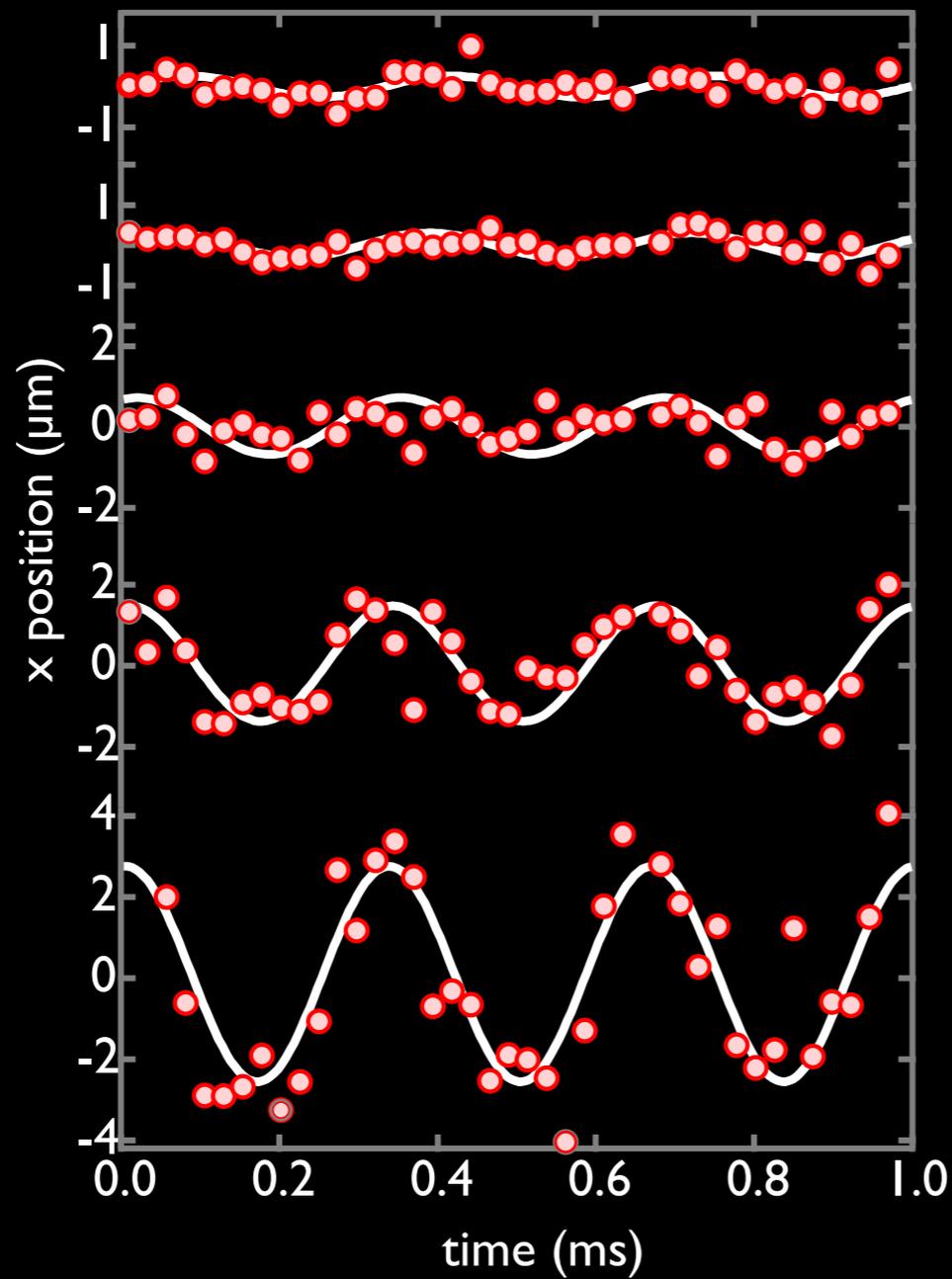
Extensive review: A.L. Fetter, RMP **81** 647 (2009)

Critical field for vortex formation



Atomic Zitterbewegung

$$\hat{H}_D = c\hat{p}_x\check{\sigma}_z + m_e c^2\check{\sigma}_x$$

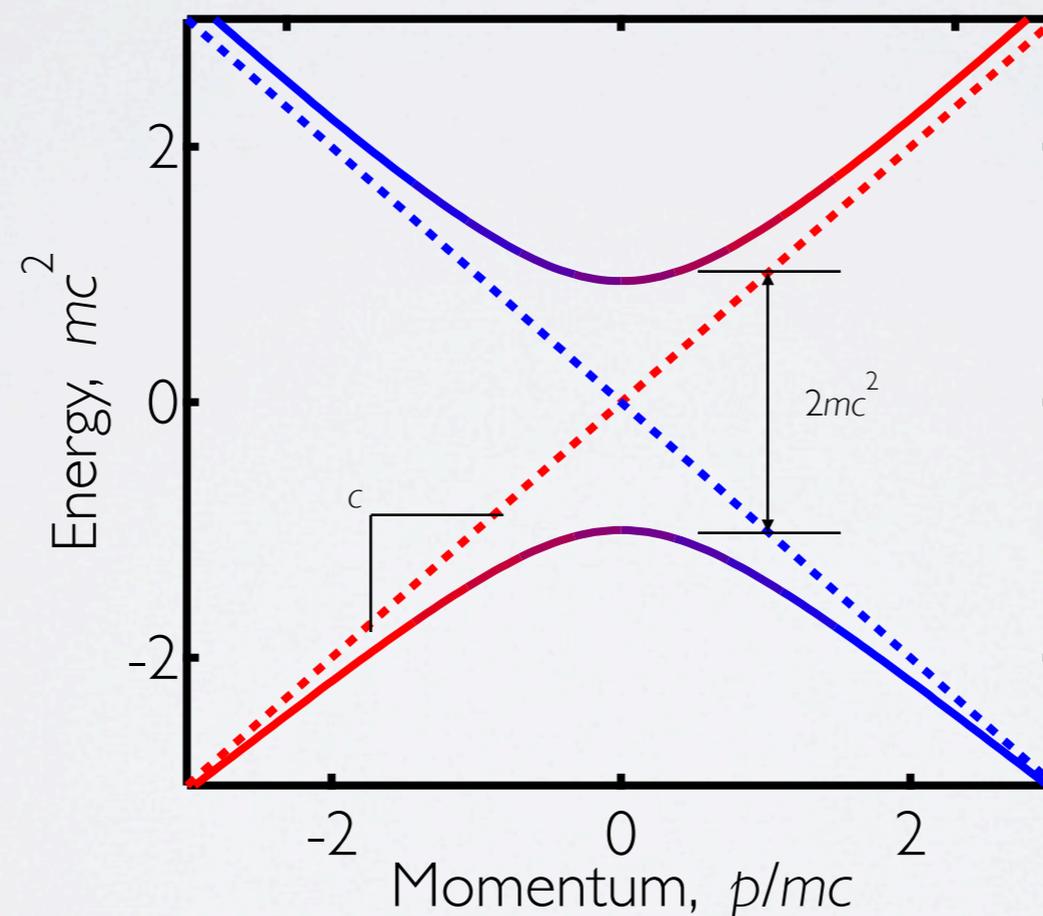


ZITTERBEWEGUNG: "JITTERING" MOTION

$$f = \frac{2mc^2}{h} = 2.5 \times 10^{20} \text{ Hz}$$

$$\delta x = \pm \frac{\lambda_C}{4\pi} = 2.4 \text{ pm}$$

$$\hat{H}_D \psi = \left(c\hat{p}_x \check{\sigma}_z + m_e c^2 \check{\sigma}_x \right) \psi,$$



Refs.

Schrodinger (1930)

David and Cserti PRB (2010)

Zitterbewegung

Expected tiny “jittering” of electrons (small and very fast)

$$f = \frac{2mc^2}{h} = 2.5 \times 10^{20} \text{ Hz}$$

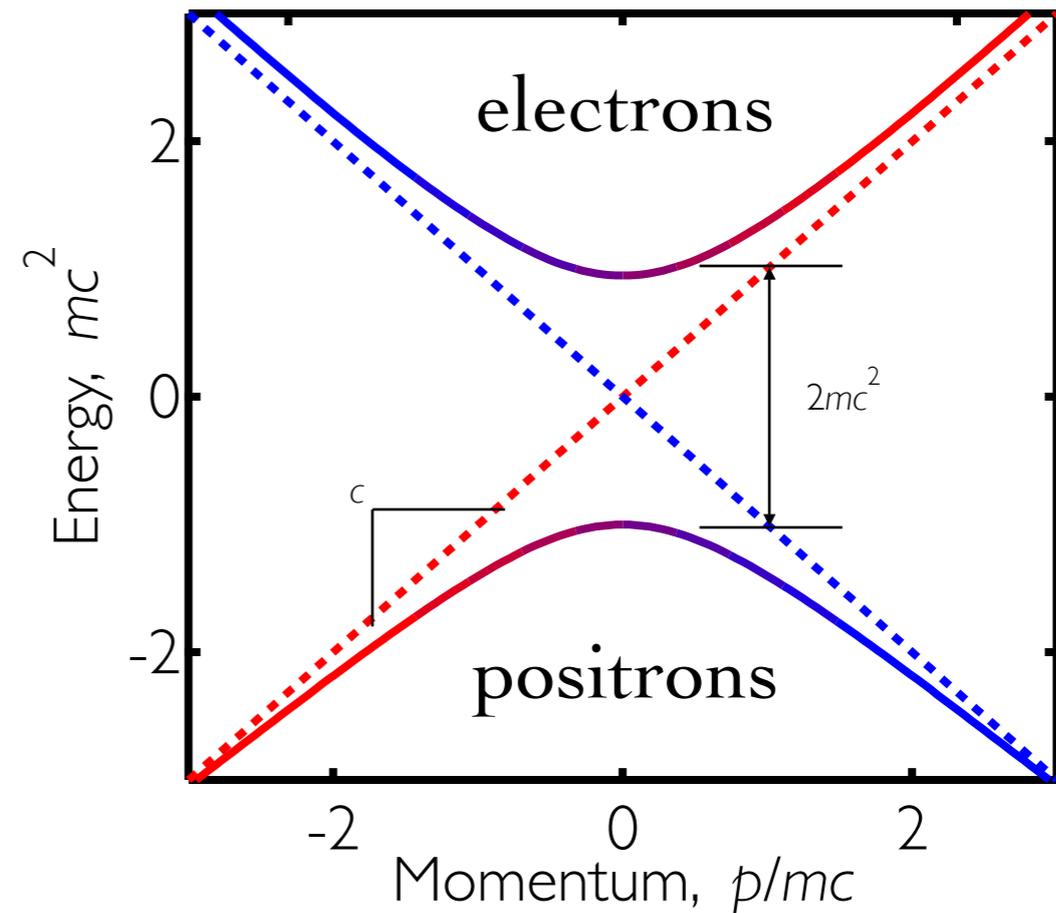
$$\delta x = \pm \frac{\lambda_C}{4\pi} = 2.4 \text{ pm}$$

$$\frac{dx}{dy} = \frac{1}{i\hbar} [\hat{x}, \hat{H}] = c\check{\sigma}_z = v$$

$$\frac{dv}{dt} = \frac{2mc^3}{\hbar} \check{\sigma}_y$$

$$\frac{d^2v}{dt^2} = \frac{4mc^4}{\hbar} \check{\sigma}_x \hat{k} - \left(\frac{2mc^2}{\hbar} \right)^2 v$$

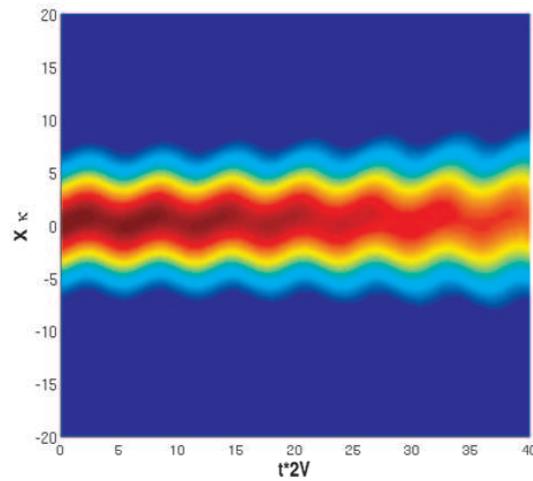
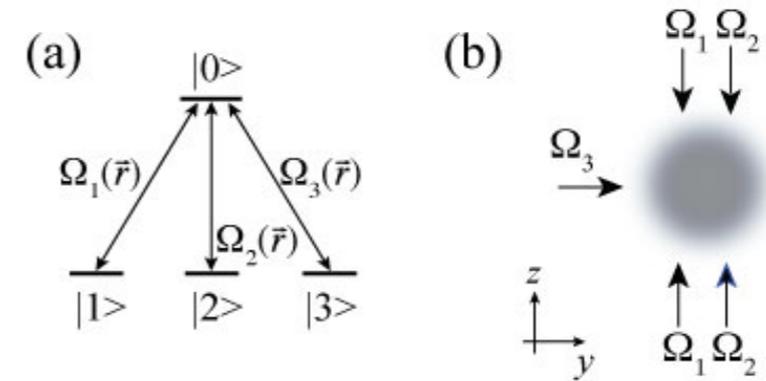
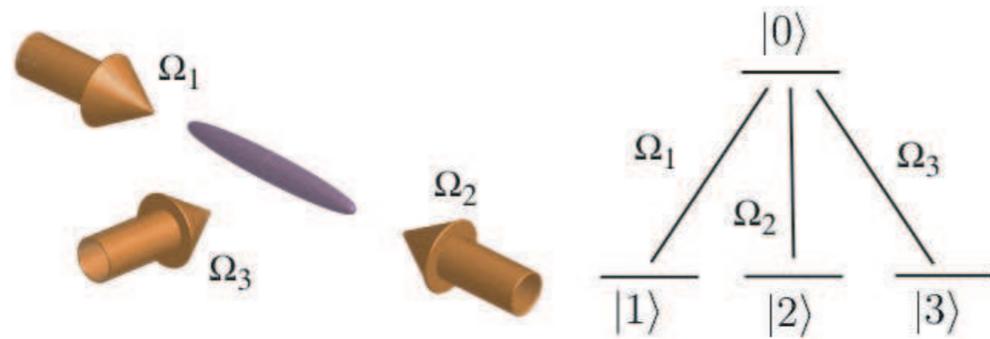
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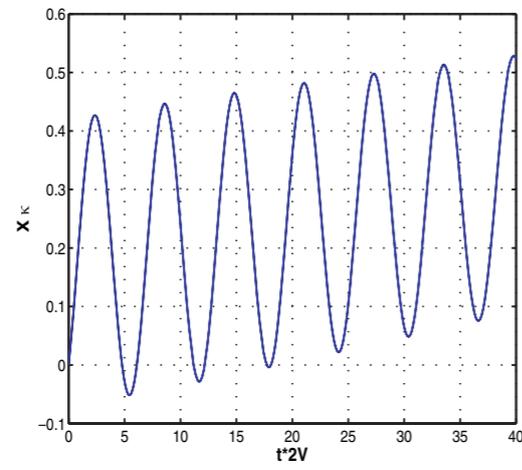
Ref.

David and Cserti PRB (2010)

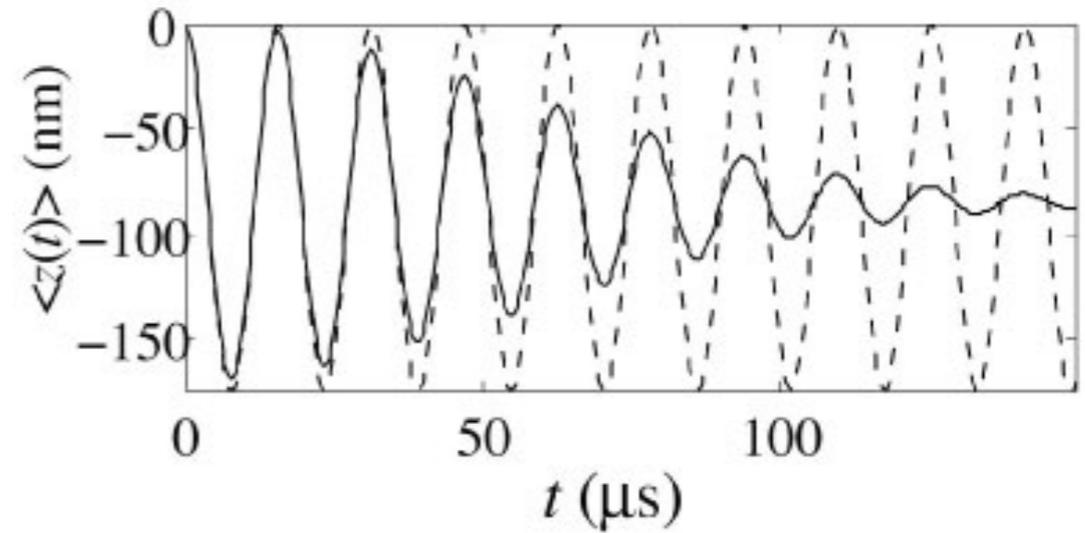
Proposals with cold atoms



(a)



(b)

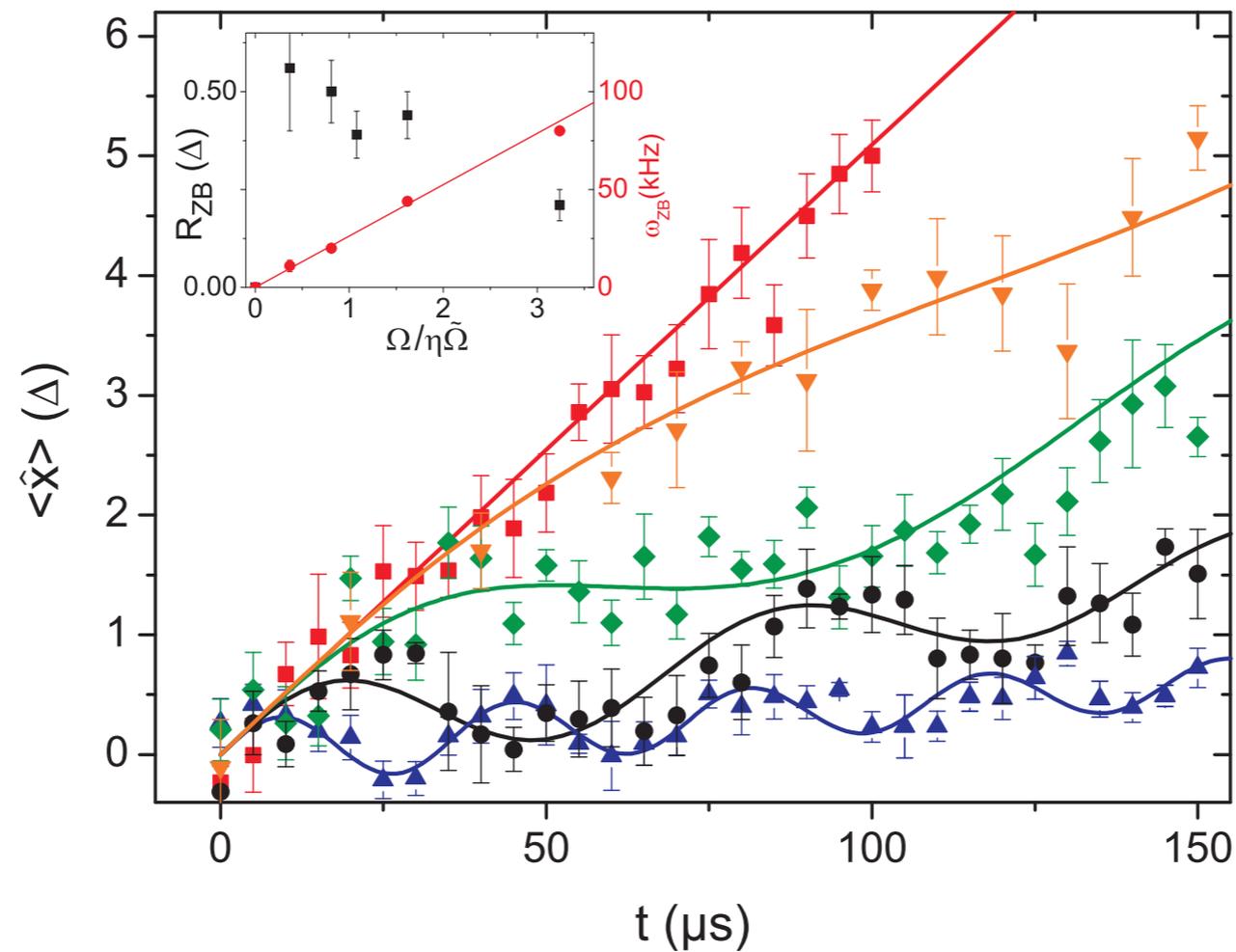


Ref.

- Vaishnav and Clark PRL (2008)
- Merkl, Zimmer, Juzeliūnas, and Ohberg EPL (2008)
- Zhang, Gong, and C. H. Oh arXiv:1208.3005 (2012)

Analog realized with individual atomic ions

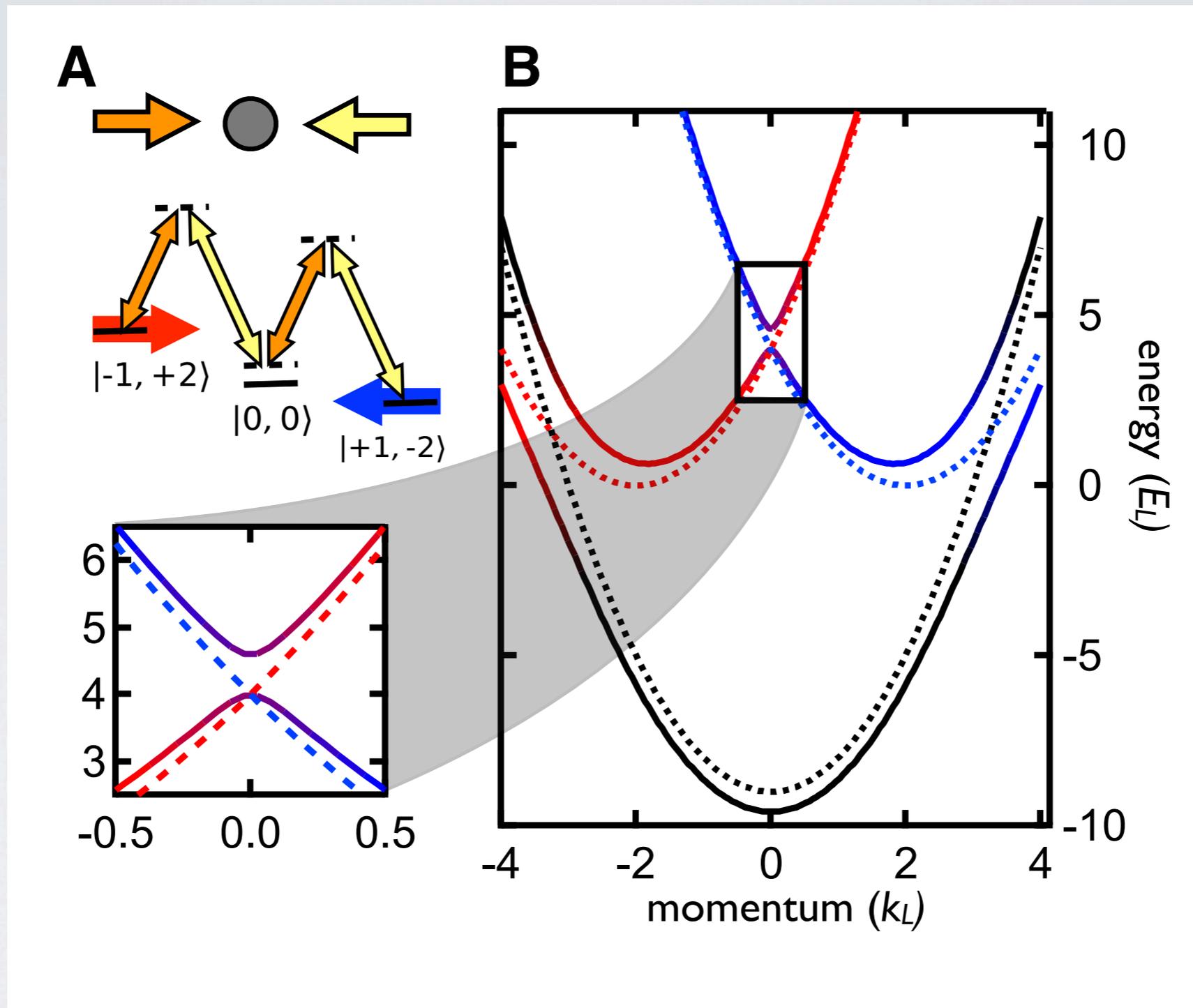
Measuring x and p quadratures in harmonic trap



Ref.

Gerritsma, et al. Nature (2010)

BEC EXPERIMENT



Theory Ref.

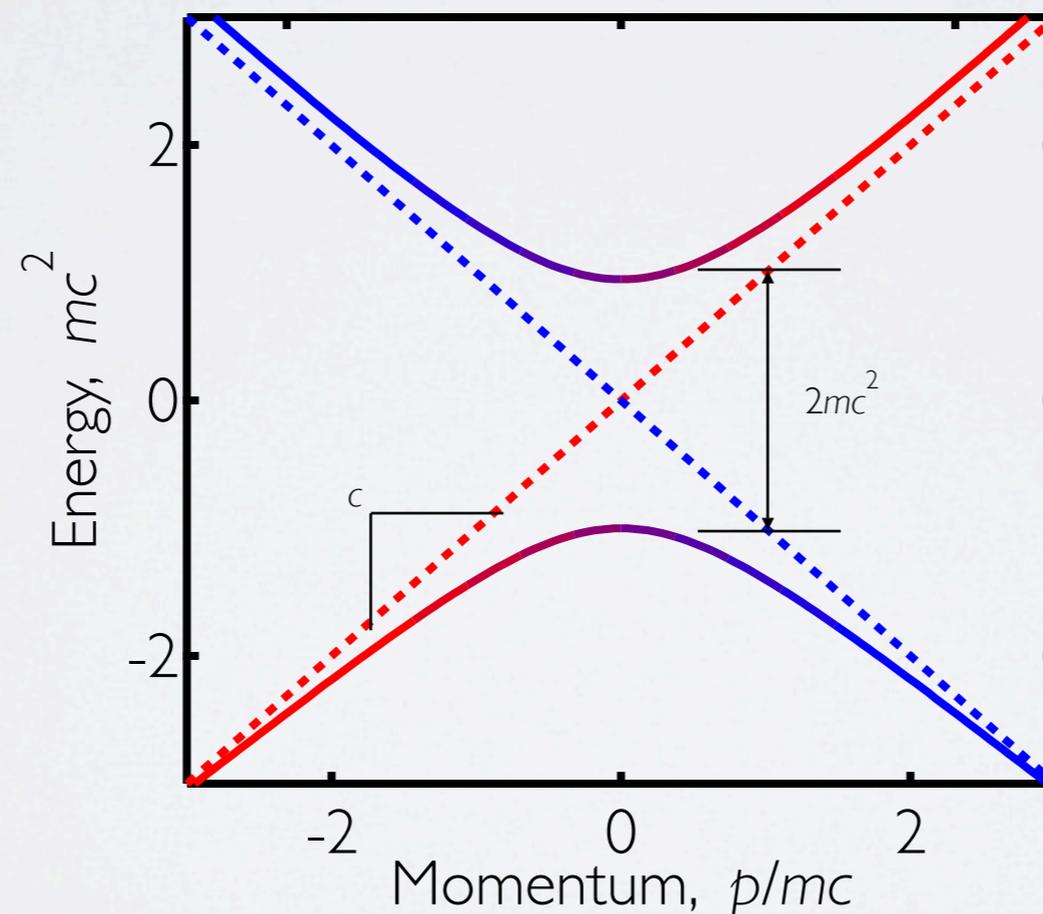
Zhang, Gong, and C. H. Oh arXiv:1208.3005 (2012)

ULTRASLOW RELATIVISTIC SYSTEM

$$2mc^2 = \hbar\Omega_2 \approx h \times 1 \text{ kHz}$$

$$c = \hbar 2k_r / m_{\text{Rb}} \approx 11 \text{ mm/s}$$

$$\hat{H}_D \psi = \left(c \hat{p}_x \check{\sigma}_z + m_e c^2 \check{\sigma}_x \right) \psi,$$



Ref.

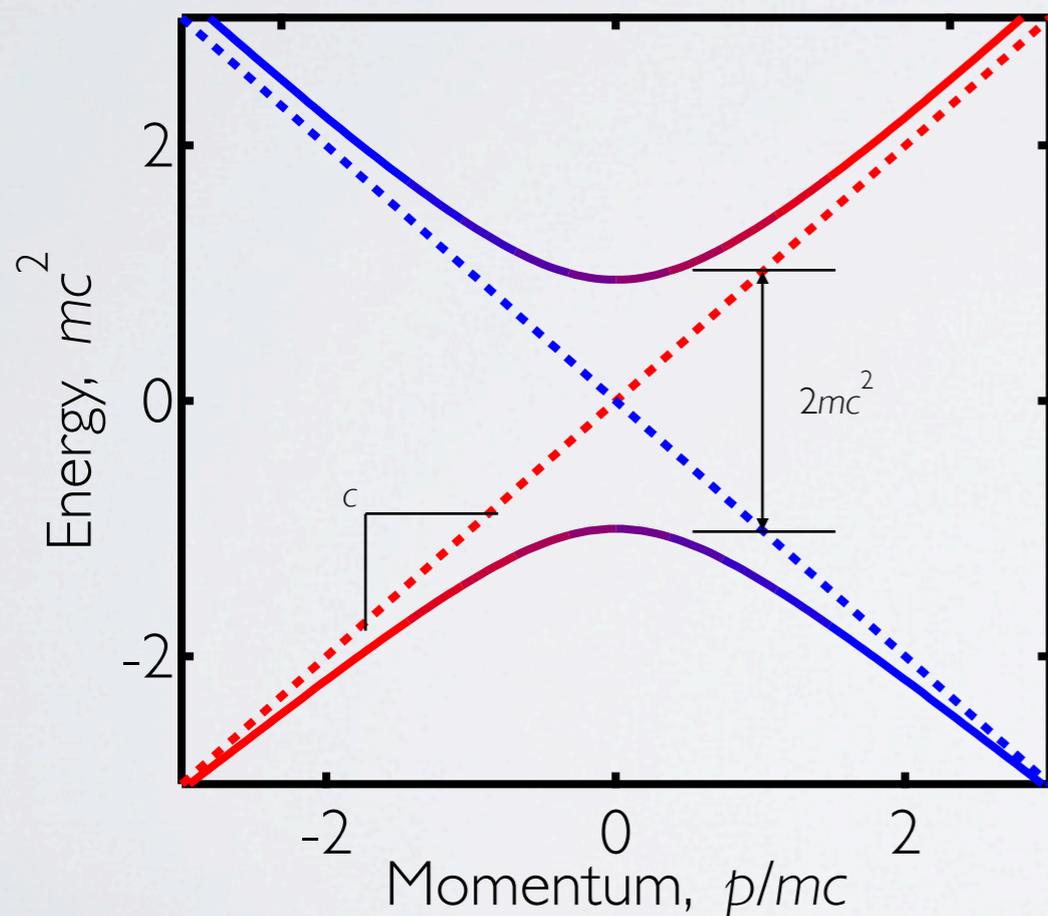
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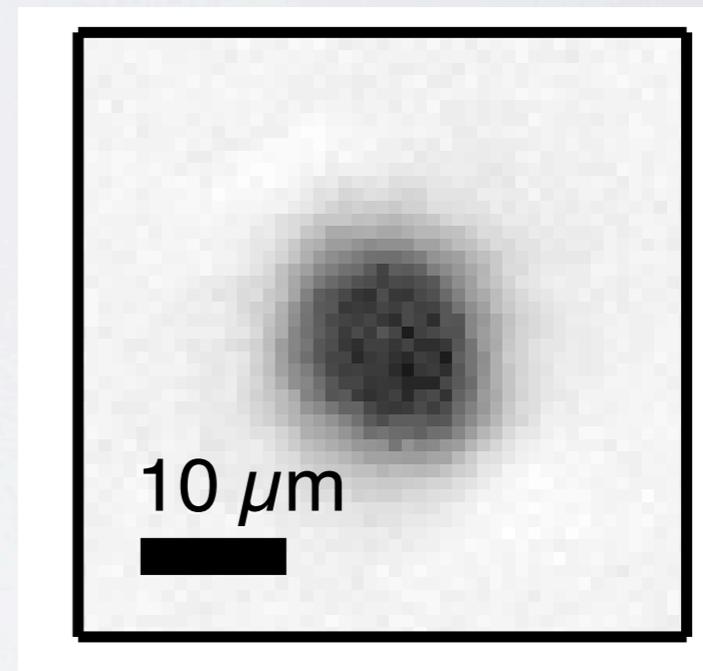
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Measure real position



Ref.

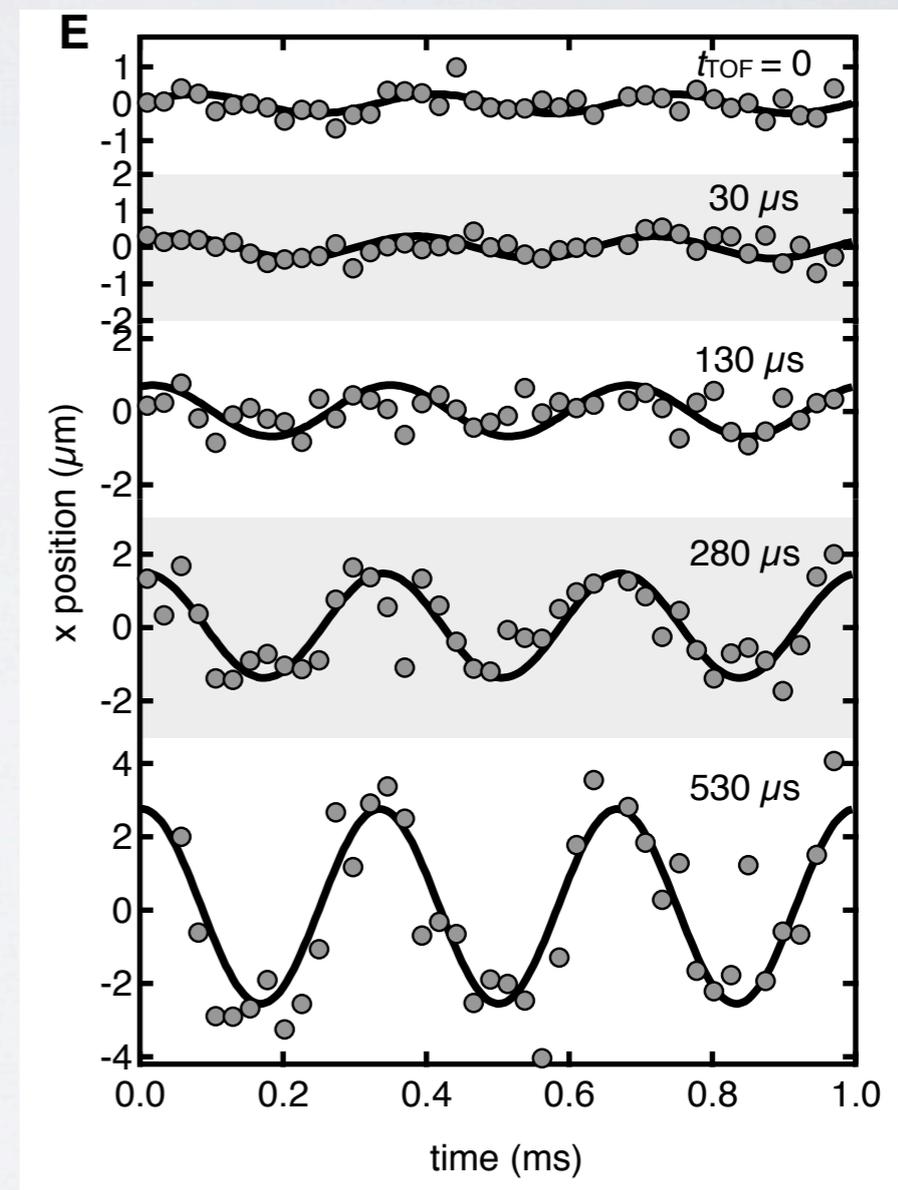
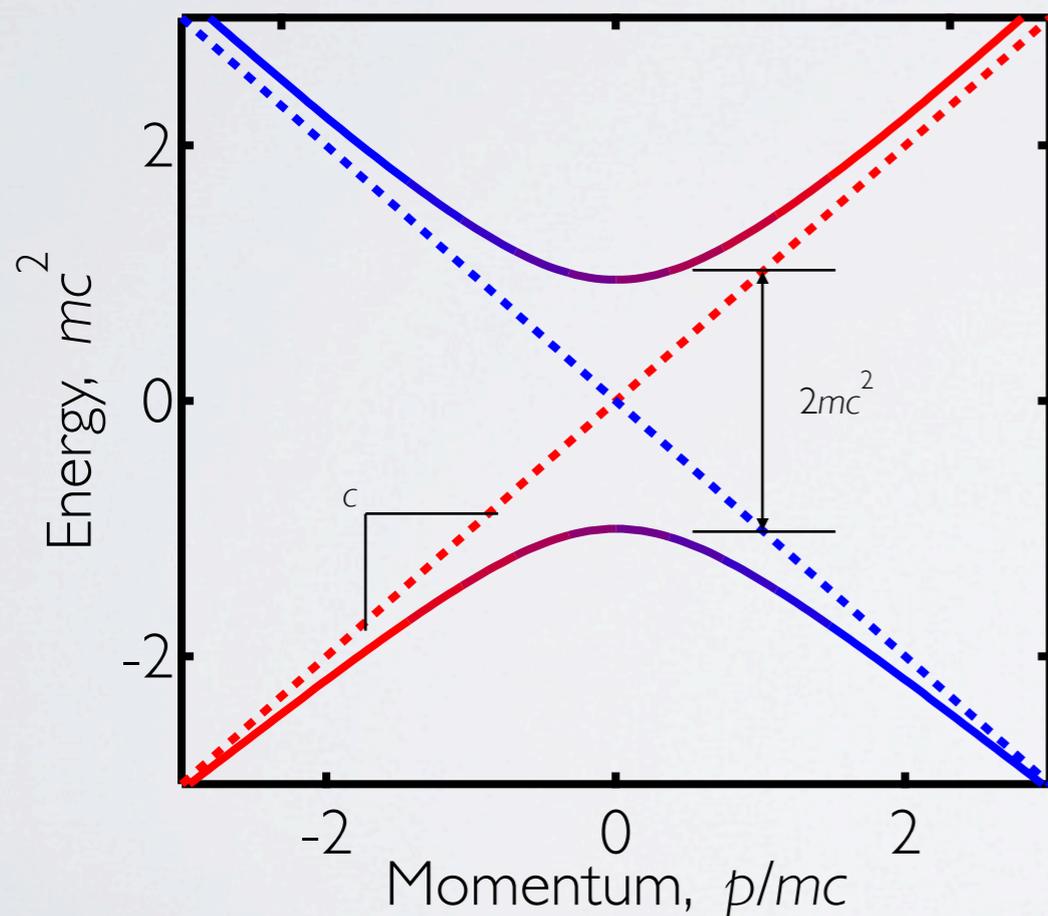
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×

∨

Ref.

L. J. Leblanc et al (in preparation)

ULTRASLOW RELATIVISTIC SYSTEM

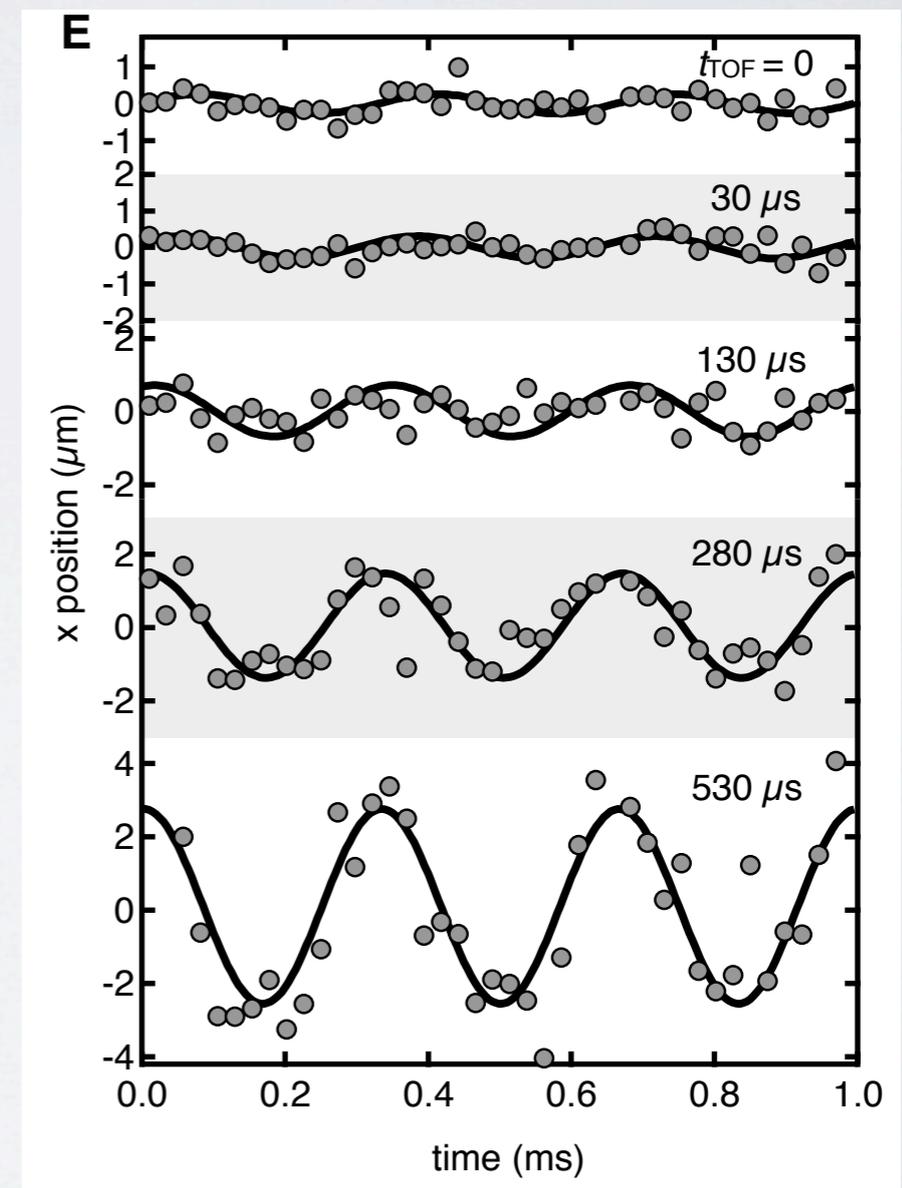
$$\langle x(t) \rangle = \frac{\lambda_C}{4\pi} \sin\left(\frac{4\pi c}{\lambda_C} t\right)$$

$$\langle \dot{x}(t) \rangle = c \cos\left(\frac{4\pi c}{\lambda_C} t\right)$$

$$2mc^2 = \hbar\Omega_2 \approx h \times 1 \text{ kHz}$$

$$c = \hbar 2k_r / m_{\text{Rb}} \approx 11 \text{ mm/s}$$

$$\lambda_C = \frac{h}{mc}$$



Ref.

L. J. Leblanc et al (in preparation)

ULTRASLOW RELATIVISTIC SYSTEM

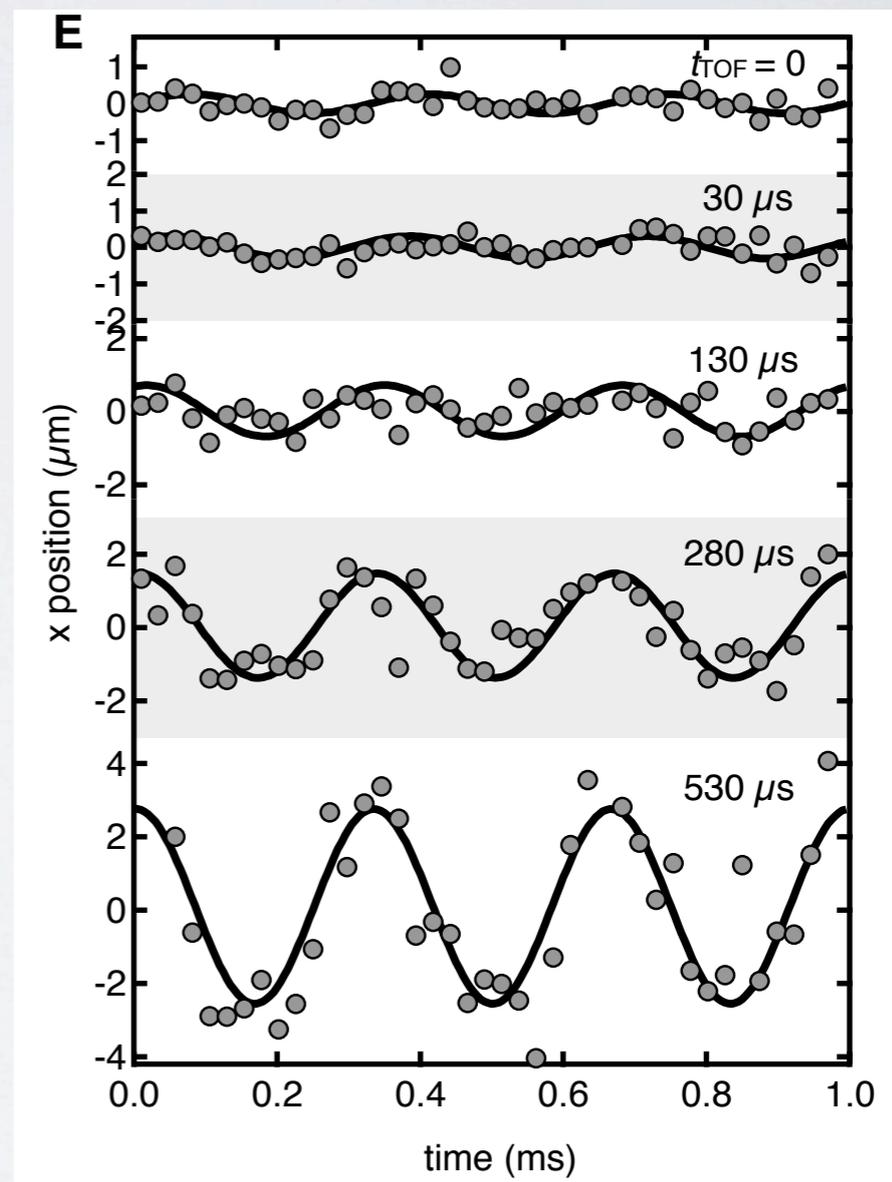
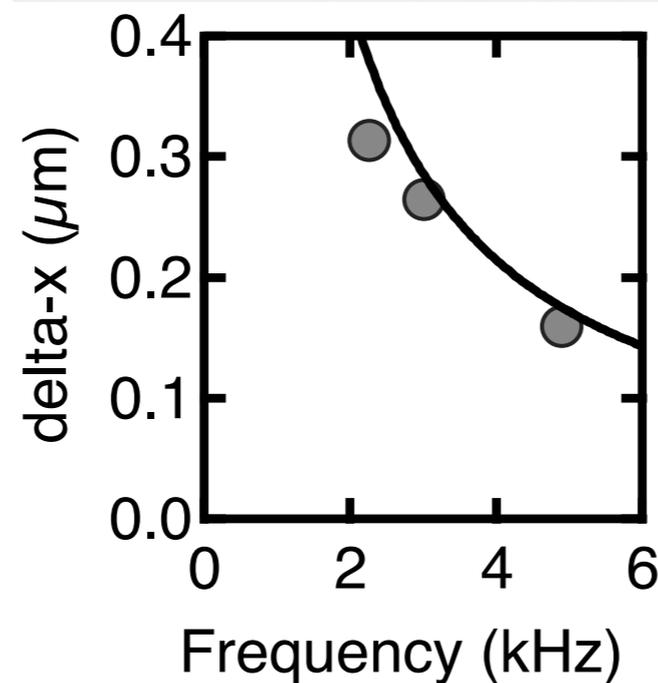
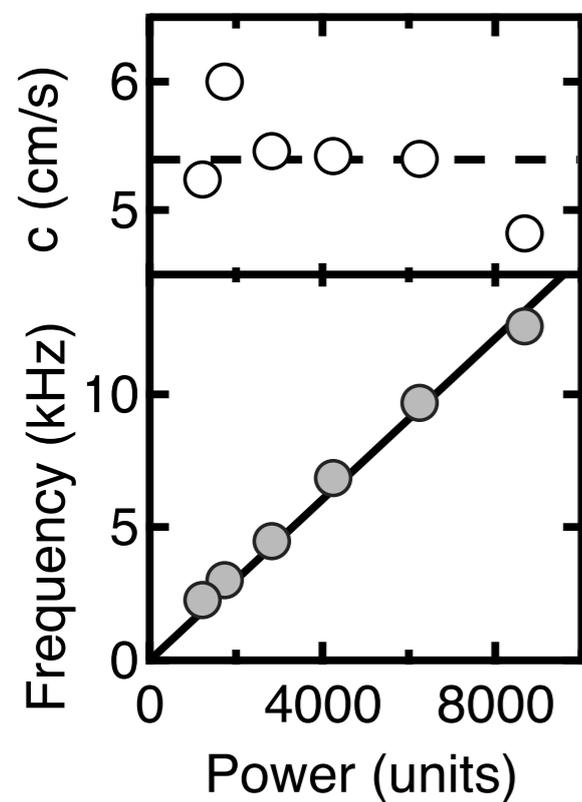
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$$\lambda_C = \frac{h}{mc}$$

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×

∨

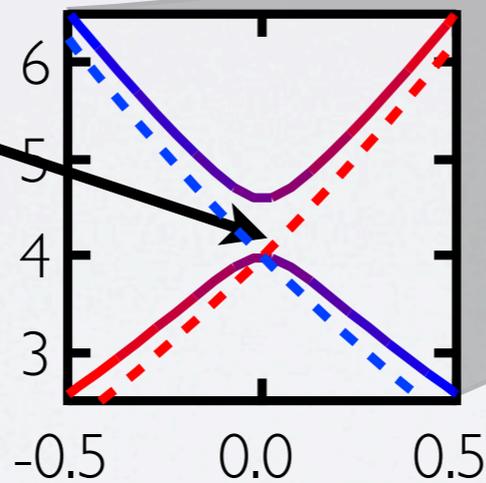
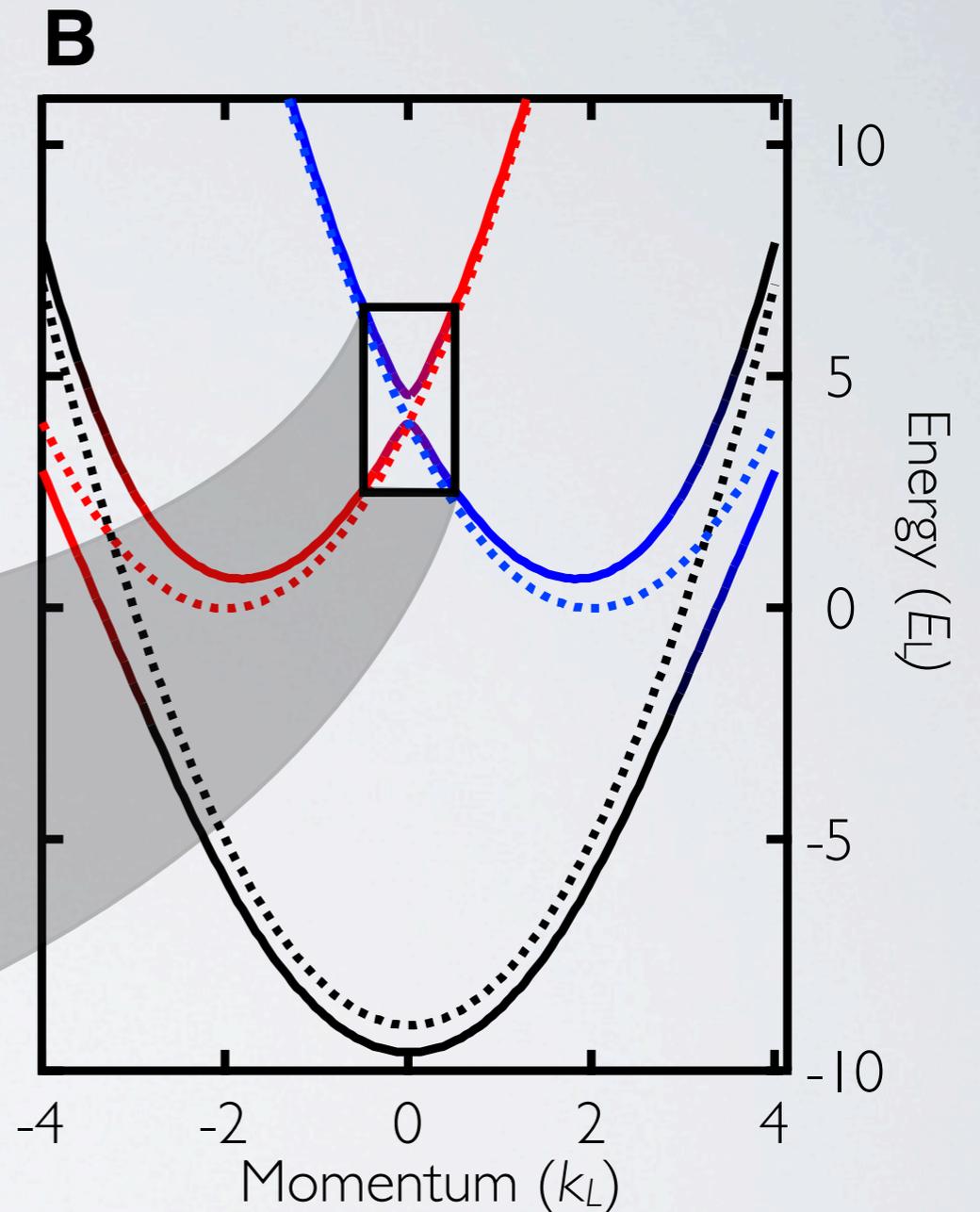
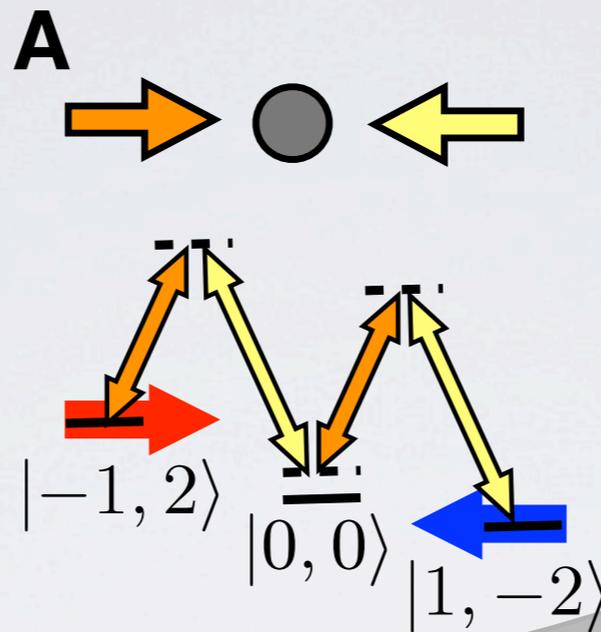
Ref.

L. J. Leblanc et al (in preparation)

PHYSICAL INTERPRETATION

Start in pure
“massless” state

$$|\psi(t=0)\rangle = |\uparrow, k = 2k_r\rangle$$



Evolve with coupling

$$|\psi(t)\rangle = \cos\left(\frac{mc^2}{\hbar}t\right) |\uparrow, k = 2k_r\rangle + i \sin\left(\frac{mc^2}{\hbar}t\right) |\downarrow, k = -2k_r\rangle$$

Ref.

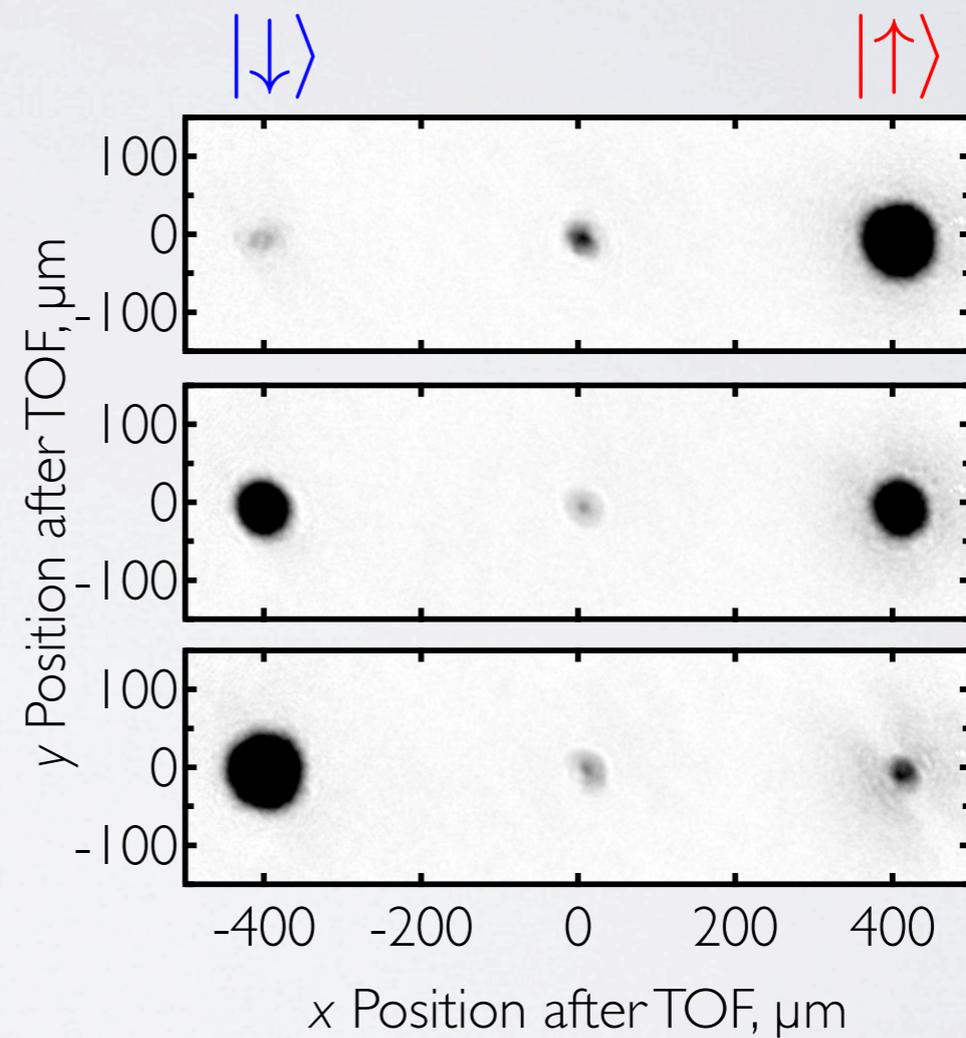
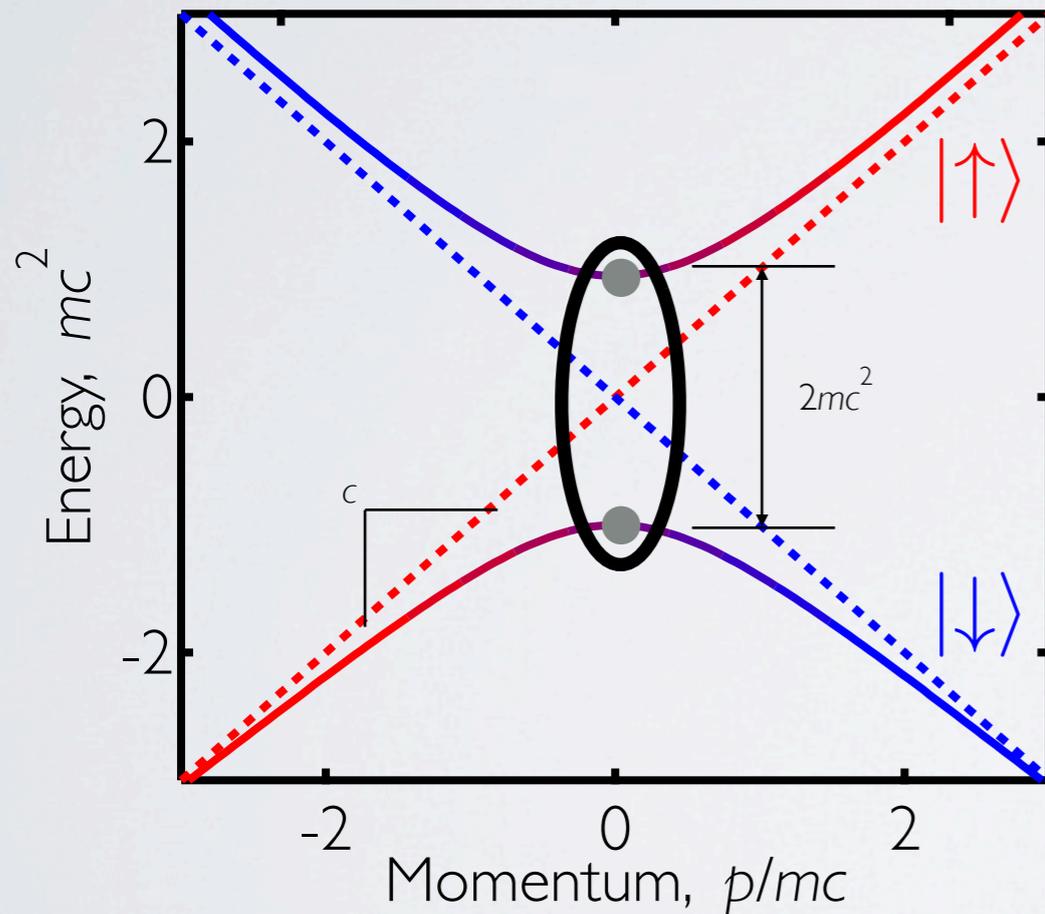
L.J. Leblanc *et al* (in preparation)

PHYSICAL INTERPRETATION

Rabi oscillations in a 2-level system

Two level system

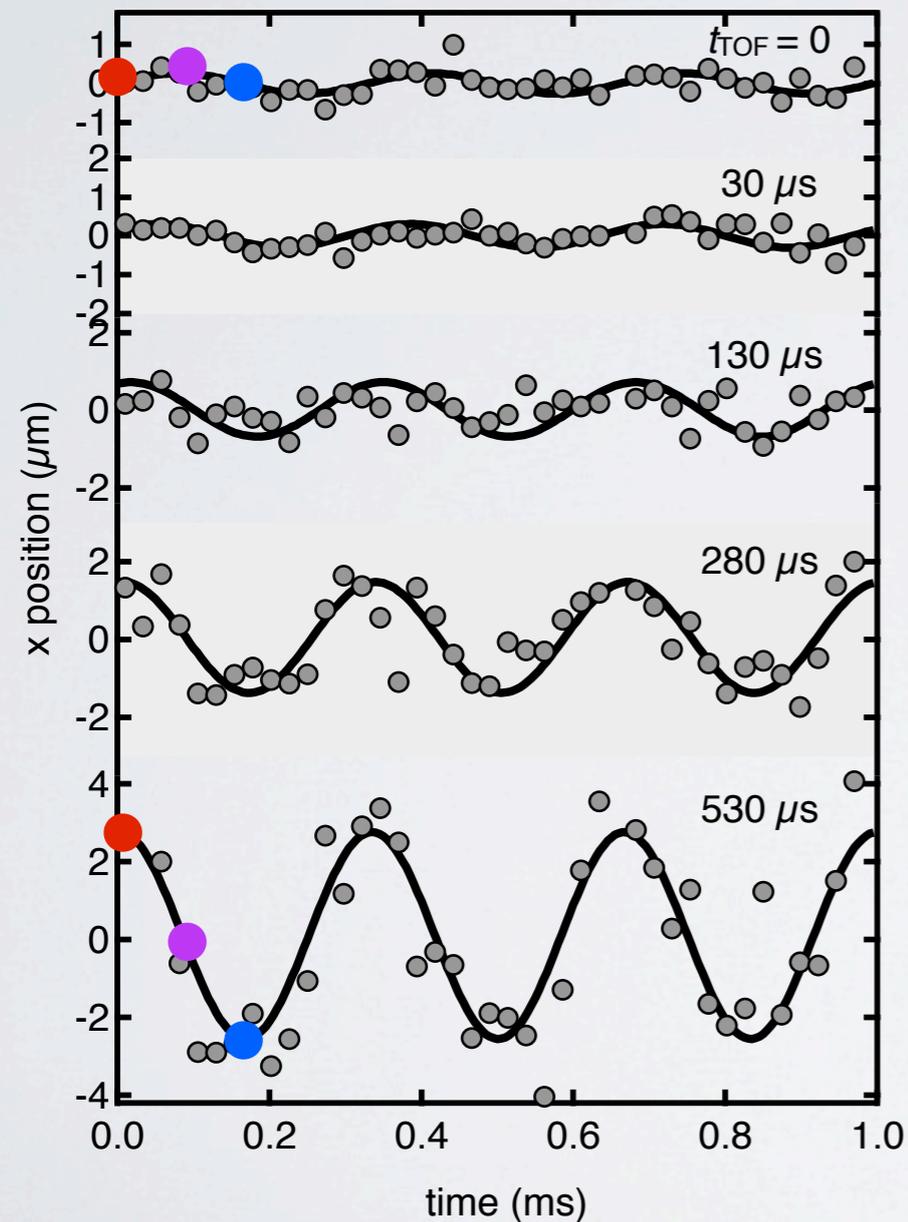
Time of flight images



PHYSICAL INTERPRETATION

Rabi oscillations in a 2-level system

Zitterbewegung



Time of flight images

